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Coulomb and Higgs branches from canonical singularities. Part I. Hypersurfaces with smooth Calabi-Yau resolutions

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ABSTRACT: Compactification of M-theory and of IIB string theory on threefold canonical singularities gives rise to superconformal field theories (SCFTs) in 5d and 4d, respectively. The resolutions and deformations of the singularities encode salient features of the SCFTs and of their moduli spaces. In this paper, we build on Part 0 of this series [1] and further explore the physics of SCFTs arising from isolated hypersurface singularities. We study in detail these canonical isolated hypersurface singularities that admit a smooth Calabi-Yau (crepant) resolution. Their 5d and 4d physics is discussed and their 3d reduction and mirrors (the magnetic quivers) are determined in many cases. As an explorative tool, we provide a `Mathematica` code which computes key quantities for any canonical isolated hypersurface singularity, including the 5d rank, the 4d Coulomb branch spectrum and central charges, higher-form symmetries in 4d and 5d, and crepant resolutions.

KEYWORDS: Conformal Field Models in String Theory, Extended Supersymmetry

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Contents

1	Introduction	1
2	Hypersurface singularities and SCFTs: a practical summary	4
2.1	Deformation and 4d Coulomb branch spectrum	4
2.2	Crepant resolution and 5d Coulomb branch	8
3	Link topology and higher-form symmetries	11
3.1	Topology of the link	12
3.2	Defects and higher-form symmetries in 5d from M-theory	13
3.3	Defects and higher-form symmetries in 4d from type IIB	15
4	Classification of canonical IHS and Mathematica code	16
4.1	Classification of canonical isolated hypersurface singularities in \mathbb{C}^4	16
4.2	Computing the basic data of any canonical IHS: Mathematica code	21
5	Smoothable models	23
5.1	Tabulating the smoothable models	23
5.2	The $\text{AD}[E_7, E_7]$ model	25
5.3	4d SCFTs: $D_p^b(G)$ -trinions	28
5.4	4d SCFTs with $c - a < 0$	28
5.4.1	Rank 2: $\{1, 1\}\{2, 5, 5, 5\}$	30
5.4.2	Rank 2: $\{1, 1\}\{2, 3, 8, 8\}$	31
5.4.3	Rank 3: $\{1, 1\}\{2, 3, 9, 9\}$	32
5.4.4	Rank 3: $\{1, 1\}\{3, 4, 4, 4\}$	33
5.4.5	Rank 3: $\{2, 1\}\{2, 7, 7, 3\}$	34
5.4.6	Rank 5: $\{1, 1\}\{2, 3, 10, 10\}$	35
5.4.7	Rank 5: $\{2, 1\}\{2, 3, 15, 7\}$	35
5.4.8	Rank 6: $\{1, 1\}\{2, 5, 6, 6\}$	36
5.5	4d SCFTs with $a = c$	37
5.5.1	Rank 6: $\{2, 1\}\{4, 5, 5, 2\}$	38
A	Crepantly resolvable geometries with smooth divisors and $b_3 = 0$	39
A.1	$\text{AD}[D_{2n}, D_{2n}]$ theories	40
A.1.1	Deformation and magnetic quiver	40
A.1.2	Resolution and 5d SCFT interpretation	41
A.2	Rank 1	41
A.3	Rank 3: $\{4, 1\}\{3, 3, 9, 2\}$	42
A.3.1	Deformations and magnetic quiver	42
A.3.2	Resolution and 5d SCFT	43
A.4	Rank 3: $\{12, 1\}\{7, 2, 3, 2\}$	43
A.4.1	Deformations and magnetic quiver	44

A.4.2	Resolution and 5d SCFT	44
A.5	Rank 6: $\{18, 1\}\{2, 3, 2, 5\}$	45
A.5.1	Deformations and 4d SCFT	45
A.5.2	Resolution and 5d SCFT	45
A.6	Rank 6: $\{12, 2\}\{13, 3, 2, 3\}$	46
A.6.1	Deformations and magnetic quiver	47
A.6.2	Resolution and 5d SCFT	47
A.7	Rank 9: $\{4, 1\}\{4, 2, 4, 4\}$	47
A.8	Rank 10: $\{10, 1\}\{5, 2, 2, 6\}$	49
B	Tables: smoothable models	50

1 Introduction

Supersymmetric quantum field theories (SQFT) can often be embedded within string theory by using either branes, geometry, or a mixture of both (distinct approaches are often related by string dualities). Over the years, string-theory techniques have led to a deeper understanding of the rich landscape of superconformal field theories (SCFT) in space-time dimension $d > 2$, and especially for $d > 4$. These SCFTs can also be analysed using SQFT methods. A typical strategy is to analyse the low-energy effective field theories obtained after breaking conformal invariance (explicitly or spontaneously). Here, we are interested in the geometric engineering approach, wherein conformal field theories correspond to certain geometric singularities, and the massive phases correspond to smoothing the singularity.

This paper is part of a series started in [1], wherein we explore the geometric engineering of 4d and 5d SCFTs at threefold isolated hypersurface singularities (IHS)

$$\mathbf{X} = \{F(x) = 0 \mid x \in \mathbb{C}^4\}. \quad (1.1)$$

We assume that \mathbf{X} is a canonical threefold singularity, which is essentially the Calabi-Yau condition. Consequently, Type II string theory or M-theory on this geometric background preserves 8 supercharges.¹ Then, we expect the existence of two distinct maps from geometries to theories in 4d and 5d, respectively

$$\text{IIB} : \mathbf{X} \mapsto \mathcal{T}_{\mathbf{X}}^{4\text{d}}, \quad \text{M-theory} : \mathbf{X} \mapsto \mathcal{T}_{\mathbf{X}}^{5\text{d}}. \quad (1.2)$$

On the one hand, the infrared limit of Type IIB string theory on \mathbf{X} is expected to give us a 4d $\mathcal{N} = 2$ SCFT [2, 3]. On the other hand, the infrared limit of M-theory on \mathbf{X} should similarly give rise to a 5d SCFT [4, 5]. The 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$ and the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5\text{d}}$ are closely related upon compactification to 3d, as discussed in [1, 6, 7].

¹This is a subtle point, since any discussion of supersymmetry implicitly assumes the existence of a Ricci-flat metric. Our discussion of the singularities will remain entirely algebraic, thus bypassing (and postponing) any deeper questions about the metric.

There are two types of smoothing of \mathbf{X} we may consider: deformations or crepant resolutions (i.e. resolutions which preserve the Calabi-Yau condition). This corresponds to turning on some complex structure moduli or some Kähler moduli, respectively. In Type IIB, families of deformation are ‘classical’, in that they are not subject to worldsheet nor D-brane instanton corrections. The general family of deformed singularities, denoted by $\hat{\mathbf{X}}$, gives us the Seiberg-Witten geometry on the extended Coulomb branch of the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$. Similarly, families of crepant resolutions $\pi : \tilde{\mathbf{X}} \rightarrow \mathbf{X}$, denoted by $\tilde{\mathbf{X}}$, are often ‘classical’ in M-theory, where they correspond to the real extended Coulomb branch of the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$. By contrast, deformations in M-theory and crepant resolutions in IIB are quantum-corrected, due to M2-brane instantons wrapping 3-cycles and to worldsheet, D1- and D3-brane instantons wrapping exceptional curves and divisors, respectively. These ‘quantum’ geometries in M-theory and Type IIB correspond to going onto the Higgs branch of $\mathcal{T}_{\mathbf{X}}^{5d}$ or $\mathcal{T}_{\mathbf{X}}^{4d}$, respectively. Note that this class of geometric engineering of SCFTs reverses the old QFT slogan about the Higgs branch being protected from quantum corrections (which is true for Lagrangian theories, but misleading otherwise): from the viewpoint of string theory, the Coulomb branch is ‘classical’ and the Higgs branch is ‘quantum’.

Topologically, the crepant resolutions of canonical IHS are rather more subtle than their deformations. While the latter always have the homotopy type of a bouquet of 3-spheres, crepant resolutions of canonical IHS can have a rich topology, including compact 2-, 3- and 4-cycles. Moreover, crepant resolutions, even if they exist, are not always smooth. We might have residual terminal singularities in the resolved threefold $\tilde{\mathbf{X}}$, or it might happen that $\tilde{\mathbf{X}}$ is smooth but the exceptional divisors themselves are not. In this paper, we will focus on isolated hypersurface singularities whose crepant resolutions are smooth. They form a small but important subset of all canonical IHS, to which we can apply a number of techniques to study the corresponding SCFT moduli spaces.

Part of the motivation for this paper is to explain some simple tools to study the key properties of deformed and resolved canonical hypersurface singularities (details which were omitted in [1]). All the necessary computations can be easily implemented on a computer. We attached an ancillary *Mathematica* notebook `Basic-Data-IHS.nb` in the supplementary material of this paper. The code takes as an input any canonical isolated hypersurface singularity, whose classification [3, 8, 9] will be streamlined and updated in section 4.1 below. (See also [10] for a related discussion.) We then provide two functions:

1. `BasicIHSdata` $[F(x_1, x_2, x_3, x_4), \{x_1, x_2, x_3, x_4\}]$: this outputs all the basic characteristics of the 4d and 5d SCFTs such as ranks, flavor ranks, dimensions of the Higgs branches, 4d Coulomb branch spectrum and central charges a and c , and the higher form symmetries.
2. `SmoothOperator` $[F(x_1, x_2, x_3, x_4), \{x_1, x_2, x_3, x_4\}]$: determines a resolution of the hypersurface singularity.

This computational tool will hopefully be useful to the practitioner.

We will then apply these methods to explore hypersurface singularities that admit a crepant resolution that is smooth (without residual terminal singularities). It is interesting

to organise the discussion in terms of the 5d rank, r , which is the number of exceptional divisors in the crepant resolution. We scanned through a large database of canonical IHS, which included a large number theories of low r . While it appears that there exists an infinite number of canonical IHS at any r (certainly, this is true at small r), we would like to put forward the following:

Conjecture (1). For any $r > 0$, there exists a finite number of canonical IHS whose crepant resolution is smooth.

We call such a singularity with a smooth resolution a ‘smoothable model’. In this paper, we list all smoothable models for $r \leq 4$ — this list is complete to the best of our current knowledge. For $r \geq 5$, the scan through our dataset becomes impractical, but we can focus on a subset of ‘fully smooth models’, which are the smoothable models with exceptional divisors which are themselves smooth. We can scan for such models using the following:

Conjecture (2). For any smoothable model with smooth exceptional divisors, the associated 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ has an integral CB spectrum (i.e. the conformal dimensions of the CB operators satisfy $\Delta \in \mathbb{Z}$).

This conjecture is based on observation in many examples. If true, this is a remarkable mathematical relation between the smoothness of the exceptional divisors and the spectrum of the deformed singularity. Assuming Conjecture (2), we can efficiently explore the space of all ‘fully smooth’ models up to $r = 10$, by first restricting to models with an integral 4d CB spectrum.

For all these smoothable models, we performed a detailed analysis of their 4d and 5d properties, including a determination of their magnetic quivers or quiverines whenever possible. The use of magnetic quivers (MQ) for 5d SCFTs was initiated in [11], where they were motivated by brane-web constructions. The MQ associated to a 5d SCFT is a 3d $\mathcal{N} = 4$ theory whose Coulomb branch is isomorphic to the Higgs branch of the 5d SCFT. Magnetic quivers have been a very powerful tool to explore the moduli spaces of superconformal theories with 8 supercharges, with most approaches based on brane-webs [12–24]. From a geometric perspective, the magnetic quivers were studied e.g. in [1, 6, 25–27].

Further studies of more generic canonical singularities and of their 4d and 5d interpretation, as well as a more in-depth discussion of the 3d electric and magnetic quiver(in)s, will be presented in part 2 of this series [7].

This paper is organised as follows. In section 2, we start with a lightning review of the geometry of isolated hypersurface singularities and of their relation to 5d and 4d SCFTs. In section 3, we discuss the topology of the link of the singularity, and its interpretation in terms of higher-form symmetries of the 5d and 4d SCFTs. In section 4, we review the classification of canonical IHS and explain how to use the ancillary `Mathematica` code. In section 5, the (fully) smoothable models of rank $r \leq 10$ are discussed in detail, highlighting some examples with noteworthy properties, such as the models with 4d central charges satisfying $a > c$ and $a = c$, and a fully resolvable model that corresponds to the $\text{AD}[E_7, E_7]$ SCFT in 4d, for which we propose a 5d IR quiver gauge theory. An appendix contains details on other smoothable models, including the infinite series of type $\text{AD}[D_{2n}, D_{2n}]$.

2 Hypersurface singularities and SCFTs: a practical summary

In this section, we review key aspects of the geometry of canonical threefold hypersurface singularities, \mathbf{X} , and of its relation to the SCFTs $\mathcal{T}_{\mathbf{X}}^{4d}$ and $\mathcal{T}_{\mathbf{X}}^{5d}$. By definition, a canonical singularity [28] admits a resolution $\pi : \tilde{\mathbf{X}} \rightarrow \mathbf{X}$ such that

$$K_{\tilde{\mathbf{X}}} = \pi^* K_{\mathbf{X}} + \sum_k a_k S_k, \quad S_k \subset \pi^*(0), \quad (2.1)$$

with $a_k \in \mathbb{R}$ such that $a_k \geq 0, \forall k$. Here, K_X denotes the canonical divisor of X , and S_k denote the exceptional divisors. A canonical singularity is called terminal if $a_k > 0, \forall k$.

In this paper, we focus on canonical singularities that admit a (complete) crepant resolution — that is, $a_k = 0, \forall k$.² We also focus exclusively on quasi-homogeneous isolated hypersurface singularities (IHS) in \mathbb{C}^4 ,

$$\mathbf{X} \cong \{F(x) = 0 \mid x \in \mathbb{C}^4\}, \quad (2.2)$$

which are defined by a single quasi-homogeneous polynomial with a unique critical point at the origin — $F \in \mathbb{C}[x_1, x_2, x_3, x_4]$ such that $dF(x) = F(x) = 0$ if and only if $x = 0$.

2.1 Deformation and 4d Coulomb branch spectrum

Consider the quasi-homogeneous IHS \mathbf{X} , with the \mathbb{C}^* -action:

$$F(\lambda^q x) = \lambda F(x), \quad \text{with} \quad x_i \rightarrow \lambda^{q_i} x_i, \quad i = 1, \dots, 4, \quad (2.3)$$

for some scaling weights $q_i \in \mathbb{Q}_{>0}$. Let us recall key properties of the deformation of such singularities [29], and its interpretation in terms of the Coulomb branch of the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$.

Canonical singularity condition. The quasi-homogeneous IHS \mathbf{X} is canonical if and only if the weights q_i are such that:

$$\sum_{i=1}^4 q_i > 1. \quad (2.4)$$

If this inequality is not satisfied, \mathbf{X} is worse than canonical. This condition can also be written in terms of the reduced central charge [30]:

$$\hat{c} \equiv \sum_{i=1}^4 (1 - 2q_i) < 2. \quad (2.5)$$

Here, \hat{c} is the reduced central charge of the 2d $\mathcal{N} = (2, 2)$ LG model with four fields x_i and superpotential $W = F(x)$.

²In previous works [1, 6], we studied canonical singularities without crepant resolution, which are either terminal singularities or singularities that admit a *partial* crepant resolution ($a_k = 0$ for some but not all k).

Milnor ring and multiplicity. A deformation of an hypersurface singularity $\mathbf{X} \cong \{F = 0\}$ is a (smooth) hypersurface of the generic form:

$$\widehat{F}(x) = F(x) + \sum_{l=1}^{\mu} t_l x^{\mathbf{m}_l} = 0, \quad (2.6)$$

where $\{x^{\mathbf{m}_l}\}_{l=1}^{\mu}$ is a basis for the so-called Milnor ring of the singularity:

$$\mathcal{M}(F) = \mathbb{C}[x_1, x_2, x_3, x_4] / \mathcal{J}, \quad \mathcal{J} \equiv (\partial_{x_1} F, \dots, \partial_{x_4} F). \quad (2.7)$$

The integer

$$\mu \equiv \dim \mathcal{M}(F), \quad (2.8)$$

is the multiplicity of the singularity (also known as the Milnor number). For quasi-homogeneous singularities, $\mathcal{M}(F)$ has a basis in terms of μ distinct monomials:

$$x^{\mathbf{m}_l} \equiv x_1^{\mathbf{m}_{l,1}} \cdots x_4^{\mathbf{m}_{l,4}}, \quad l = 1, \dots, \mu, \quad (2.9)$$

with the integers $\mathbf{m}_{l,i} \geq 0$. These monomials can be determined for any F by standard Gröbner basis algorithms implemented e.g. in SINGULAR. Two hypersurface singularities $\{F_1 = 0\}$ and $\{F_2 = 0\}$ are biholomorphically equivalent if and only if their Milnor rings are isomorphic [31]. The multiplicity of \mathbf{X} is related to the scaling weights according to

$$\mu = \prod_{i=1}^4 \left(\frac{1}{q_i} - 1 \right). \quad (2.10)$$

The spectrum of \mathbf{X} . The monomials (2.9) are assigned weights $Q_l \in \mathbb{Q}$ according to

$$Q_l = q(\mathbf{m}_l) \equiv \sum_{i=1}^4 q_i \mathbf{m}_{l,i}. \quad (2.11)$$

Let us (partially) order the monomials by their weights, with the lowest weight being $Q_1 = 0$ for $x^{\mathbf{m}_1} = 1$. Let us also introduce the shifted weights

$$\ell_l = Q_l + \sum_{i=1}^4 q_i - 1, \quad (2.12)$$

called the spectral numbers. These rational numbers are valued in the open interval $(0, 2)$. The so-called *spectrum* of the quasi-homogeneous singularity is then given by the ordered set:

$$\mathcal{S}_{\mathbf{X}} = \{\ell_l\}_{l=1}^{\mu}, \quad \ell_1 \leq \ell_2 \leq \dots \leq \ell_{\mu}. \quad (2.13)$$

The ordered monomials $x^{\mathbf{m}_l}$ and $x^{\mathbf{m}_{\mu-l+1}}$ of $\mathcal{M}(F)$ form pairs such that

$$\ell_l + \ell_{\mu-l+1} = 2 \quad \text{if} \quad \ell_l < 1. \quad (2.14)$$

There can also be unpaired monomials such that $\ell_l = 1$. The spectral numbers define a grading of the Milnor ring

$$\mathcal{M}(F) = \bigoplus_{\ell} \mathcal{M}(F)_{\ell}, \quad (2.15)$$

There is also a simple formula for the Poincaré series of the graded ring

$$P_{\mathcal{M}(F)}(q) \equiv \sum_{\ell} \dim \mathcal{M}(F)_{\ell} q^{\ell} = \frac{1}{q} \prod_{i=1}^4 \frac{q^{q_i} - q}{1 - q^{q_i}}, \quad (2.16)$$

which reduces to (2.10) in the limit $q \rightarrow 1$. The Poincaré series gives an efficient way of computing the spectrum for quasi-homogeneous hypersurface singularities, bypassing the need for finding an explicit basis of $\mathcal{M}(F)$. We denote by $\widehat{\mathbf{X}}$ a (generic) deformation of \mathbf{X} .

Another important number that characterizes the singularity \mathbf{X} is its modality, $m(\mathbf{X})$, which is defined as the number of deformations of \mathbf{X} which preserve its multiplicity. For quasi-homogeneous singularities, this is given by

$$m(\mathbf{X}) = \# \left\{ x^{\mathbf{m}} \mid \ell_l \geq \sum_{i=1}^4 q_i \right\}. \quad (2.17)$$

The corresponding deformation parameters t_l are the ‘moduli’ of the singularity.

The Coulomb branch spectrum of $\mathcal{T}_{\mathbf{X}}^{4d}$. In the Type IIB engineering, the family of deformations $\{\widehat{\mathbf{X}}_t\}$ is interpreted as the Seiberg-Witten geometry of the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ [2, 3]. The spectrum of \mathbf{X} is equivalent to the spectrum of (extended) Coulomb branch (CB) dimensions of $\mathcal{T}_{\mathbf{X}}^{4d}$, as follows. To each monomial $x^{\mathbf{m}}$, one assigns the CB dimension

$$\Delta_l \equiv \Delta(t_l) = \frac{1 - Q_l}{\sum_{i=1}^4 q_i - 1} = \frac{\sum_{i=1}^4 q_i - \ell_l}{\sum_{i=1}^4 q_i - 1}, \quad (2.18)$$

which are the scaling dimensions of the deformation parameters t_l divided by $\sum_i q_i - 1$. There are three classes of deformations, depending on Δ_l :

- $\Delta_l > 1$: these geometric deformations are interpreted as VEVs of CB operators of conformal dimensions Δ_l and $U(1)_r$ charge $r = 2\Delta_l$

$$t_l = \langle (\mathcal{E}_{r,(0,0)})_l \rangle. \quad (2.19)$$

We denote by \widehat{r} the number of such operators.

- $\Delta_l = 1$: these geometric deformations are interpreted as the mass deformations of $\mathcal{T}_{\mathbf{X}}^{4d}$

$$\delta S = t_l \int d^4x \int d^4\theta (\widehat{B}_1)_l, \quad l = 1, \dots, f. \quad (2.20)$$

They correspond to conserved currents, i.e. flavor symmetries. We may equivalently view the deformation parameters as VEVs for the scalars ϕ_l in the $U(1)$ background vectors multiplets spanning a maximal torus of the flavour symmetry group

$$t_l = \langle \phi_l \rangle \equiv m_l. \quad (2.21)$$

In particular, f is the rank of the flavor symmetry group of $\mathcal{T}_{\mathbf{X}}^{4d}$.

- $\Delta_l < 1$: these geometric deformations correspond to the \hat{r} monomials which are paired with the ones with $\Delta_{\mu-l} > 1$, such that $\Delta_l + \Delta_{\mu-l} = 2$. They are interpreted as chiral deformations of the SCFT by the CB operators themselves

$$\delta S = t_l \int d^4x \int d^4\theta (\mathcal{E}_{r,(0,0)})_{\mu-l}. \quad (2.22)$$

In particular, if $\Delta_l = 0$ ($\Delta_{\mu-l} = 2$), this is a marginal deformation of the SCFT — indeed, this corresponds to deforming the IHS equation by a monomial with $Q_l = 1$, thus preserving the scaling symmetry, which is necessary for conformal invariance. Note that we can also have deformations with $0 < \Delta_l < 1$, corresponding to relevant deformations of the SCFT ($1 < \Delta_{\mu-l} < 2$), or else deformations with $\Delta_l < 0$, corresponding to irrelevant deformations ($\Delta_{\mu-l} > 2$). The relevant deformations only exists if there are generators of the Milnor ring with scaling dimensions in the small range

$$\frac{\hat{c}}{2} < Q_{\mu-l} < 1. \quad (2.23)$$

Note also that the modality (2.17) is the total number of marginal and irrelevant operators of the 4d SCFT; thus, the ‘moduli’ of the singularity are precisely those deformations that do not trigger an RG flow away from the fixed point $\mathcal{T}_{\mathbf{X}}^{4d}$.

Let us denote by $\mathcal{B}_{\hat{r}}$ the Coulomb branch spanned by (2.19). In particular, \hat{r} is the rank of the SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$. We define the extended Coulomb branch $\mathcal{M}_C[\mathcal{T}_{\mathbf{X}}^{4d}]$ as the total space of the Coulomb branch fibered over the mass parameters (2.21), namely

$$\mathcal{B}_{\hat{r}} \rightarrow \mathcal{M}_C[\mathcal{T}_{\mathbf{X}}^{4d}] \rightarrow \{m\}, \quad (2.24)$$

The extended CB obviously has complex dimension $\hat{r} + f$. Note also the relation

$$\mu = 2\hat{r} + f. \quad (2.25)$$

The multiplicity μ gives the dimension of the charge lattice Γ of BPS particles on the CB, which arise as D3-branes wrapping the compact 3-cycles of the deformed hypersurface

$$\Gamma \cong H_3(\hat{\mathbf{X}}, \mathbb{Z}) \cong \mathbb{Z}^\mu. \quad (2.26)$$

Let us also denote by $\mathcal{M}_{3,3}$ the intersection form on the third homology

$$(\mathcal{M}_{3,3})_{\alpha\beta} = S_\alpha^3 \cdot S_\beta^3, \quad (2.27)$$

where the 3-cycles $\{S_\alpha^3\}_{\alpha=1}^\mu$ give some basis of $H_3(\hat{\mathbf{X}}, \mathbb{Z})$.

Central charges. From the Coulomb branch spectrum, we can compute the conformal anomaly coefficients a and c according to [3, 32]

$$a = \frac{R(A)}{4} + \frac{R(B)}{6} + \frac{5\hat{r}}{24}, \quad c = \frac{R(B)}{3} + \frac{\hat{r}}{6}. \quad (2.28)$$

Here, we defined

$$R(A) = \sum_{\Delta_l > 1} (\Delta_l - 1), \quad R(B) = \frac{\mu}{4(\sum_{i=1}^4 q_i - 1)} = \frac{1}{4}\mu\Delta_{\max}, \quad (2.29)$$

with $\Delta_{\max} \equiv \Delta_1$ the largest CB operator dimension. Here, we also used the fact that there is no hypermultiplets on the Coulomb branch of $\mathcal{T}_{\mathbf{X}}^{4d}$ because $H_2(\widehat{\mathbf{X}}, \mathbb{Z}) = 0$ [3]. From a and c , we may also compute the ‘effective number’ of hypermultiplets and of vector multiplets of the SCFT

$$n_h = -16a + 20c, \quad n_v = 8a - 4c. \quad (2.30)$$

For Lagrangian SCFTs, this gives the actual number of matter fields and gauge fields, respectively. For further discussions of the spectrum of 4d SCFTs in related contexts, see e.g. [33–36].

2.2 Crepant resolution and 5d Coulomb branch

For any canonical threefold singularity, there exists a crepant resolution $\pi : \tilde{\mathbf{X}} \rightarrow \mathbf{X}$, such that the resolved singularity $\tilde{\mathbf{X}}$ has at worst terminal singularities [37, 38]. In M-theory, the resolution $\tilde{\mathbf{X}}$ corresponds to probing the extended Coulomb branch of the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$ [39]. (See also e.g. [40–44] for some further discussions.)

Topology of the crepant resolution $\tilde{\mathbf{X}}$. We will call ‘a crepant resolution’ any (partial) resolution $\pi : \tilde{\mathbf{X}} \rightarrow \mathbf{X}$ that introduces r exceptional divisors with $a_k = 0$ in (2.1). This is expected to preserve 5d $\mathcal{N} = 1$ (or 4d $\mathcal{N} = 2$) supersymmetry. A smooth (or ‘complete’) crepant resolution, corresponding to $a_k = 0, \forall k$, in (2.1), is a crepant resolution that is smooth.

In general, there may be many distinct crepant resolutions of \mathbf{X} , but some of their key topological properties are invariant. The set of residual terminal singularities is such an invariant, and so are the Betti numbers

$$\begin{aligned} b_1(\tilde{\mathbf{X}}) &= 0, & b_2(\tilde{\mathbf{X}}) &= r + f, \\ b_3(\tilde{\mathbf{X}}) &= b_3, & b_4(\tilde{\mathbf{X}}) &= r. \end{aligned} \quad (2.31)$$

Here, r is the number of exceptional divisors in the crepant resolution, and $f = \rho(\mathbf{X})$ is the rank of the divisor class group of the singularity \mathbf{X} . We can also understand f as the number of 2-cycles of $\tilde{\mathbf{X}}$ that are dual to non-compact divisors. They correspond to U(1) flavor symmetries of the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$. Note also that the crepant resolution $\tilde{\mathbf{X}}$ generally has three-cycles, as we will explain momentarily.

There is a remarkable relation between f as defined here and the integer f defined in (2.20) from the spectrum of the singularity. It is a non-trivial mathematical fact that the two quantities agree [45]

$$f = \rho(\mathbf{X}) = \#\{\ell_l \mid \ell_l = 1\}. \quad (2.32)$$

Physically, this can be interpreted in terms of generalised geometric transitions, wherein f 2-cycles in the resolution $\tilde{\mathbf{X}}$ become f 3-cycles in the deformation $\hat{\mathbf{X}}$ [1].³ A third way to compute f is from the five-dimensional link of the singularity, as we will review in section 3.

Newton polytope and weightings. For any quasi-homogeneous IHS, the Betti numbers (2.31) can be computed in the following way [45]. Let us define the Newton polytope $\mathcal{N}(F)$ associated to $\{F(x) = 0\}$ as the polytope in the lattice of monomials $M \cong \mathbb{Z}^4$

$$\mathcal{N}(F) = \{(\mathbf{m}_{n,1}, \mathbf{m}_{n,2}, \mathbf{m}_{n,3}, \mathbf{m}_{n,4}) \in M = \mathbb{Z}^4 \mid \mathbf{m}_n \in F\}. \quad (2.33)$$

Here, $\mathbf{m}_n \in F$ means that the monomial $x^{\mathbf{m}_n}$ enters the polynomial $F(x)$ with non-zero coefficient. Let us also define the weightings α as the maps $\alpha : M \rightarrow \mathbb{Z}$, which can be viewed as vectors $(\alpha_1, \alpha_2, \alpha_3, \alpha_4)$ in the dual lattice $N \cong \text{Hom}(M, \mathbb{Z})$. The evaluation map on a monomial \mathbf{m} reads

$$\alpha(\mathbf{m}) = \sum_{i=1}^4 \alpha_i \mathbf{m}_i, \quad (2.34)$$

and the evaluation map on a polynomial F is defined as

$$\alpha(F) = \min_{\mathbf{m} \in F} \alpha(\mathbf{m}). \quad (2.35)$$

We then define the sets

$$W_+(F) = \left\{ \alpha \in N^+ \mid \forall \mathbf{m} \in \mathcal{N}(F), \alpha(\mathbf{m} - (1, 1, 1, 1)) \geq -1 \right\}. \quad (2.36)$$

and

$$W(F) = W_+(F) \cup \{(1, 0, 0, 0), (0, 1, 0, 0), (0, 0, 1, 0), (0, 0, 0, 1)\}. \quad (2.37)$$

N^+ is the open cone of N satisfying $\alpha_i > 0$ ($i = 1, \dots, 4$). We can then define a non-compact toric fourfold T_Σ with toric fan Σ . The elements of $W(F)$ are in one-to-one correspondence with the 1d cone of Σ , and one must choose a set of 2d, 3d and 4d cones of Σ , such that Σ is maximally triangulated. The resolved threefold $\tilde{\mathbf{X}}$ is then exactly the anti-canonical hypersurface $-K(T_\Sigma)$ of T_Σ , and each triangulation corresponds to a different (toric) resolution of \mathbf{X} . Finally, each element of $W_+(F)$ corresponds to a compact exceptional divisor $S_i \subset \tilde{\mathbf{X}}$, which may be reducible.

Computation of r and b_3 . The topological numbers r and b_3 are independent of the choice of resolution, and they can be computed as follows [45, 46] (analogously to the Batyrev formula [47]). Consider the dual polytope Γ of $W(F)$, defined as

$$\Gamma = \left\{ \mathbf{m} \in M^+ \mid \forall \alpha \in W(F), \alpha(\mathbf{m} - (1, 1, 1, 1)) \geq -1 \right\}. \quad (2.38)$$

The dual face $\Gamma_\alpha \subset \Gamma$ of $\alpha \in W(F)$, defined as

$$\Gamma_\alpha = \{ \mathbf{m} \in \Gamma \mid \alpha(\mathbf{m} - (1, 1, 1, 1)) = -1 \} \quad (2.39)$$

³The proper mathematical understanding of this relation involves discussing the mixed Hodge structure of the singularity, which we will review in more detail in [7].

can have different dimensions, denoted by $\dim(\Gamma_\alpha)$. The explicit formula for r is then

$$r = \sum_{\alpha \in W_+(F)} \begin{cases} 0 & \text{if } \dim(\Gamma_\alpha) = 0, \\ l'(\Gamma_\alpha) + 1 & \text{if } \dim(\Gamma_\alpha) = 1, \\ 1 & \text{if } \dim(\Gamma_\alpha) > 1, \end{cases} \quad (2.40)$$

where $l'(\Gamma_\alpha)$ is number of interior lattice points in Γ_α . When $\dim(\Gamma_\alpha) = 1$ and $l'(\Gamma_\alpha) > 0$, the compact divisor corresponding to α has $l'(\Gamma_\alpha) + 1$ irreducible components.

To compute b_3 , the number of 3-cycles in the crepant resolution, we first note that each compact divisor associated to the weighting $\alpha \in W_+(F)$ has the structure of \mathbb{P}^1 fibration over a complex curve with genus $g_C(\alpha)$. The numbers $g_C(\alpha)$ are computed as

$$g_C(\alpha) = \begin{cases} l'(\Gamma_\alpha) & \text{if } \dim(\Gamma_\alpha) = 2, \\ 0 & \text{if } \dim(\Gamma_\alpha) \neq 2. \end{cases} \quad (2.41)$$

The formula for b_3 is then

$$b_3 = 2 \sum_{\alpha \in W_+(F)} g_C(\alpha). \quad (2.42)$$

In other words, there is a one-to-one correspondence between the 3-cycles and the 1-cycles at the base of the ruled exceptional divisors. Note that the resolved threefold $\tilde{\mathbf{X}}$ is itself simply-connected.

Resolution sequence. Finally, let us discuss the resolution sequence itself. A toric resolution sequence is defined as follows. For any IHS, we start with the non-compact ambient space \mathbb{C}^4 with coordinates (x_1, x_2, x_3, x_4) , with the set of weightings $\{(1, 0, 0, 0), (0, 1, 0, 0), (0, 0, 1, 0), (0, 0, 0, 1)\}$. Each step in the resolution sequence is a weighted blow-up of the ambient space at the locus $y_1 = \dots = y_k = 0$, where $y_i = 0$ is the toric divisor S_i^Σ in the ambient space corresponding to the weighting α_i . It takes one of the following forms [45, 48, 49]:

1. $(y_1^{(a_1)}, y_2^{(a_2)}, y_3^{(a_3)}, y_4^{(a_4)}; \delta)$: weighted blow-up of the locus $y_1 = y_2 = y_3 = y_4 = 0$ with weight (a_1, a_2, a_3, a_4) . Here, the only possible weights are $(3, 2, 1, 1)$, $(2, 1, 1, 1)$ and $(1, 1, 1, 1)$. The exceptional divisor $\{\delta = 0\}$ in the ambient space is associated to the weighting

$$\alpha(\delta) = \sum_{i=1}^4 a_i \alpha(y_i). \quad (2.43)$$

In the hypersurface equation, one replaces

$$y_1 \rightarrow y_1 \delta^{a_1}, \quad y_2 \rightarrow y_2 \delta^{a_2}, \quad y_3 \rightarrow y_3 \delta^{a_3}, \quad y_4 \rightarrow y_4 \delta^{a_4}, \quad (2.44)$$

and then divide the equation by $\delta^{a_1+a_2+a_3+a_4-1}$. There is a new SR ideal generator $y_1 y_2 y_3 y_4$ after the resolution.

2. $(y_1, y_2, y_3; \delta)$: blow up the locus $y_1 = y_2 = y_3 = 0$. The exceptional divisor $\delta = 0$ in the ambient space is associated to the weighting

$$\alpha(\delta) = \alpha(y_1) + \alpha(y_2) + \alpha(y_3). \quad (2.45)$$

In the hypersurface equation, one replaces

$$y_1 \rightarrow y_1 \delta, \quad y_2 \rightarrow y_2 \delta, \quad y_3 \rightarrow y_3 \delta, \quad (2.46)$$

and then divide the equation by δ^2 . There is a new SR ideal generator $y_1 y_2 y_3$ after the resolution.

3. $(y_1, y_2; \delta)$: blow up the locus $y_1 = y_2 = 0$. The exceptional divisor $\delta = 0$ in the ambient space is associated to the weighting

$$\alpha(\delta) = \alpha(y_1) + \alpha(y_2). \quad (2.47)$$

In the hypersurface equation, one replaces

$$y_1 \rightarrow y_1 \delta, \quad y_2 \rightarrow y_2 \delta, \quad (2.48)$$

and then divide the equation by δ . There is a new SR ideal generator $y_1 y_2$ after the resolution.

After the resolution sequence, one can also compute the topological quantities associated to each exceptional divisor. In particular, we can compute the intersection numbers of the exceptional divisors S_i inside $\tilde{\mathbf{X}}$, according to

$$S_i \cdot S_j \cdot S_k|_{\tilde{\mathbf{X}}} = (-K_\Sigma) \cdot S_i^\Sigma \cdot S_j^\Sigma \cdot S_k^\Sigma. \quad (2.49)$$

This determines the prepotential on the 5d CB of $\mathcal{T}_{\mathbf{X}}^{5d}$ [39].

3 Link topology and higher-form symmetries

The five-manifold L_5 at the boundary of a non-compact Calabi-Yau three-fold encodes global symmetries of the field theories, both in 4d and 5d. In particular, the torsion part of $H_*(L_5, \mathbb{Z})$ encodes the higher-form symmetries [50] of both the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$ and the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ [1, 51–54]. For isolated hypersurface singularities, L_5 is simply connected, and the 5d field theory $\mathcal{T}_{\mathbf{X}}^{5d}$ does not have any 1-form symmetries. On the other hand, we often have non-zero torsion in $H_2(L_5, \mathbb{Z})$, which can be computed from the weights q_i of the quasi-homogeneous polynomial $F(x)$ [55, 56]. In M-theory, this translates into 3-form symmetries of $\mathcal{T}_{\mathbf{X}}^{5d}$ that act on M5-branes wrapping relative 3-cycles (and/or its dual 0-form symmetry acting on wrapped M2-branes). In Type IIB string theory, this torsion group is related to the 1-form symmetry of the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$, with the charged line operators arising from wrapped D3-branes.

In the rest of this section, we first review an explicit formula for the link homology, and we then briefly discuss its 5d and 4d interpretation.

3.1 Topology of the link

The link of the singularity \mathbf{X} is the 5-dimensional compact manifold $L_5(\mathbf{X})$ obtained by intersecting the hypersurface with a seven-sphere S_ε^7 centered at the origin of \mathbb{C}^4

$$L_5(\mathbf{X}) = \left\{ x \in \mathbb{C}^4 \left| F(x) = 0, \quad \sum_{i=1}^4 |x_i|^2 = \varepsilon \right. \right\}, \quad (3.1)$$

for a small radius $\varepsilon > 0$. For quasi-homogeneous singularities, we may pick any value of ε due to the scaling symmetry. Therefore, by sending $\varepsilon \rightarrow \infty$, we see that the link is also the *boundary manifold* of \mathbf{X} and of any local desingularization of the singularity; in particular, for any smooth deformation $\hat{\mathbf{X}}$ of \mathbf{X} , we have $\partial \hat{\mathbf{X}} \cong L_5(\mathbf{X})$. The five-manifold L_5 is connected and simply-connected, and its homology groups are

$$\begin{aligned} H_0(L_5, \mathbb{Z}) &\cong H_5(L_5, \mathbb{Z}) \cong \mathbb{Z}, & H_2(L_5, \mathbb{Z}) &\cong \mathbb{Z}^f \oplus \mathfrak{t}_2 \\ H_1(L_5, \mathbb{Z}) &\cong H_4(L_5, \mathbb{Z}) \cong 0, & H_3(L_5, \mathbb{Z}) &\cong \mathbb{Z}^f, \end{aligned} \quad (3.2)$$

where

$$\mathfrak{t}_2 = \text{Tor } H_2(L_5, \mathbb{Z}) \quad (3.3)$$

denotes the torsion part of the second homology. This leads to our third definition of the ‘flavor rank’ f , as the Betti number of the link

$$f = b_2(L_5) = b_3(L_5). \quad (3.4)$$

This gives us an intuitive way of understanding the ‘generalised geometric transition’ mentioned below (2.32), generalising the conifold transition [57]: as we resolve or deform the singularity \mathbf{X} , we have f exceptional 2-cycles that survive ‘at the tip’ inside $\tilde{\mathbf{X}}$, or f 3-cycles that survive inside $\hat{\mathbf{X}}$, respectively.

The homology of $L_5(\mathbf{X})$: explicit formulas. The non-trivial quantities f and \mathfrak{t}_2 in (3.2) can be computed from the scaling weights q_i in (2.3) as follows. Let us define the quantities

$$v_i \equiv \frac{1}{q_i}, \quad a_i \equiv \text{Numerator}(v_i). \quad (3.5)$$

Let us consider the set $I = \{1, 2, 3, 4\}$, and let I_s denote any particular subset $I_s \subset I$ of s elements. Then, we have the explicit formula [55, 56]

$$f = b_2(L_5) = \sum_{s=0}^4 \sum_{I_s \subset I} (-1)^{4-s} \frac{\prod_{j \in I_s} v_j}{\text{LCM}_{j \in I_s}(a_j)}. \quad (3.6)$$

The sum includes all subsets I_s of I , including the trivial subset $I_0 = \{\}$, which contributes +1. More generally, we have the following formula for the torsion part, \mathfrak{t}_2 [56]. For every $I_s \subset I$, we define

$$\kappa(I_s) = \sum_{t=0}^s \sum_{I_t \subset I_s} (-1)^{s-t} \frac{\prod_{i \in I_t} v_i}{\text{LCM}_{j \in I_t}(a_j)}, \quad (3.7)$$

and

$$k(I_s) = \begin{cases} \kappa(I_s) & \text{if } 4 - s \in 2\mathbb{Z} + 1, \\ 0 & \text{if } 4 - s \in 2\mathbb{Z}. \end{cases} \quad (3.8)$$

We then define a set of integers, $C(I_s)$, inductively, starting with

$$C(I_0) = C(\{\}) = \gcd_{i \in I} (a_i) = \gcd(a_1, \dots, a_4), \quad (3.9)$$

and

$$C(I_s) = \frac{\gcd_{i \in I - I_s} (a_i)}{\prod_{I_t \subsetneq I_s} C(I_t)}. \quad (3.10)$$

Then, for the integers $j > 1$, we define

$$d_j = \prod_{I_s | k(I_s) \geq j} C(I_s), \quad j = 1, \dots, J, \quad J \equiv \max_{I_s \subset I} k(I_s). \quad (3.11)$$

We then have:

$$\mathfrak{t}_2 \cong \mathbb{Z}_{d_1} \oplus \dots \oplus \mathbb{Z}_{d_J}, \quad f = \kappa(I). \quad (3.12)$$

The torsion group \mathfrak{t}_2 from the deformation $\widehat{\mathbf{X}}$. Using a long-exact sequence for the (relative) homology groups of the deformed threefold $\widehat{\mathbf{X}}$, as reviewed in the next subsection, one can show that the torsion group \mathfrak{t}_2 is also isomorphic to the cokernel of the map $\mathcal{M}_{3,3} : \mathbb{Z}^\mu \rightarrow \mathbb{Z}^\mu$ given by the intersection form on $H_3(\widehat{\mathbf{X}}, \mathbb{Z})$, namely

$$\mathfrak{t}_2 \cong \mathbb{Z}^\mu / \mathcal{M}_{3,3} \mathbb{Z}^\mu. \quad (3.13)$$

We computed this cokernel in many examples, allowing us to check the general formula (3.12) for \mathfrak{t}_2 , which appeared as a conjecture in [56], on a case-by-case basis.⁴

3.2 Defects and higher-form symmetries in 5d from M-theory

Consider the canonical singularity \mathbf{X} , any desingularisation \mathbf{X}_6 of \mathbf{X} (which will be either the deformation $\widehat{\mathbf{X}}$ or the crepant resolution $\widetilde{\mathbf{X}}$), and the boundary five-manifold $L_5 \cong \partial \mathbf{X}_6$. M-theory on \mathbf{X}_6 provides us with a natural set of defect operators of the 5d low-energy theory, from membranes wrapping non-compact k -cycles. The wrapped M2- and M5-branes give rise to defect operators of dimension $q = k - 3$ and $q = 6 - k$, respectively, which can be charged under a non-trivial q -form symmetry. The sets of such charged ‘electric’ and ‘magnetic’ defect operators of dimension q are denoted by

$$\Gamma_{\text{M2}}^{(q)} = \mathfrak{h}_{(k=3-q)}, \quad \Gamma_{\text{M5}}^{(q)} = \mathfrak{h}_{(k=6-q)}, \quad (3.14)$$

respectively. They are computed as

$$\mathfrak{h}_{(k)} \equiv \text{Tor}(H_k(\mathbf{X}_6, L_5, \mathbb{Z}) / \text{im}(f_k)), \quad (3.15)$$

⁴A proof of (3.12) is given in [55] for Type I singularities and in [58] for Type XIX singularities, in the notation of section 4.1 below. As far as we are aware, the general case remains a conjecture.

where f_k is the map of $f_k : H_k(\mathbf{X}_6, \mathbb{Z}) \rightarrow H_k(\mathbf{X}_6, L_5, \mathbb{Z})$ that sends compact k -cycles to relative k -cycles. It sits inside a long exact sequence for the relative homology (with integral coefficients)

$$\cdots \rightarrow H_k(L_5) \xrightarrow{h_k} H_k(\mathbf{X}_6) \xrightarrow{f_k} H_k(\mathbf{X}_6, L_5) \xrightarrow{g_k} H_{k-1}(L_5) \xrightarrow{h_{k-1}} H_{k-1}(\mathbf{X}_6) \rightarrow \cdots \quad (3.16)$$

The objects counted by $\mathfrak{h}_{(k)}$ are the defects whose charge cannot be screened by the dynamical particles that arise from membranes wrapping compact k -cycles of \mathbf{X}_6 . For $q \leq 1$, for definiteness, let us define the ‘defect groups’ of the 5d theory, in its massive phase determined by \mathbf{X}_6

$$\mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(q)} \cong \Gamma_{M2}^{(q)} \oplus \Gamma_{M5}^{(3-q)}. \quad (3.17)$$

The defect operators in $\Gamma_{M2}^{(q)}$ and $\Gamma_{M5}^{(3-q)}$ can have a non-trivial Dirac pairing given by the intersection of the relative $(3-k)$ - and $(3+k)$ -cycles inside \mathbf{X}_6 . The higher-form symmetry of the 5d theory is determined after one chooses a polarisation of the defect group (3.17),

$$\Lambda_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(q)} \subset \mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(q)}, \quad (3.18)$$

by fixing a maximal set of mutually commuting defects [51, 53, 54, 59–63]. The higher-form symmetry group itself is given by the Pontryagin dual group

$$\widehat{\Lambda}_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(q)} = \text{Hom}(\Lambda_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(q)}, \text{U}(1)). \quad (3.19)$$

Note that it is generally a mixture of q - and $(3-q)$ -form symmetries. We may also choose, for simplicity, the ‘purely electric’ or the ‘purely magnetic’ polarisation, consisting in picking one of the two summands in (3.17).

Given a general local Calabi-Yau threefold in M-theory, there are two defects groups of interest

$$\mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(0)} \cong \mathfrak{h}_{(3)} \oplus \mathfrak{h}_{(3)}, \quad \mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{5d}}^{(1)} \cong \mathfrak{h}_{(2)} \oplus \mathfrak{h}_{(4)}, \quad (3.20)$$

corresponding to the ‘0-form symmetry’ and the ‘1-form symmetry’, respectively, allowing for a slight abuse of language. Using the exact sequence (3.16), we can write the groups $\mathfrak{h}_{(k)}$ in terms of the homology of the boundary:

$$\mathfrak{h}_{(k)} = \text{Tor}(\text{im}(g_k)) = \text{Tor}(\text{ker}(h_{k-1})) \subset \text{Tor } H_{k-1}(L_5). \quad (3.21)$$

The charged defect operators wrap relative k -cycles Σ_k with non-trivial boundary on L_5 , $[\partial \Sigma_k] \in \text{Tor } H_{k-1}(L_5)$. For L_5 any connected, orientable closed manifold, there are two independent torsion groups in its homology

$$\mathfrak{t}_1 \equiv \text{Tor } H_2(L_5, \mathbb{Z}) \cong \text{Tor } H_3(L_5, \mathbb{Z}), \quad \mathfrak{t}_2 \equiv \text{Tor } H_2(L_5, \mathbb{Z}). \quad (3.22)$$

Let us now see how these torsion groups are related to higher-form symmetries in 5d.

1-form symmetry and Coulomb-branch defects. For the sake of generality, let us first discuss the possibility of a non-trivial 1-form symmetry in 5d [51, 53, 63]. We will assume that the smooth threefold \mathbf{X}_6 is simply-connected and that it has no torsion in homology, which is the case for all examples we have studied (such as toric threefolds or hypersurface singularities). It directly follows from (3.16) and (3.21) that the ‘electric’ part of the defect group is isomorphic to the torsion part of the first homology group of the boundary.

$$\mathfrak{h}_{(2)} \cong \mathfrak{t}_1. \quad (3.23)$$

Alternatively, we can use the definition (3.15) of $\mathfrak{h}_{(2)}$ as the cokernel of the map f_2 , which gives

$$\mathfrak{h}_{(2)} \cong \text{Tor}(\mathbb{Z}^{b_4}/\mathcal{M}_{4,2}\mathbb{Z}^{b_2}), \quad (3.24)$$

where b_j denote the Betti numbers of \mathbf{X}_6 . Similarly, we have

$$\mathfrak{h}_{(4)} \cong \text{Tor}(\mathbb{Z}^{b_2}/\mathcal{M}_{2,4}\mathbb{Z}^{b_4}) \cong \mathfrak{h}_{(2)}, \quad (3.25)$$

The 1-form (and/or 2-form) symmetry is then determined by the intersection pairing between the compact 2-cycles and 4-cycles of \mathbf{X}_6 . The charged defect operators naturally live on the Coulomb branch of the 5d SCFT, in which case $\mathbf{X}_6 = \tilde{\mathbf{X}}$, $b_2 = r + f$ and $b_4 = f$. For the 5d SCFTs $\mathcal{T}_{\tilde{\mathbf{X}}}^{5d}$ of interest in this paper, which arise at hypersurface singularities, the link manifold L_5 is simply connected, and therefore we find that $\mathfrak{h}_{(2)} = \mathfrak{h}_{(4)} = 0$.

0- and 3-form symmetry and Higgs-branch defects. The 0-form/3-form symmetries of $\mathcal{T}_{\tilde{\mathbf{X}}}^{5d}$ act on defect operators that are more naturally defined on the Higgs branch of $\mathcal{T}_{\tilde{\mathbf{X}}}^{5d}$. Let us then take $\mathbf{X}_6 = \hat{\mathbf{X}}$ a deformed singularity, and assume that $H_2(\hat{\mathbf{X}}) = 0$ for simplicity, which is the case, in particular, for the generic deformation of an hypersurface singularity. We then find that

$$\mathfrak{h}_{(3)} \cong \mathfrak{t}_2 \cong \text{Tor}(\mathbb{Z}^{b_3}/\mathcal{M}_{3,3}\mathbb{Z}^{b_3}), \quad (3.26)$$

with $b_3 = \mu = 2\hat{r} + f$ for an IHS. It would be interesting to have a more detailed understanding of these charged HB defect operators from the 5d field-theory point of view. This is left as a challenge for future work.

3.3 Defects and higher-form symmetries in 4d from type IIB

In the Type IIB compactification on \mathbf{X}_6 , the natural set of charged defect operators of the 4d field theory arise from wrapped D-branes. We may define the defect group

$$\mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(0)} \cong \Gamma_{D1}^{(0)} \oplus \Gamma_{D5}^{(2)} \oplus \Gamma_{D3}^{(2)} \oplus \Gamma_{D3}^{(0)} \cong \mathfrak{h}_{(2)} \oplus \mathfrak{h}_{(4)} \oplus \mathfrak{h}_{(2)} \oplus \mathfrak{h}_{(4)}, \quad (3.27)$$

corresponding to D1-, D5- and D3-branes wrapping 2- and 4-cycles, as well as the defect group

$$\mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)} \cong \Gamma_{D3}^{(1)} \cong \mathfrak{h}_{(3)} \cong \mathfrak{t}_2, \quad (3.28)$$

from D3-brane wrapping relative 3-cycles inside \mathbf{X}_6 . Since our focus is on hypersurface singularities, we will not consider further the interesting possibilities suggested by (3.27). In

other words, our 4d theories $\mathcal{T}_{\mathbf{X}}^{4d}$ will have no defects charged under non-trivial 0-form or 2-form symmetries. On the other hand, the defect group (3.28) is generally non-trivial, and it encodes all the (‘electric’ and ‘magnetic’) line operators of $\mathcal{T}_{\mathbf{X}}^{4d}$ [1, 53] — see also [10, 62–64] for related discussions. It is interesting to note that, for 4d SCFTs $\mathcal{T}_{\mathbf{X}}^{4d}$ engineered at IHS, non-trivial 1-form symmetries can only arise if the SCFT is non-isolated, as shown in [10].⁵

The actual one-form symmetry $\widehat{\Lambda}_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)}$ of $\mathcal{T}_{\mathbf{X}}^{4d}$ is obtained upon choosing a polarisation

$$\Lambda_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)} \subset \mathbf{D}_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)}, \quad \widehat{\Lambda}_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)} = \text{Hom}(\Lambda_{\mathcal{T}_{\mathbf{X}}^{4d}}^{(1)}, \text{U}(1)). \quad (3.29)$$

For our hypersurface singularities, the torsion in the second homology of the link always takes the form:

$$\mathfrak{t}_2 \cong \mathfrak{f} \oplus \mathfrak{f}, \quad (3.30)$$

so that there exists a ‘purely electric’ choice of polarisation for which the one-form symmetry group is isomorphic to \mathfrak{f} [1].

4 Classification of canonical IHS and Mathematica code

In this section, we review the complete classification of canonical isolated singularities, which is equivalent to the Kreuzer-Skarke classification of 2d $\mathcal{N} = (2, 2)$ Landau-Ginzburg (LG) superpotentials [65]. This 2d perspective clarifies some aspects of the discussion in [3, 8].

We then present an ancillary `Mathematica` code [66] (available in the supplementary material of this paper) which implements the computations explained in the previous sections for any such singularity. The code computes the spectrum of \mathbf{X} , the basic topological data of its crepant resolution $\widetilde{\mathbf{X}}$, as well as the homology of the link L_5 . This directly gives all the ‘basic data’ associated to the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$ and the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$. The code also implements the crepant resolution procedure explained in section 2.2.

4.1 Classification of canonical isolated hypersurface singularities in \mathbb{C}^4

Consider a quasi-homogeneous IHS in \mathbb{C}^4 , defined by the polynomial

$$F(x_1, x_2, x_3, x_4) \in \mathbb{C}[x_1, x_2, x_3, x_4], \quad (4.1)$$

with the scaling weights q_i as in (2.3). Such quasi-homogeneous polynomials $F(x)$ were classified by Kreuzer and Skarke (KS) many years ago [65]. A classification of isolated hypersurface threefold singularities was also given by Yau and Yu (YY) [8], which was used by Xie and Yau in their studies of geometrically-engineered 4d $\mathcal{N} = 2$ SCFTs [3]. The YY classification is slightly incomplete, however [9].

⁵This claim was checked systematically in [10], and it is easily checked with the ‘experimental’ methods of this paper. In particular, in our finite database of models discussed in section 5, 56% of the models are isolated (no marginal deformations) and all those isolated models have $\mathfrak{t}_2 = 0$. Only 15% of all models (amounting to 35% of the models with marginal deformations) have a non-trivial one-form symmetry.

It is very useful to think of the polynomial (4.1) as defining a 2d Landau-Ginzburg (LG) model with four chiral superfields x_i and a superpotential $W = F(x)$ [2]. Any quasi-homogeneous IHS then defines a 2d $\mathcal{N} = (2, 2)$ SCFT as well as the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$, with the two perspectives related by a 4d/2d correspondence [67], provided that the reduced central charge of the 2d theory satisfies [2, 68, 69]:

$$\hat{c} \equiv \sum_{i=1}^4 (1 - 2q_i) < 2. \quad (4.2)$$

In fact, the IHS is canonical if and only if (4.2) holds [28, 70]. Such LG models were recently classified by Davenport and Melnikov [9].⁶ Let us review their result, while keeping the YY notation for ease of comparison. One finds that the YY classification needs to be refined, by introducing ‘subtypes’ for some of the 19 ‘types’. Since this classification is essentially equivalent to the KS classification of quasi-homogeneous polynomials, we will call it the KS-YY classification.

The full KS-YY classification. The YY classification distinguishes between 19 types of polynomials $F(x)$, denoted by the roman numerals I to XIX, and we denote the subtypes necessary for the full classification by subscripts.⁷ Then, for a given type and subtype, the positive integers (a, b, c, d) fully determine the singularity (up to a choice of solution to some linear Diophantine equation, in some cases, which we discuss momentarily). For instance, the singularity IX₂(2, 3, 4, 2) is given by $F = x_1^2 + x_2^3 x_4 + x_3^4 x_4 + x_2 x_4^2 + x_1 x_2 x_3^2$. The complete classification is given in tables 1, 2 and 3.

It is very important to note that this classification includes many redundancies: several singularities of distinct types can be physically equivalent — two quasi-homogeneous singularities are physically equivalent if and only if they have the same scaling weights, up to ordering. From the 2d LG model perspective, these scaling dimensions (q_i) are the fundamental quantities, and any two superpotentials $W = F_1(x)$ and $W = F_2(x)$ that realise the same scaling dimensions must flow to the same 2d SCFT. From the 4d SCFT point of view, two singularities with the same scaling weights obviously give rise to the same 4d CB spectrum, as per (2.16).

In order to explain the KS-YY classification, let us introduce some useful terminology and notation [9, 65]. First of all, an (LG) field x_i that appears in $F(x)$ is called a ‘root’ if it appears as $x_i^{m_i}$, and it is called a ‘pointer’ if it appears as $x_i^{m_i} x_j$, for some positive integer m_i . One can show that every x_i in $F(x)$ must be either a root or a pointer. We say that a pointer x_i ‘points at’ the field x_j . One can then introduce a convenient graphical notation for the polynomial. Each x_i , root or pointer, is denoted by a dot accompanied with the integer m_i , and for each pointer we draw an arrow from x_i to x_j . For instance, the polynomial in two variables $F = x_1^{m_1} + x_1 x_2^{m_2}$ is represented by

$$\begin{array}{ccc} m_1 & & m_2 \\ \circ & \longleftarrow & \circ \end{array} \quad (4.3)$$

⁶Note that [9] also considered LG models with five fields such that $\hat{c} < 2$. We only consider four fields.

⁷We also use the notation {type, subtype} with type $\in \{1, 2, \dots, 19\}$, as in the `Mathematica` notebook.

Type		$F(x_1, x_2, x_3, x_4)$	Skeleton and links				Skeleton
I	$\{1, 1\}$	$x_1^a + x_2^b + x_3^c + x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$(S_{1,1})^{\oplus 4}$
II	$\{2, 1\}$	$x_1^a + x_2^b + x_3^c + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$(S_{1,1})^{\oplus 2} \oplus S_{2,1}$
III	$\{3, 1\}$	$x_1^a + x_2^b + x_3^c x_4 + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$		$(S_{1,1})^{\oplus 2} \oplus S_{2,2}$
IV	$\{4, 1\}$	$x_1^a + x_1 x_2^b + x_3^c + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$S_{2,1} \oplus S_{2,1}$
V	$\{5, 1\}$	$x_1^a x_2 + x_1 x_2^b + x_3^c + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} b \\ \circ \end{smallmatrix}$		$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$S_{2,2} \oplus S_{2,1}$
VI	$\{6, 1\}$	$x_1^a x_2 + x_1 x_2^b + x_3^c x_4 + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} b \\ \circ \end{smallmatrix}$		$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$		$S_{2,2} \oplus S_{2,2}$
VII	$\{7, 1\}$	$x_1^a + x_2^b + x_2 x_3^c + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$S_{1,1} \oplus S_{3,1}$
VIII ₁	$\{8, 1\}$	$x_1^a + x_2^b + x_2 x_3^c + x_2 x_4^d + x_3^p x_4^q$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$			$S_{1,1} \oplus S_{3,2}$
VIII ₂	$\{8, 2\}$	$x_1^a + x_2^b + x_2 x_3^c + x_2 x_4^d + x_1 x_3^p x_4^q$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$			$S_{1,1} \oplus S_{3,2}$
IX ₁	$\{9, 1\}$	$x_1^a + x_2^b x_4 + x_3^c x_4 + x_2 x_4^d + x_2^p x_3^q$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$			$S_{1,1} \oplus S_{3,4}$
IX ₂	$\{9, 2\}$	$x_1^a + x_2^b x_4 + x_3^c x_4 + x_2 x_4^d + x_1 x_2^p x_3^q$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$			$S_{1,1} \oplus S_{3,4}$
X	$\{10, 1\}$	$x_1^a + x_2^b x_3 + x_3^c x_4 + x_2 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix}$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$			$S_{1,1} \oplus S_{3,3}$
XI	$\{11, 1\}$	$x_1^a + x_1 x_2^b + x_2 x_3^c + x_3 x_4^d$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} b \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} c \\ \circ \end{smallmatrix} \longleftarrow$	$\begin{smallmatrix} d \\ \circ \end{smallmatrix}$	$S_{4,1}$
XII ₁	$\{12, 1\}$	$x_1^a + x_1 x_2^b + x_1 x_3^c + x_2 x_4^d + x_2^p x_3^q$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$				$S_{4,2}$
XII ₂	$\{12, 2\}$	$x_1^a + x_1 x_2^b + x_1 x_3^c + x_2 x_4^d + x_2^p x_3^q x_4$	$\begin{smallmatrix} a \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} b \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} c \\ \circ \end{smallmatrix} \begin{smallmatrix} \curvearrowright \\ \curvearrowleft \end{smallmatrix} \begin{smallmatrix} d \\ \circ \end{smallmatrix}$				$S_{4,2}$

Table 1. The KS-YY classification of isolated hypersurface singularities in \mathbb{C}^4 : types I to XIII. Note that types I–VI exhaust the possibilities using skeleton with up to $n = 2$ fields, while the cases I–X exhaust the possibilities using skeletons with up to $n = 3$ fields.

Type		$F(x_1, x_2, x_3, x_4)$	Skeleton and links	Skeleton
XIII ₁	{13, 1}	$x_1^a + x_1x_2^b + x_2x_3^c + x_2x_4^d + x_3^px_4^q$		$S_{4,3}$
XIII ₂	{13, 2}	$x_1^a + x_1x_2^b + x_2x_3^c + x_2x_4^d + x_1x_3^px_4^q$		$S_{4,3}$
XIV ₁	{14, 1}	$x_1^a + x_1x_2^b + x_1x_3^c + x_1x_4^d + x_2^px_3^qx_4^r + x_2^ux_4^v$		$S_{4,4}$
XIV ₂	{14, 2}	$x_1^a + x_1x_2^b + x_1x_3^c + x_1x_4^d + x_2^px_3^qx_4^r + x_2^ux_4^v$		$S_{4,4}$
XIV ₃	{14, 3}	$x_1^a + x_1x_2^b + x_1x_3^c + x_1x_4^d + x_2^px_3^qx_4^r + x_2^ux_4^v$		$S_{4,4}$
XIV ₄	{14, 4}	$x_1^a + x_1x_2^b + x_1x_3^c + x_1x_4^d + x_2^px_3^qx_4^r + x_2^ux_3x_4^v$		$S_{4,4}$
XV ₁	{15, 1}	$x_1^ax_2 + x_1x_2^b + x_1x_3^c + x_3x_4^d + x_2^px_3^q$		$S_{4,6}$
XV ₂	{15, 2}	$x_1^ax_2 + x_1x_2^b + x_1x_3^c + x_3x_4^d + x_2^px_3^qx_4$		$S_{4,6}$

Table 2. The KS-YY classification of isolated hypersurface singularities in \mathbb{C}^4 : types XIII to XVI.

With four fields x_1, \dots, x_4 , there are 19 distinct possible ‘skeletons’, corresponding to the 19 classes in the YY classification. As we see from the tables, some of the skeletons are reducible. The basic building blocks are the skeletons denoted by $S_{n,\alpha}$ in [9], where n denotes the number of fields connected by arrows. For $n = 1, 2, 3, 4$, there are 1, 2, 4 and 9 distinct skeletons, respectively [9, 65].

Whenever a skeleton contains a pair of pointers x_i and x_j that point to the same node x_k (with $i \neq j \neq k$), the corresponding polynomial is not an isolated singularity. One should introduce a ‘link’ or a ‘pointing link’, which is an additional monomial involving x_i and x_j such that the singularity is isolated. The link is denoted by the dashed line connecting the two nodes

$$x_i^{m_i}x_k + x_j^{m_j}x_k + x_i^px_j^q : \begin{array}{ccccc} m_i & & m_k & & m_j \\ \circ & \xrightarrow{\quad} & \circ & \xleftarrow{\quad} & \circ \\ & & \text{---} & & \end{array} \quad (4.4)$$

The link, and therefore the singularity, exists only if there exists positive integers p and q such that the polynomial is quasi-homogeneous. We must then solve a Diophantine

Type		$F(x_1, x_2, x_3, x_4)$	Skeleton and links	Skeleton
XVI ₁	{16, 1}	$x_1^a x_2 + x_1 x_2^b + x_1 x_3^c + x_1 x_4^d$ $+ x_2^p x_3^q + x_3^r x_4^s + x_2^u x_4^v$		$S_{4,9}$
XVI ₂	{16, 2}	$x_1^a x_2 + x_1 x_2^b + x_1 x_3^c + x_1 x_4^d$ $+ x_2^p x_3^q x_4 + x_3^r x_4^s + x_2^u x_4^v$		$S_{4,9}$
XVI ₃	{16, 3}	$x_1^a x_2 + x_1 x_2^b + x_1 x_3^c + x_1 x_4^d$ $+ x_2^p x_3^q x_4 + x_3^r x_4^s + x_2^u x_3 x_4^v$		$S_{4,9}$
XVII ₁	{17, 1}	$x_1^a x_2 + x_1 x_2^b + x_2 x_3^c + x_1 x_4^d$ $+ x_1^p x_3^q + x_2^r x_4^s$		$S_{4,8}$
XVII ₂	{17, 2}	$x_1^a x_2 + x_1 x_2^b + x_2 x_3^c + x_1 x_4^d$ $+ x_1^p x_3^q + x_2^r x_3 x_4^s$		$S_{4,8}$
XVII ₃	{17, 3}	$x_1^a x_2 + x_1 x_2^b + x_2 x_3^c + x_1 x_4^d$ $+ x_1^p x_3^q + x_2^r x_4^s$		$S_{4,8}$
XVIII ₁	{18, 1}	$x_1^a x_3 + x_1 x_2^b + x_2 x_3^c + x_2 x_4^d + x_3^p x_4^q$		$S_{4,7}$
XIX	{19, 1}	$x_1^a x_3 + x_1 x_2^b + x_3^c x_4 + x_2 x_4^d$		$S_{4,5}$

Table 3. The KS-YY classification of isolated hypersurface singularities in \mathbb{C}^4 : types XVII to XIX.

equation determined by the scaling weight q_i of the four fields. For instance, in (4.4), if q_k is the scaling weight of x_j , we must have $p \frac{1-q_k}{m_i} + q \frac{1-q_k}{m_j} = 1$.

A pointing link is a link that is itself a pointer to another field y , of the form $x_i^p x_j^q y$. A link associated to a field x_k can only point to a field y that is not x_k nor any of the fields that are pointed at by the fields x_i and x_j involved in the link. For instance, we can have

$$x_i^{m_i} x_k + x_j^{m_j} x_k + x_i^p x_j^q y : \quad \begin{array}{ccccccc} m_i & & m_k & & m_j & & n \\ \circ & \xrightarrow{\quad} & \circ & \xleftarrow{\quad} & \circ & & \circ \\ & & & & & & \\ & & & & & & \end{array} \quad (4.5)$$

We denote the pointed link as shown here. As for simple links, pointed links can only exist if they preserve the quasi-homogeneity of the polynomial.

4.2 Computing the basic data of any canonical IHS: Mathematica code

Given any $F(x)$ defining a canonical threefold singularity \mathbf{X} , we can easily compute all the mathematical and physical quantities introduced in the previous sections, which we call ‘the basic data’ associated to \mathbf{X} . We implemented these computations on a computer using `Mathematica` [66], and we provided the `Mathematica` notebook `Basic-Data-IHS.nb` as an ancillary file in the supplementary material of this paper. Here, we describe the main routines in that notebook.

Basic characteristics of IHSs. The code implements many computations which are collated through the main routine:

$$\text{BasicIHSdata}[F(x_1, x_2, x_3, x_4), \{x_1, x_2, x_3, x_4\}]. \quad (4.6)$$

The code also generate any allowed polynomial $F(x)$ given a type, subtype and the positive integers (a, b, c, d) :

$$\text{FfromABCD}[\{t, s\}, \{a, b, c, d\}], \quad (4.7)$$

with $\{t, s\} = \{\text{type}, \text{subtype}\}$. The ‘basic data’ routine can be called directly as:

$$\text{BasicIHSdataFromABCD}[\{t, s\}, \{a, b, c, d\}]. \quad (4.8)$$

The output is given as a list with the following information:

- $F(x)$ itself;
- the scaling weights (q_i) ;
- r : the rank of the 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5\text{d}}$ (that is, the number of exceptional divisors in $\tilde{\mathbf{X}}$);
- f : the flavor rank preserved on the Coulomb branch of $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$ and $\mathcal{T}_{\mathbf{X}}^{5\text{d}}$;
- $d_H = \hat{r} + f$: the quaternionic Higgs branch dimension of $\mathcal{T}_{\mathbf{X}}^{5\text{d}}$;
- \hat{r} : the rank of the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$;
- $\hat{d}_H = r + f$: the quaternionic Higgs branch dimension of $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$;
- $\mu = 2\hat{r} + f$: the multiplicity (Milnor number) of \mathbf{X} ;
- modality: the modality $m(\mathbf{X})$ of \mathbf{X} ;
- n_h : the effective number of hypermultiplets of $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$;
- n_v : the effective number of vector multiplets of $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$;
- b_3 : the number of 3-cycles in the resolved threefold $\tilde{\mathbf{X}}$;
- $\Delta A \equiv \Delta \mathcal{A}_r$: the difference between the ‘virtual dimension’ $n_h - n_v$ of the HB of $\mathcal{T}_{\mathbf{X}}^{4\text{d}}$ and the actual HB dimension, $\Delta \mathcal{A}_r = n_h - n_v - \hat{d}_H$;

- g_C : the list $\{g_{1,2}, \dots\}$ of genera for the higher-genus ($g > 0$) complex curves at the base of the exceptional divisors inside $\tilde{\mathbf{X}}$;
- specL : the spectrum $\{\ell_l\}$ of the Milnor ring $\mathcal{M}(F)$;
- CB spectrum: the Coulomb branch spectrum of $\mathcal{T}_{\mathbf{X}}^{4d}$;
- a and c : the conformal anomaly coefficients of $\mathcal{T}_{\mathbf{X}}^{4d}$;
- marginal Defs: the number of marginal deformations of $\mathcal{T}_{\mathbf{X}}^{4d}$;
- $\text{Tor}H_2$: the torsional part of the second homology of the link L_5 , $\text{Tor } H_2(L_5, \mathbb{Z}) = \mathfrak{t}_2$.

Smoothing operator. The crepant resolution procedure described in section 2.2 is encoded in an additional routine, which constructs the (partially) resolved Calabi-Yau three-fold \mathbf{X} of any IHS $\{F(x_1, x_2, x_3, x_4) = 0\}$. Using the function:

$$\text{SmoothOperator}[F(x_1, x_2, x_3, x_4), \{x_1, x_2, x_3, x_4\}], \quad (4.9)$$

one gets an output in form of $\{\text{Resolution sequence, Resolved equation, Patches}\}$. Note that there can be several distinct resolution sequences for a given IHS, and the code only generates one of them.

The set “Resolution sequence” contains the crepant resolution sequence of the IHS, in the form $\{\text{Blp}_1, \text{Blp}_2, \dots, \text{Blp}_n\}$. Each Blp_i is a blow-up of the toric ambient space, given in the form $\{\{\{y_1, y_2, \dots, y_k\}, \{w_1, w_2, \dots, w_k\}\}, \delta\}$, which corresponds to the blow-up of the locus $y_1 = y_2 = \dots = y_k = 0$ with weights (w_1, w_2, \dots, w_k) , with the exceptional divisor $\delta = 0$. Using the notations of section 2.2, Blp_i is also written as $(y_1^{(w_1)}, y_2^{(w_2)}, \dots, y_k^{(w_k)}; \delta)$, and the superscripts w_i are omitted if they all equal to one.

The set “Patches” contains the coordinate patches of the toric ambient space T_Σ after the blow-ups, which are in one-to-one correspondence with the 4d cones of the toric fan Σ . This coordinate patch structure also extends to the anti-canonical hypersurface $\tilde{\mathbf{X}} = -K(T_\Sigma)$. On each coordinate patch $\{y_1, y_2, y_3, y_4\}$, where y_i are the local coordinates of $\tilde{\mathbf{X}}$, one can set $y_1 = y_2 = y_3 = y_4 = 0$ unless the intersection locus is not contained in $\tilde{\mathbf{X}}$.

We always use the notation $\delta_i = 0$ ($\delta[i] = 0$ in the code) for the exceptional divisors inside the blow-up of T_Σ . They may not intersect $\tilde{\mathbf{X}} = -K(T_\Sigma)$ eventually, and they may be reducible in $\tilde{\mathbf{X}}$. Hence, the total number of δ_i does not equal to the 5d rank r in general.

Note that this code also works in the generic case with residual terminal singularities. In these cases, the resolution sequence is a partial crepant resolution that gives rise to $\tilde{\mathbf{X}}$ with at most terminal singularities.

As an example, take $F = x_1^5 + x_2^2 x_1 + x_3^5 + x_4^2 x_3$. We input $\text{SmoothOperator}[x_1^5 + x_2^2 x_1 + x_3^5 + x_4^2 x_3, \{x_1, x_2, x_3, x_4\}]$, and get the output $\{\text{Resolution sequence, Resolved equation, Patches}\}$ where

$$\begin{aligned} \text{Resolution sequence} = \\ = \{ \{ \{ \{x_1, x_2, x_3, x_4\}, \{1, 1, 1, 1\}\}, \delta[1] \}, \{ \{ \{x_2, x_4, \delta[1]\}, \{1, 1, 1\}\}, \delta[2] \} \}, \end{aligned} \quad (4.10)$$

and after applying this to the singular IHS, we obtain the resolved equation (including the proper transforms):

$$\text{Resolved equation} = x_1 x_2^2 + x_3 x_4^2 + (x_1^5 + x_3^5) \delta[1]^2. \quad (4.11)$$

The patches are then given by

$$\begin{aligned} \text{Patches} &= \{\{x_1, x_2, x_3, \delta[1]\}, \{x_2, x_4, x_1, \delta[2]\}, \{x_2, \delta[1], x_1, \delta[2]\}, \{x_4, \delta[1], x_1, \delta[2]\}, \\ &= \{x_1, x_3, x_4, \delta[1]\}, \{x_2, x_4, x_3, \delta[2]\}, \{x_2, \delta[1], x_3, \delta[2]\}, \{x_4, \delta[1], x_3, \delta[2]\}\}. \end{aligned} \quad (4.12)$$

The resolution sequence and the resolved equation match (A.7)–(A.8).

5 Smoothable models

In the previous section, we spelled out a general approach to determine the basic properties of the SCFTs $\mathcal{T}_{\mathbf{X}}^{5d}$ and $\mathcal{F}_{\mathbf{X}}^{4d}$ for any IHS \mathbf{X} . We also explained how to explicitly construct the crepant resolution $\pi: \tilde{\mathbf{X}} \rightarrow \mathbf{X}$. This allows us to determine the smoothness properties of the resolved model $\tilde{\mathbf{X}}$ (whether it contains residual singularities or not), as well as possible 5d infrared gauge theory descriptions. In earlier works [1, 6], we studied interesting models with terminal singularities (whether \mathbf{X} was a terminal singularity with $r = 0$, or a singularity with $r > 0$ that contains remnant terminal singularities after a crepant resolution). In this section, we would like to focus on *smoothable models* — that is, singularities \mathbf{X} whose crepant resolution $\tilde{\mathbf{X}}$ is a smooth CY threefold (below, we also call such a CY threefold ‘a smooth model’).

5.1 Tabulating the smoothable models

We found it useful to organise the smoothable singularities \mathbf{X} into three categories:

1. Smooth models with smooth exceptional divisors and $b_3(\tilde{\mathbf{X}}) = 0$
2. Smooth models with smooth exceptional divisors and $b_3(\tilde{\mathbf{X}}) > 0$
3. Smooth models with singular divisors (for any $b_3(\tilde{\mathbf{X}})$).

It is also useful to organise the models by their 5d rank, r , i.e. by the number of exceptional divisors. For low rank, we gathered (conjecturally) all such models by brute force, by generating a large dataset of 15,142 physically distinct models with $r > 0$.⁸ The number of models for $r \leq 10$ in that dataset is given in table 4.

For $r = 1, 2, 3$, we resolved all the models explicitly and found all the smoothable models in our dataset. For larger r , this becomes impractical, but we can sharply reduce the number of cases to consider by realising that a necessary condition for having a smoothable model is that $\Delta\mathcal{A} \in \mathbb{Z}$, because $\Delta\mathcal{A} \notin \mathbb{Z}$ can only arise from residual terminal

⁸We considered all canonical IHS of type I to XIX with (a, b, c, d) ‘small enough’, and then identified the many models with identical scaling weights. For definiteness, we chose the cap $a, b, c, d \leq K$ with $K = 60$ for the 19 types (and for each subtype) in the KS-YY classification. This gives us a total of 39,094 physically distinct models. 61% of these models have $r = 0$, and we do not consider them further in this paper.

r	1	2	3	4	5	6	7	8	9	10
# of models	1167	649	693	666	506	549	486	495	491	461
# of $\Delta\mathcal{A}_r \in \mathbb{Z}$ models	161	128	143	131	113	154	98	117	137	117
# of smoothable models	12	24	32	23	32	39	23	41	?	?
# of $\mathbb{Z}\text{CB}$ models	3	6	10	9	12	18	7	9	12	12
# of ‘fully smooth’ models	3	5	10	7	9	14	5	8	7	7

Table 4. Number of low-rank models in our dataset. Here, $\mathbb{Z}\text{CB}$ refers to the models with integral CB dimensions and ‘fully smooth’ refers to the smoothable models with smooth exceptional divisors.

singularities [1]. We are also particularly interested in the smoothable models with smooth exceptional divisors (the ‘fully smooth models’, for short). We then make the following conjecture, based on explicit analysis of our dataset:

Conjecture 5.1 *For any smoothable model \mathbf{X} with smooth exceptional divisors, the associated 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ has an integral CB spectrum.*

Note that the converse is not true. On the other hand, we propose the following:

Conjecture 5.2 *Any model \mathbf{X} with an integral CB spectrum is smoothable (not necessarily with smooth divisors).*

In summary, we have the following inclusions of sets, at any fixed rank $r > 0$:

$$\{\text{‘fully smooth’}\} \subset \{\text{integral CB}\} \subset \{\text{smoothable}\} \subset \{\Delta\mathcal{A} \in \mathbb{Z}\} \subset \{\text{rank-}r \text{ models}\}, \quad (5.1)$$

where the first and second inclusions are conjectures. The number of models at fixed rank appears to be infinite, but we conjecture that the number of smoothable models of a fixed rank is finite. For small r , we also believe that all these models are included in our dataset.⁹ The numbers of models with various properties (in our dataset) are shown in table 4. For $r \leq 10$, we considered all the models with integral CB spectrum and extracted all the ‘fully smooth’ models from that set. They are discussed in detail in the following, and in appendix.

Models with $b_3 = 0$. We first discuss the smoothable models with smooth exceptional divisors with $b_3(\mathbf{X}) = 0$. We studied all such models with rank $r \leq 10$ in our dataset (assuming Conjecture 5.1 is true). We also considered the infinite series $\text{AD}[D_{2n}, D_{2n}]$, which engineers the 4d SCFT of type $[G, G'] = [D_{2n}, D_{2n}]$ [67]. These $b_3 = 0$ models are summarized in tables 8 and 9.

This class of models have the simplest 5d interpretation, as they correspond to 5d SCFTs $\mathcal{T}_{\mathbf{X}}^{5d}$ whose Coulomb-branch low-energy physics is described entirely by r 5d $\text{U}(1)$

⁹We do not have a proof that our lists of ‘smoothable’ and ‘fully smooth’ models up to $r = 10$ is complete within the infinite class of models defined by canonical IHSs. This is largely due to the fact that the correlation between rank r and the hypersurface equation types is not one we understand systematically at this point. It would be very interesting to determine more direct relations between the (a, b, c, d) ’s defining the equation $F(x) = 0$ and the rank r , if possible.

vector multiples, with a prepotential determined by the intersection numbers (2.49) in the smooth resolved threefold. Similarly, the 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ is such that the Higgs-branch low-energy physics consists of only $\hat{d}_H = r + f$ hypermultiplets [1].

Models with $b_3 > 0$. Similarly, the smoothable models with smooth exceptional divisors but $b_3 > 0$ are listed until rank $r = 10$ in tables 10, 11, 12, 13 and 14.

Note that, in all these tables, we color the labels $\{t, s\}\{a, b, c, d\}$ in (i) black for the cases with a Lagrangian $\mathcal{T}_{\mathbf{X}}^{4d}$; (ii) blue for the cases with a $D_p^b(G)$ trinion description studied in [6]; (iii) red for the cases with a generalised $D_p^b(G)$ description of $\mathcal{T}_{\mathbf{X}}^{4d}$ (to be discussed below); (iv) orange for the cases without a known description of $\mathcal{T}_{\mathbf{X}}^{4d}$ yet. For the case (i), the Lagrangian quiver description of $\mathcal{T}_{\mathbf{X}}^{4d}$ is written down in the table. For the other cases, we simply list the CB spectrum. Many of these models were already discussed in [1, 6].

The physical interpretation of the 3-cycles in $\tilde{\mathbf{X}}$ was already discussed in [1]. The 5d SCFT $\mathcal{T}_{\mathbf{X}}^{5d}$ has an ‘enhanced Coulomb branch’ with $\frac{1}{2}b_3$ free hypermultiplets at any generic point on the CB. The 4d SCFT $\mathcal{T}_{\mathbf{X}}^{4d}$ has an Higgs branch whose low-energy physics includes $\frac{1}{2}b_3$ free abelian vector multiplets in addition to the \hat{d}_H hypermultiplets.

Once we have a resolved threefold with 3-cycles, it is generally possible to perform a sort of geometric transition that results in a distinct local geometry [71], as we will see in some examples below. Such transitions change the asymptotic of the threefold. Physically, they correspond to mass deformations that lift the Higgs branch component of the enhanced CB. To precisely characterise the physics of enhanced Coulomb branches in 5d is an interesting question, which is left for future work (for a discussion of 4d enhanced Coulomb branches, see [72]).

Smoothable models with singular divisors. Finally, one can consider the smoothable models whose exceptional divisors are not all smooth. The interpretation of such singularities in the divisors (while the threefold itself is smooth) appears to be more subtle, and we hope to come back to this in future work. Here, we simply list all the smoothable models up to rank $r = 4$ in the tables 15, 16, 17, 18, 19, 20 and 21.

5.2 The $\text{AD}[E_7, E_7]$ model

We now illustrate one example of ‘fully smooth’ model with $b_3 = 0$, which turns out to have $r = 5$. We consider the geometry that is associated to the Argyres-Douglas-type theory $[G, G'] = [E_7, E_7]$ [67]. This is the $\text{Type}\{4, 1\}\{3, 3, 3, 3\}$ singularity, with basic data:

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_1^3 + x_2^3x_1 + x_3^3 + x_4^3x_3$	5	7	$\text{SU}(2)^6 \times \text{U}(1)$	28	21	12	0	0	—

(5.2)

Deformations and 4d SCFT. In this case, the CB spectrum of $\mathcal{T}_{\mathbf{X}}^{4d}$ contains operators with $\Delta = 7$ and $\Delta = 9$, but there is no operator with $\Delta = 8$. Although the CB spectrum is fully integral,

$$\Delta = \{2^4, 3^5, 4^4, 5^3, 6^2, 7^2, 9\}, \quad (5.3)$$

it cannot match any set of Casimir operators of Lie groups. Hence there is no Lagrangian description of $\mathcal{T}_{\mathbf{X}}^{4d}$. There is, however, a generalised quiver using $D_p^b(G)$ building blocks.

The conjectured 4d generalised $D_p(G)$ description is:¹⁰

$$[1] - \text{SU}(3) - \text{SU}(5) - \text{SU}(7) - D_2(\text{SO}(20)) - \text{SO}(8) - D_2(\text{SO}(8)), \quad (5.4)$$

where the $D_2(\text{SO}(2N))$ factor contributes the CB operators of dimensions $\Delta = 3, 5, 7, \dots, N-1$, and the $\text{SO}(2N)$ gauge groups contribute $\Delta = 2, 4, \dots, 2N-2, N$. The flavor symmetry of $D_2(\text{SO}(20))$ is $\text{SO}(22) \times \text{U}(1)$, of which we gauge an $\text{SU}(7) \times \text{SO}(8)$ subgroup. Note that $D_2(\text{SO}(8))$ is the E_6 MN theory [73, 74].

Resolutions and 5d SCFT. This is a fully resolvable IHS, and a resolution sequence is given by:

$$(x_1, x_2, x_3, x_4; \delta_1), (x_1, x_3, \delta_1; \delta_2), (\delta_1, \delta_2; \delta_3). \quad (5.5)$$

After the resolution, the equation

$$x_1^3 \delta_2 + x_2^3 x_1 \delta_1 + x_3^3 \delta_2 + x_4^3 x_3 \delta_1 = 0 \quad (5.6)$$

is smooth. The exceptional divisor $\delta_1 = 0$ has three components:

$$\delta_1 = 0 : \quad x_1^3 + x_4^3 = 0. \quad (5.7)$$

The exceptional divisors $\delta_2 = 0$ and $\delta_3 = 0$ are irreducible, thus resulting in total in a rank $r = 5$ theory. The intersection numbers of the compact surfaces are shown on the l.h.s. of figure 1. The three components of $\delta_1 = 0$ are S_2, S_3 and S_4 , which are all \mathbb{P}^2 s. S_1 corresponds to $\delta_3 = 0$ and S_5 corresponds to $\delta_2 = 0$.

The conventions for the intersection diagrams are as follows (similarly to the ones in [75, 76]):

- (1) Each oval denotes an irreducible compact divisor S_i . The number in the bracket after S_i is the self-triple intersection number S_i^3 .
- (2) The line segment between ovals S_i and S_j corresponds to the complete intersection curve $S_i \cdot S_j$. The numbers in the boxes at the end of the line segment describe the self-intersection numbers of the curves inside their respective divisors. The number on S_i corresponds to $S_i \cdot S_j^2$ while the number of S_j corresponds to $S_j \cdot S_i^2$.
- (3) For each triangle among three ovals S_i, S_j and S_k , we use a boxed number to denote the triple intersection number $S_i \cdot S_j \cdot S_k$.

To see the structure of the curves more clearly, we may do four flops by blowing up S_2, S_3, S_4, S_5 once and blowing down S_1 four times, which gives the local model shown on the right-hand-side of figure 1. Here, S_1 is a gdp_6 of type $4\mathbf{A}_1$, which is also known as the (resolution of) Cayley's cubic surface [77]. We plot the Mori cone generators of S_1 and their

¹⁰We thank Simone Giacomelli for discussions of this model and for suggesting this particular quiver.

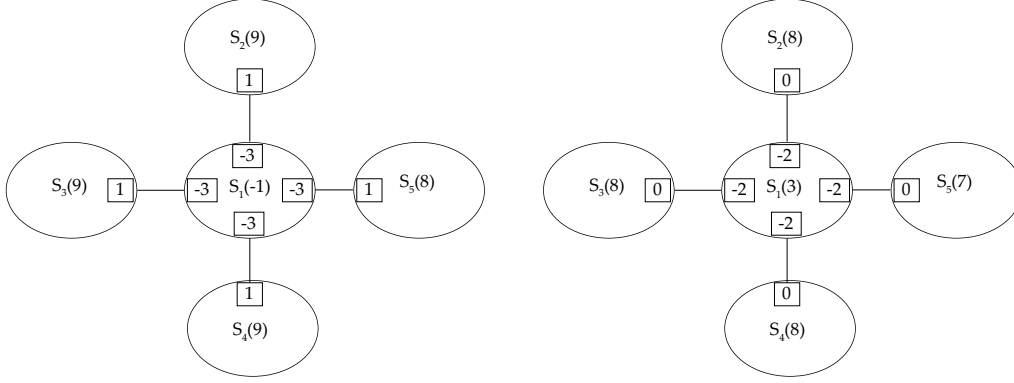


Figure 1. Resolved $\text{AD}[E_7, E_7]$ geometry: compact surfaces are S_i and gluing curves are links between these, with self-intersection numbers shown in the boxes. The l.h.s. is the model after resolution of the IHS, and the r.h.s. after the flop.

representations under the Picard group generators h, e_i ($h^2 = 1$, $h \cdot e_i = 0$, $e_i \cdot e_j = -\delta_{ij}$), showing the four (-2) -curves $S_1 \cdot S_2$, $S_1 \cdot S_3$, $S_1 \cdot S_4$ and $S_1 \cdot S_5$:

$$(5.8)$$

We also plot the curves on S_2 , S_3 , S_4 and S_5 :

$$(5.9)$$

Now we can combine the curves on each S_i and get the combined fiber diagram (CFD) introduced in [49, 78–80] of \mathcal{T}_X^{5d} by applying the procedure in [80]. The CFDs are — in brief — a collection of (usually rational) curves, which encode the flavor symmetries (in terms of the subgraph of green-colored nodes) and the (-1) -self-intersection curves that can be flopped and the associated hypermultiplet decoupled.

The CFD is constructed by connected to the same $S_i \cdot S_j$ are combined, with the proper multiplicities. For example, on S_1 the six (-1) -curves between the four (-2) -curves have multiplicities two, which correspond to (-2) nodes in the CFD. Hence the CFD can be read off as:



The UV flavor symmetry algebra is therefore $\mathfrak{g}_F = \mathfrak{su}(2)^6 \oplus \mathfrak{u}(1)$.

On the geometry, we can also identify a ruling structure by taking $S_1 \cdot S_2$ and $S_1 \cdot S_5$ to be the section curves, $S_1 \cdot S_3$ and $S_1 \cdot S_4$ to be the fiber. From the spectrum of massless hypermultiplets in the IR gauge theory limit, we can identify the following 5d gauge theory description:

$$\begin{array}{c} 1\mathbf{AS} + 1\mathbf{F} \\ | \\ \text{SU}(2)_0 \text{ --- } \text{SU}(4)_{-\frac{1}{2}} \text{ --- } \text{SU}(2)_0 \end{array} \quad (5.11)$$

If we decouple the two $\text{SU}(2)$ gauge factors by sending the volume of S_3 and S_4 to infinity, the surfaces S_1, S_2, S_5 give rise to IR gauge theory $\text{SU}(4)_{-\frac{1}{2}} + 1\mathbf{AS} + 5\mathbf{F}$. It would be interesting to see whether this IR gauge theory description can inform the computation of the magnetic quiver for this theory, and to then match this with the conjectured 4d quiver (5.4).

5.3 4d SCFTs: $D_p^b(G)$ -trinions

We now turn our attention ‘fully smooth’ models that have $b_3 > 0$. The 5d interpretation becomes more difficult in general, due to existence of an enhanced Coulomb branch. As before, the 4d SCFTs $\mathcal{T}_{\mathbf{X}}^{5d}$ may not have a Lagrangian description, as can be seen from tables 10, 11, 12, 13 and 14. Some of these models have an interpretation in terms of the $D_p^b(G)$ -trinions introduced in [6]. Using the notation for the trinions introduced therein, namely:

$$\begin{array}{c} D_{p_3}^{b_3}(G) \\ | \\ D_{p_2}^{b_2}(G) \text{ --- } G \text{ --- } D_{p_1}^{b_1}(G) \end{array} \quad (5.12)$$

the models with 4d trinion description are summarized in table 5.

5.4 4d SCFTs with $c - a < 0$

Another interesting class of models are those for which the conformal anomalies of $\mathcal{T}_{\mathbf{X}}^{4d}$ satisfy the inequality $c < a$. Such 4d SCFTs are slightly peculiar, and appear to be more

Type (a, b, c, d)	F	Trinion Representation	CB Spectrum
$\{1, 1\}\{3, 4, 4, 4\}$	$x_2^4 + x_3^4 + x_4^4 + x_1^3$	$ \begin{array}{c} D_2^{12} \\ \\ D_4^{12} \text{ --- } E_6 \text{ --- } D_4^{12} \end{array} $	$2^6, 3^7, 5^3, 6^6, 8, 9^3, 12$
$\{1, 1\}\{2, 3, 10, 10\}$	$x_3^{10} + x_4^{10} + x_2^3 + x_1^2$	$ \begin{array}{c} D_1^{24} \\ \\ D_{10}^{30} \text{ --- } E_8 \text{ --- } D_{10}^{30} \end{array} $	$2^7, 3^8, 5^6, 6^9, 8^5, 9^8, 11^4, 12^7, 14^3, 15^6, 17^2, 18^5, 20, 21^4, 24^3, 27^2, 30$
$\{1, 1\}\{2, 5, 6, 6\}$	$x_3^6 + x_4^6 + x_2^5 + x_1^2$	$ \begin{array}{c} D_1^{20} \\ \\ D_6^{30} \text{ --- } E_8 \text{ --- } D_6^{30} \end{array} $	$2^3, 3^4, 4^5, 5^4, 7^2, 8^3, 9^4, 10^5, 12, 13^2, 14^3, 15^4, 18, 19^2, 20^3, 24, 25^2, 30$
$\{2, 1\}\{3, 6, 9, 2\}$	$x_3^9 + x_4^2 x_3 + x_2^6 + x_1^3$	$ \begin{array}{c} D_2^{18} \\ \\ D_3^{18} \text{ --- SO}(20) \text{ --- } D_6^{18} \end{array} $	$2^5, 3^4, 4^7, 5^3, 6^5, 7^4, 8^3, 9^3, 10^4, 11, 12^3, 13, 14, 15, 16, 18$
$\{2, 1\}\{4, 4, 8, 2\}$	$x_3^8 + x_4^2 x_3 + x_1^4 + x_2^4$	$ \begin{array}{c} D_2^{16} \\ \\ D_4^{16} \text{ --- SO}(18) \text{ --- } D_4^{16} \end{array} $	$2^8, 4^8, 5^2, 6^6, 8^6, 9, 10^3, 12^3, 14, 16$
$\{2, 1\}\{2, 5, 5, 6\}$	$x_3^5 + x_4^6 x_3 + x_1^2 + x_2^5$	$ \begin{array}{c} D_1^{20} \\ \\ D_6^{24} \text{ --- } E_8 \text{ --- } D_5^{30} \end{array} $	$2^7, 4^5, 6^7, 8^6, 10^4, 12^6, 14^4, 16^2, 18^4, 20^2, 22, 24^2, 26, 30$
$\{2, 1\}\{3, 6, 12, 2\}$	$x_3^{12} + x_4^2 x_3 + x_1^3 + x_2^6$	$ \begin{array}{c} D_2^{24} \\ \\ D_6^{24} \text{ --- SO}(26) \text{ --- } D_3^{24} \end{array} $	$2^9, 4^9, 5^2, 6^8, 8^8, 9, 10^6, 12^6, 13, 14^4, 16^4, 18^2, 20^2, 22, 24$

Table 5. Models with integral spectrum and a 4d SCFT description as a $D_p^b(G)$ trinion.

Type	$F(x, y, z, t)$	$r = \widehat{d}_H$	f	$d_H = \widehat{r}$	$24(c - a)$	b_3	\mathfrak{t}_2
$\{1, 1\}\{2, 5, 5, 5\}$	$x^2 + y^5 + z^5 + t^5$	2	0	32	-4	12	\mathbb{Z}_2^{12}
$\{1, 1\}\{2, 3, 8, 8\}$	$x^2 + y^3 + z^8 + t^8$	2	0	49	-1	6	\mathbb{Z}_3^6
$\{1, 1\}\{2, 3, 9, 9\}$	$x^2 + y^3 + z^9 + t^9$	3	0	64	-4	14	\mathbb{Z}_2^{14}
$\{1, 1\}\{3, 4, 4, 4\}$	$x^3 + y^4 + z^4 + t^4$	3	0	27	-3	12	\mathbb{Z}_3^6
$\{2, 1\}\{2, 7, 7, 3\}$	$x^2 + y^7 + z^7 + t^3 z$	3	0	45	-3	12	\mathbb{Z}_2^{12}
$\{1, 1\}\{2, 3, 10, 10\}$	$x^2 + y^3 + z^{10} + t^{10}$	5	0	81	-3	16	\mathbb{Z}_3^8
$\{2, 1\}\{2, 3, 15, 7\}$	$x^2 + y^3 + z^{15} + t^7 z$	5	0	91	-1	12	\mathbb{Z}_2^{12}
$\{1, 1\}\{2, 5, 6, 6\}$	$x^2 + y^5 + z^6 + t^6$	6	0	50	-2	16	\mathbb{Z}_5^4

Table 6. Table of 4d SCFTs with $c < a$ engineered from IHS in IIB.

‘rare’; from the point of view of the superconformal index, $c - a < 0$ is analogous to having $c < 0$ for a 2d CFT — see e.g. [81–84]. It turns out that there are only 8 canonical IHS for which $c < a$ in our database, and they are all smoothable models with $b_3 > 0$. Their basic data is shown in table 6.

The 4d Coulomb branch spectrum of all these singularities except for the type $\{1, 1\}\{2, 3, 8, 8\}$ is integer-valued, $\Delta_\alpha \in \mathbb{Z}$. (And these 7 models are actually ‘fully smooth’.) Note that we have:

$$24(c - a) = n_h - n_v = r + f - \frac{1}{2}b_3, \quad (5.13)$$

in agreement with the physical interpretation of the 3-cycles as vector multiplets in IIB. We will now provide the corresponding 4d SCFTs $\mathcal{T}_{\mathbf{X}}^{4d}$ (whose description is simply related to the electric and magnetic quiverines in 3d, since $f = 0$ [1, 6]). We also study the resolution of these singularities and identify the 5d gauge theory phase if it exists. Note that most of these theories are non-Lagrangian 4d SCFTs, and that in general we do not have a Lagrangian description of the 3d quiverines either.

5.4.1 Rank 2: $\{1, 1\}\{2, 5, 5, 5\}$

This case was also discussed in [1]. We have the Lagrangian description of $\mathcal{T}_{\mathbf{X}}^{4d}$ as the quiver:

$$\begin{array}{ccccccc}
 & & \text{Spin}(1) & & & & \\
 & & \bullet & & & & \\
 & & | & & & & \\
 & & \text{Sp}(2) & & & & \\
 & & \bullet & & & & \\
 & & | & & & & \\
 \text{Sp}(3) & & \bullet & & \text{Sp}(4) & & \text{Sp}(3) & & \text{Sp}(2) & & \text{Sp}(1) \\
 \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\
 \text{Spin}(5) & \text{Spin}(11) & \text{Spin}(9) & \text{Spin}(7) & \text{Spin}(5) & \text{Spin}(3) & & & & &
 \end{array} \quad (5.14)$$

Note that the six $\text{Spin}(2n + 1)$ ($n > 0$) gauge groups give rise to a $\mathfrak{f} = \mathbb{Z}_2^6$ 1-form symmetry (the vector of $\text{Spin}(2n + 1)$ does not break the \mathbb{Z}_2 center), in agreement with the geometric determination of \mathfrak{t}_2 .

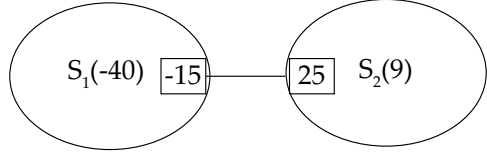
Let us also discuss the resolution of this model. Starting with the hypersurface equation $x_1^2 + x_2^5 + x_3^5 + x_4^5 = 0$, the resolution sequence is

$$(x_1^{(2)}, x_2^{(1)}, x_3^{(1)}, x_4^{(1)}; \delta_1), (x_1, \delta_1; \delta_2), (\delta_1, \delta_2; \delta_3) \quad (5.15)$$

The resolved equation is

$$x_1^2 \delta_2 + (x_2^5 + x_3^5 + x_4^5) \delta_1 = 0. \quad (5.16)$$

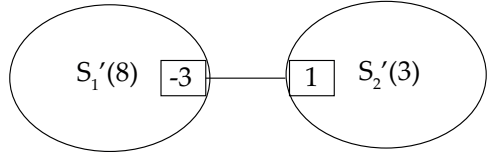
The set $\delta_1 = 0$ is empty, so there are only two exceptional divisors $S_1 : \delta_2 = 0$ and $S_2 : \delta_3 = 0$. The triple intersection numbers are



$$\begin{array}{c} \text{S}_1(-40) \quad \boxed{-15} \text{---} \boxed{25} \quad \text{S}_2(9) \end{array} \quad (5.17)$$

Hence $S_1 \cdot S_2$ is a genus-6 curve, and S_1 is a ruled surface over that curve. Note that S_2 is a \mathbb{P}^2 and $S_1 \cdot S_2$ has equation $\delta_2 = \delta_3 = x_2^5 + x_3^5 + x_4^5 = 0$, which is a degree-5 curve on \mathbb{P}^2 with genus $g = 6$. Since S_1 has 12 3-cycles, the total b_3 of the resolved CY3 is $b_3(\tilde{\mathbf{X}}) = 12$.

The presence of the 3-cycles in $\tilde{\mathbf{X}}$ makes the M-theory interpretation more subtle, as mentioned above. We can obtain a distinct and more ‘conventional’ model by performing a geometric transition, as described in [71], to obtain a new Calabi-Yau threefold with distinct asymptotics and $b_3 = 0$. In the present example, one performs a complex structure deformation of the geometry, such that the intersection curve $S_1 \cdot S_2$ has six double point singularities. Then we perform flop operations six times by blowing up six double points on S_2 . After these flops, the triple intersection numbers of the new surfaces S'_1 and S'_2 are



$$\begin{array}{c} \text{S}'_1(8) \quad \boxed{-3} \text{---} \boxed{1} \quad \text{S}'_2(3) \end{array} \quad (5.18)$$

Hence S'_1 is a \mathbb{F}_3 and S'_2 is a dP_6 . These triple intersection numbers exactly give rise to the 5d SCFT with IR description $\text{SU}(3)_{9/2} + 5\mathbf{F}$, $\text{Sp}(2) + 3\mathbf{F} + 2\mathbf{AS}$ or $G_2 + 5\mathbf{F}$, and UV flavor symmetry $G_F^{5d} = \text{Sp}(6)$. Before the flops, the geometry gives $f = 0$. After the flops, there are six new 2-cycles in the flopped geometry, which give rise to the flavor rank $f' = 6$ that matches G_F^{5d} .

5.4.2 Rank 2: $\{1, 1\}\{2, 3, 8, 8\}$

The 4d CB spectrum of this model is non-integral. In [6], a non-Lagrangian description of $\mathcal{T}_{\mathbf{X}}^{4d}$ is found, where D_p^b denotes the $D_p^b(G)$ theory [74, 85, 86], here with $G = E_6$ for the central gauge group:

$$\begin{array}{c} D_8^{12} \\ | \\ \mathcal{T}_{\mathbf{X}}^{4d} = D_8^{12} \text{---} E_6 \text{---} D_1^9 \end{array} \quad (5.19)$$

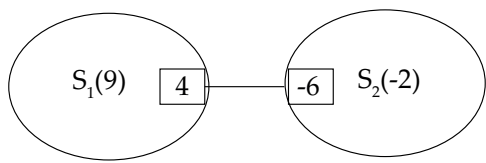
The singularity $x_1^2 + x_2^3 + x_3^8 + x_4^8 = 0$ is fully resolved by

$$(x_1^{(3)}, x_2^{(2)}, x_3^{(1)}, x_4^{(1)}; \delta_1), (x_1, x_2, \delta_1; \delta_2). \quad (5.20)$$

The resolved equation is completely smooth:

$$x_1^2 + x_2^3 \delta_2 + x_3^8 \delta_1^2 + x_4^8 \delta_1^2 = 0. \quad (5.21)$$

The triple intersection numbers of $S_1 : \delta_1 = 0$ and $S_2 : \delta_2 = 0$ are computed to be



$$S_1(9) \quad \boxed{4} \quad \text{---} \quad \boxed{-6} \quad S_2(-2) \quad (5.22)$$

These numbers appear to match the rank-2 5d SCFT with an $SU(3)_{1/2} + 9\mathbf{F}$ IR description and $G_F = SO(20)$ flavor symmetry. The intersection curve $S_1 \cdot S_2$ is an irreducible curve.

However, this model has flavor rank $f = 0$, which leads to a contradiction. The reason is that the equation of S_2 is

$$x_1^2 + x_3^8 \delta_1^2 + x_4^8 \delta_1^2 = 0. \quad (5.23)$$

At a generic locus $\delta_1 \neq 0$, $x_1 \neq 0$, the equation is a smooth genus-3 curve, because the Newton polytope with vertices $(2, 0, 0)$, $(0, 8, 0)$ and $(0, 0, 8)$ has exactly 3 interior points. On the other hand, S_2 is singular at the point

$$x_1 = x_3 = x_4 = \delta_2 = 0. \quad (5.24)$$

Note that this point is not a singular point in the resolved CY3 $\tilde{\mathbf{X}}$, but it is a singular point on S_2 .

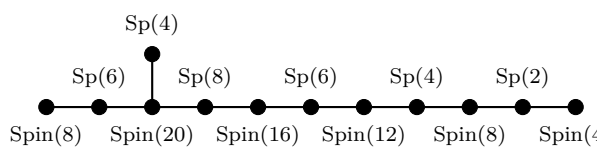
From the genus-3 curve on S_2 , we have

$$b_3(S_2) = b_3(\tilde{\mathbf{X}}) = 6, \quad (5.25)$$

although $S_1 \cdot S_2$ has genus 0. The 5d physical interpretation of this geometry is unclear.

5.4.3 Rank 3: $\{1, 1\}\{2, 3, 9, 9\}$

The 4d SCFT has the Lagrangian description:



$$\begin{array}{ccccccccc} & & \text{Sp}(4) & & & & & & \\ & & | & & & & & & \\ \text{Sp}(6) & & \text{Sp}(8) & & \text{Sp}(6) & & \text{Sp}(4) & & \text{Sp}(2) \\ | & | & | & | & | & | & | & | & | \\ \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\ \text{Spin}(8) & \text{Spin}(20) & \text{Spin}(16) & \text{Spin}(12) & \text{Spin}(8) & \text{Spin}(4) & & & \end{array} \quad (5.26)$$

It is a simple exercise to check that this quiver has a 1-form symmetry $\mathfrak{f} = \mathbb{Z}_2^7$, in agreement with the geometric determination. The resolution of the singularity $x_1^2 + x_2^3 + x_3^9 + x_4^9 = 0$ is achieved by

$$(x_1^{(3)}, x_2^{(2)}, x_3^{(1)}, x_4^{(1)}; \delta_1), (x_1, x_2, \delta_1; \delta_2), (x_1, \delta_2; \delta_3), (\delta_2, \delta_3; \delta_4). \quad (5.27)$$

The resolved equation is

$$x_1^2 \delta_3 + x_2^3 \delta_2 + (x_3^9 + x_4^9) \delta_1^3 \delta_2 = 0. \quad (5.28)$$

The set $\delta_2 = 0$ is empty. The other surface components are irreducible. The triple intersection numbers of $S_1 : \delta_1 = 0$, $S_2 : \delta_3 = 0$ and $S_3 : \delta_4 = 0$ are

$$\begin{array}{c}
 \text{S}_1(9) \quad \boxed{1} \text{ --- } \boxed{-3} \quad \text{S}_3(8) \quad \boxed{27} \text{ --- } \boxed{-15} \quad \text{S}_2(-48)
 \end{array} \quad (5.29)$$

The intersection curve $S_2 \cdot S_3$ is a genus-7 curve, and the surface S_2 is ruled over that curve. S_1 is a \mathbb{P}^2 and S_4 is a \mathbb{F}_3 . There are 14 3-cycles on S_3 , and $b_3(\mathbf{X}) = 14$. After the flop that results from blowing up S_3 at seven double points on $S_2 \cdot S_3$, the triple intersection numbers become

$$\begin{array}{c}
 \text{S}'_1(9) \quad \boxed{1} \text{ --- } \boxed{-3} \quad \text{S}'_3(1) \quad \boxed{-1} \text{ --- } \boxed{-1} \quad \text{S}'_2(8)
 \end{array} \quad (5.30)$$

The model does not appear to have a 5d gauge-theory description.

5.4.4 Rank 3: $\{1, 1\}\{3, 4, 4, 4\}$

A non-Lagrangian description of $\mathcal{T}_{\mathbf{X}}^{4d}$ was found in [6]:

$$\begin{array}{c}
 D_4^{12} \\
 | \\
 \mathcal{T}_{\mathbf{X}}^{4d} = D_2^{12} \text{ --- } E_6 \text{ --- } D_4^{12}
 \end{array} \quad (5.31)$$

The singularity $x_1^3 + x_2^4 + x_3^4 + x_4^4 = 0$ leads to the resolution sequence

$$(x_1, x_2, x_3, x_4; \delta_1), (x_1, \delta_1; \delta_2), (\delta_1, \delta_2; \delta_3), (\delta_1, \delta_3; \delta_4). \quad (5.32)$$

The resolved equation is

$$x_1^3 \delta_2^2 \delta_3 + (x_2^4 + x_3^4 + x_4^4) \delta_1 = 0. \quad (5.33)$$

The set $\delta_1 = 0$ is empty, hence there are only three exceptional divisors $S_1 : \delta_2 = 0$, $S_2 : \delta_3 = 0$ and $S_3 : \delta_4 = 0$. The triple intersection numbers are

$$\begin{array}{c}
 \text{S}_1(-16) \quad \boxed{-8} \text{ --- } \boxed{12} \quad \text{S}_2(-16) \quad \boxed{-12} \text{ --- } \boxed{16} \quad \text{S}_3(9)
 \end{array} \quad (5.34)$$

Hence the intersection curves $S_1 \cdot S_2$ and $S_2 \cdot S_3$ both have genus $g = 3$, and both S_1 and S_2 are ruled over genus-three curves. The divisor S_3 is a \mathbb{P}^2 . Since S_1 and S_2 both have 6 3-cycles, we have $b_3(\tilde{\mathbf{X}}) = 12$. If we perform flops by blowing up S_3 at three double points on $S_2 \cdot S_3$, we end up with

$$(5.35)$$

5.4.5 Rank 3: $\{2, 1\}\{2, 7, 7, 3\}$

The 4d SCFT for this model has the Lagrangian description:

$$(5.36)$$

This quiver has a 1-form symmetry $\mathfrak{f} = \mathbb{Z}_2^6$. The singularity $x_1^2 + x_2^7 + x_3^7 + x_4^3 x_3 = 0$ has a resolution sequence

$$(x_1^{(2)}, x_2^{(1)}, x_3^{(1)}, x_4^{(1)}; \delta_1), (x_1, x_4, \delta_1; \delta_2), (x_1, \delta_2; \delta_3), (\delta_2, \delta_3; \delta_4) \quad (5.37)$$

The resolved equation is

$$x_1^2 \delta_3 + x_2^7 \delta_1^3 \delta_2 + x_3^7 \delta_1^3 \delta_2 + x_4^3 x_3 \delta_2 = 0. \quad (5.38)$$

The triple intersection numbers of $S_1 : \delta_1 = 0$, $S_2 : \delta_3 = 0$, $S_3 : \delta_4 = 0$ are

$$(5.39)$$

S_3 is ruled over a genus-6 curve $S_3 \cdot S_4$. After the flop by blowing up $S_3 \cdot S_4$ at the six double points on S_4 , the new intersection numbers are

$$(5.40)$$

Since the assignment of base and ruling curves on S_i cannot match, this geometry does not have a gauge theory description.

5.4.6 Rank 5: $\{1, 1\}\{2, 3, 10, 10\}$

As we noted earlier this model has a trinion representation in 4d

$$\begin{array}{c} D_2^{24} \\ | \\ D_6^{24} \text{---SO}(26)\text{---} D_3^{24} \end{array} \quad (5.41)$$

The singularity $x_1^2 + x_2^3 + x_3^{10} + x_4^{10} = 0$ has a resolution sequence

$$\begin{aligned} (x_1^{(3)}, x_2^{(2)}, x_3^{(1)}, x_4^{(1)}; \delta_1), & \quad (x_1, x_2, \delta_1; \delta_2), & \quad (x_1, \delta_2; \delta_3), & \quad (x_2, \delta_3; \delta_4) \\ (\delta_2, \delta_3; \delta_5), & & (\delta_3, \delta_4; \delta_6), & \quad (\delta_3, \delta_6; \delta_7) \end{aligned} \quad (5.42)$$

The resolved equation is

$$x_1^2 \delta_3 + x_2^3 \delta_2 \delta_4^2 \delta_6 + (x_3^{10} + x_4^{10}) \delta_1^4 \delta_2^2 \delta_3 \delta_5^2 = 0. \quad (5.43)$$

The sets $\delta_2 = 0$ and $\delta_3 = 0$ is empty. The other surface components are irreducible. The Coulomb branch rank is $r = 5$. The triple intersection numbers of $S_1 : \delta_1 = 0$, $S_2 : \delta_4 = 0$, $S_3 : \delta_5 = 0$, $S_4 : \delta_6 = 0$ and $S_5 : \delta_7 = 0$ are

$$\begin{array}{c} \text{S}_1(9) \quad \boxed{1} \quad \text{S}_3(8) \quad \boxed{3} \quad \text{S}_5(8) \quad \boxed{20} \quad \text{S}_4(-24) \quad \boxed{14} \quad \text{S}_2(-24) \\ \text{---} \boxed{-3} \text{---} \boxed{-5} \text{---} \boxed{-14} \text{---} \boxed{-8} \end{array} \quad (5.44)$$

The surface components S_2 and S_4 are ruled over a genus-4 curve, and the other surface components are rational. On each of S_2 and S_4 , there are eight 3-cycles. In total we have $b_3(\tilde{\mathbf{X}}) = 16$. After a flop performed by blowing up four double points on $S_5 \cdot S_4$, we obtain the geometry

$$\begin{array}{c} \text{S}'_1(9) \quad \boxed{1} \quad \text{S}'_3(8) \quad \boxed{3} \quad \text{S}'_5(4) \quad \boxed{4} \quad \text{S}'_4(8) \quad \boxed{6} \quad \text{S}'_2(8) \\ \text{---} \boxed{-3} \text{---} \boxed{-5} \text{---} \boxed{-6} \text{---} \boxed{-8} \end{array} \quad (5.45)$$

5.4.7 Rank 5: $\{2, 1\}\{2, 3, 15, 7\}$

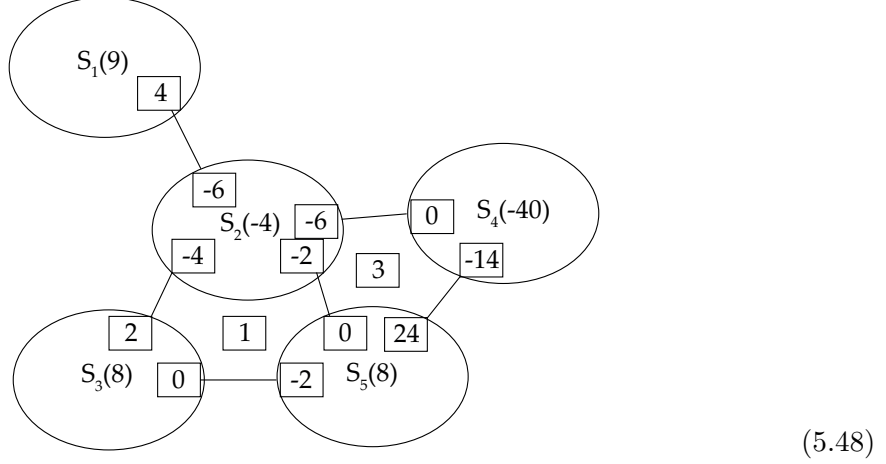
For this model, we did not find a Lagrangian nor a generalised quiver description. The equation $x_1^2 + x_2^3 + x_3^{15} + x_4^7 x_3 = 0$ leads to the resolution sequence

$$\begin{aligned} (x_1^{(3)}, x_2^{(2)}, x_3^{(1)}, x_4^{(1)}; \delta_1), & \quad (x_1, x_2, \delta_1; \delta_2), & \quad (x_1^{(2)}, x_2^{(1)}, x_4^{(1)}, \delta_2^{(1)}; \delta_3), & \quad (x_1, x_2, \delta_3; \delta_4) \\ (x_1, \delta_4; \delta_5), & & (\delta_4, \delta_5; \delta_6) \end{aligned} \quad (5.46)$$

The resolved equation is

$$x_1^2 \delta_5 + x_2^7 \delta_1^8 \delta_2^{13} \delta_3^{16} \delta_4^{21} \delta_5^{20} \delta_6^{19} + x_3^{15} \delta_1^9 \delta_2^7 \delta_3^3 \delta_4 + x_4^7 x_3 \delta_1^2 \delta_3^3 \delta_4 = 0. \quad (5.47)$$

The set $\delta_4 = 0$ is empty, and the other divisors are irreducible, hence the 5d rank is $r = 5$. The triple intersection numbers of $S_1 : \delta_1 = 0$, $S_2 : \delta_2 = 0$, $S_3 : \delta_3 = 0$, $S_4 : \delta_5 = 0$ and $S_5 : \delta_6 = 0$ are



S_4 is ruled over a genus-6 curve, the other surfaces are rational. There are 12 3-cycles on S_4 , and $b_3(\tilde{\mathbf{X}}) = 12$.

5.4.8 Rank 6: $\{1, 1\}\{2, 5, 6, 6\}$

This 4d SCFT for this model has a trinion description [6]:

$$\begin{array}{c} D_1^{20} \\ | \\ \mathcal{T}_{\mathbf{X}}^{4d} = D_6^{30} \text{ --- } E_8 \text{ --- } D_6^{30} \end{array} \quad (5.49)$$

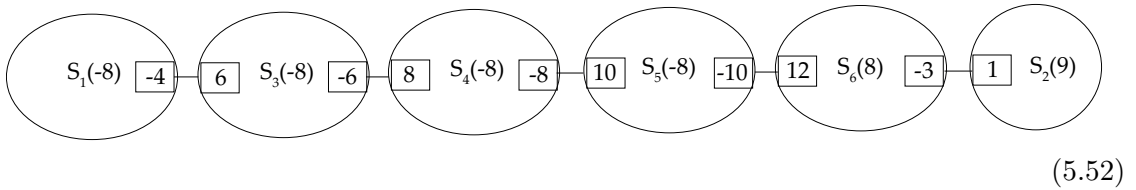
The singularity $x_1^2 + x_2^5 + x_3^6 + x_4^6 = 0$ leads to the resolution sequence

$$\begin{aligned} & (x_1^{(2)}, x_2^{(1)}, x_3^{(1)}, x_4^{(1)}; \delta_1), \quad (x_1, \delta_1; \delta_2), \quad (x_2, \delta_2; \delta_3), \quad (\delta_1, \delta_2; \delta_4) \\ & (\delta_2, \delta_3; \delta_5), \quad (\delta_2, \delta_5; \delta_6), \quad (\delta_2, \delta_6; \delta_7), \quad (\delta_2, \delta_7; \delta_8) \end{aligned} \quad (5.50)$$

The resolved equation is

$$x_1^2 \delta_2 + x_2^5 \delta_1 \delta_3^4 \delta_5^3 \delta_6^2 \delta_7 + (x_3^6 + x_4^6) \delta_1^2 \delta_2 \delta_4^2 = 0. \quad (5.51)$$

The sets $\delta_1 = 0$ and $\delta_2 = 0$ are empty, and the other divisors are irreducible, so that $r = 6$. The triple intersection numbers of $S_1 : \delta_3 = 0$, $S_2 : \delta_4 = 0$, $S_3 : \delta_5 = 0$, $S_4 : \delta_6 = 0$, $S_5 : \delta_7 = 0$ and $S_6 : \delta_8 = 0$ are



The divisors S_1, S_3, S_4 and S_5 are all ruled over a genus-2 curve. S_6 is \mathbb{F}_3 and S_2 is \mathbb{P}^2 . S_1, S_3, S_4 and S_5 each has four 3-cycles, and thus $b_3(\tilde{\mathbf{X}}) = 16$. After the flop performed by blowing up two double points on $S_6 \cdot S_5$, we obtain the geometry

$$(5.53)$$

5.5 4d SCFTs with $a = c$

Another noteworthy class of models consists of the smoothable models such that $a = c$. They all satisfy

$$f = 0, \quad r = \frac{1}{2}b_3 = \hat{d}_H, \quad \hat{r} = d_H. \quad (5.54)$$

Therefore, any such model with $r > 0$ is either smoothable or it has residual terminal singularities that correspond to some $a = c$ model.¹¹ They also all have higher-form symmetries (\mathfrak{t}_2 non-trivial). These cases include the following infinite series, which are called $E_k^{(l)}$ models in [6]. We have $E_6^{(3n-2)}$ and $E_6^{(3n-1)}$ models:

Type(a, b, c, d)	$E_k^{(l)}$	F	r	f	d_H	$a = c$	b_3	\mathfrak{t}_2
$\{1, 1\}\{3, 3, 3, 3n+1\}$	$E_6^{(3n-2)}$	$x_1^3 + x_2^3 + x_3^3 + x_4^{3n+1}$	n	0	$12n$	$2n(3n+2)$	$2n$	\mathbb{Z}_{3n+1}^2
$\{1, 1\}\{3, 3, 3, 3n+2\}$	$E_6^{(3n-1)}$	$x_1^3 + x_2^3 + x_3^3 + x_4^{3n+2}$	n	0	$12n+4$	$(2n+2)(3n+1)$	$2n$	\mathbb{Z}_{3n+2}^2

The model $E_6^{(l)}$ have a 4d $G = \mathrm{SU}(l+3)$ trinion description [6]:

$$\begin{array}{c} D_3 \\ | \\ D_3 \text{ --- } \mathrm{SU}(l+3) \text{ --- } D_3 \end{array} \quad (5.55)$$

We similarly have the $E_7^{(4n-3)}$ and $E_7^{(4n-1)}$ models:

Type(a, b, c, d)	$E_k^{(l)}$	F	r	f	d_H	$a = c$	b_3	\mathfrak{t}_2
$\{1, 1\}\{2, 4, 4, 4n+1\}$	$E_7^{(4n-3)}$	$x_1^2 + x_2^4 + x_3^4 + x_4^{4n+1}$	n	0	$18n$	$6n(2n+1)$	$2n$	\mathbb{Z}_{4n+1}^2
$\{1, 1\}\{2, 4, 4, 4n+3\}$	$E_7^{(4n-1)}$	$x_1^2 + x_2^4 + x_3^4 + x_4^{4n+3}$	n	0	$18n+9$	$6(n+1)(2n+1)$	$2n$	\mathbb{Z}_{4n+3}^2

as well as the $E_8^{(6n-5)}$ and $E_8^{(6n-1)}$ models:

Type(a, b, c, d)	$E_k^{(l)}$	F	r	f	d_H	$a = c$	b_3	\mathfrak{t}_2
$\{1, 1\}\{2, 3, 6, 6n+1\}$	$E_8^{(6n-5)}$	$x_1^2 + x_2^3 + x_3^6 + x_4^{6n+1}$	n	0	$30n$	$10n(3n+1)$	$2n$	\mathbb{Z}_{6n+1}^2
$\{1, 1\}\{2, 3, 6, 6n+5\}$	$E_8^{(6n-1)}$	$x_1^2 + x_2^3 + x_3^6 + x_4^{6n+5}$	n	0	$30n+20$	$10(n+1)(3n+2)$	$2n$	\mathbb{Z}_{6n+5}^2

¹¹When there are several terminal singularities, the residual IR theory on the 4d HB can be a direct sum of irreducible theories so that the total central charges satisfy $a = c$.

The $E_7^{(l)}$ and $E_8^{(l)}$ models have a 4d $SU(l+4)$ and $SU(l+6)$ trinion description, respectively:

$$\begin{array}{ccc}
 & D_4 & \\
 & | & \\
 D_2 & \text{---} SU(l+4) \text{---} & D_4
 \end{array}
 \qquad
 \begin{array}{ccc}
 & D_6 & \\
 & | & \\
 D_2 & \text{---} SU(l+6) \text{---} & D_3
 \end{array}
 \quad (5.56)$$

These trinion models were also recently discussed in [87]. We also find an interesting infinite family of $a = c$ models with $r = 0$:

Type(a, b, c, d)	F	r	f	d_H	$a = c$	b_3	t_2
$\{2, 1\}\{2, 2k+1, 2k+1, 2\}$	$x_1^2 + x_2^{2k+1} + x_3^{2k+1} + x_4^2 x_3$	0	0	$2k(k+1)$	$\frac{1}{9}k(k+1)(4k+5)$	0	\mathbb{Z}_2^{2k}

Finally, we find six ‘sporadic’ models with $r > 0$:

Type(a, b, c, d)	F	r	f	d_H	$a = c$	b_3	t_2
$\{1, 1\}\{2, 3, 7, 14\}$	$x_1^2 + x_2^3 + x_3^7 + x_4^{14}$	3	0	78	286	6	\mathbb{Z}_3^6
$\{1, 1\}\{2, 3, 7, 21\}$	$x_1^2 + x_2^3 + x_3^7 + x_4^{21}$	6	0	120	860	12	\mathbb{Z}_2^{12}
$\{2, 1\}\{3, 5, 5, 2\}$	$x_1^3 + x_2^5 + x_3^5 + x_3 x_4^2$	2	0	24	34	4	\mathbb{Z}_3^4
$\{2, 1\}\{3, 7, 7, 2\}$	$x_1^3 + x_2^7 + x_3^7 + x_3 x_4^2$	6	0	48	176	12	\mathbb{Z}_3^6
$\{2, 1\}\{4, 5, 5, 2\}$	$x_1^4 + x_2^5 + x_3^5 + x_3 x_4^2$	6	0	36	126	6	\mathbb{Z}_4^4
$\{3, 1\}\{2, 11, 3, 4\}$	$x_1^2 + x_2^{11} + x_4^4 x_3 + x_4 x_3^3$	5	0	60	230	10	\mathbb{Z}_2^{10}

(5.57)

All of the six models have a crepant resolution with no terminal singularities. Among these models, the type $\{2, 1\}\{4, 5, 5, 2\}$ has a crepant resolution with no singular divisor, see the entry in table 12. The other five models have singular divisors in their crepant resolution. We will now discuss the ‘fully smooth’ model in more detail.

5.5.1 Rank 6: $\{2, 1\}\{4, 5, 5, 2\}$

The theory does not have known quivernines $EQ^{(4)}$ and $MQ^{(5)}$. The resolution sequence of the singularity $x_1^4 + x_2^5 + x_3^5 + x_3 x_4^2 = 0$ reads

$$\begin{aligned}
 & (x_1, x_2, x_3, x_4; \delta_1), \quad (x_4, \delta_1; \delta_2), \quad (x_3, \delta_1; \delta_3), \quad (x_1, \delta_2; \delta_4), \quad (\delta_1, \delta_2; \delta_5) \\
 & (\delta_2, \delta_4; \delta_6), \quad (\delta_2, \delta_6; \delta_7), \quad (\delta_2, \delta_7; \delta_8).
 \end{aligned}
 \quad (5.58)$$

The resolved equation is

$$x_1^4 \delta_1 \delta_4^3 \delta_6^2 \delta_7 + x_2^5 \delta_1^2 \delta_2 \delta_3 \delta_5^2 + x_3^5 \delta_1^2 \delta_2 \delta_3 \delta_5^2 + x_3 x_4^2 \delta_2 = 0. \quad (5.59)$$

The sets $\delta_1 = 0$ and $\delta_2 = 0$ are empty, while the other exceptional divisors $\delta_i = 0$ are irreducible. The triple intersection numbers of $S_1 : \delta_3 = 0$, $S_2 : \delta_4 = 0$, $S_3 : \delta_5 = 0$,

$S_4 : \delta_6 = 0$, $S_5 : \delta_7 = 0$ and $S_6 : \delta_8 = 0$ are

(5.60)

The compact divisors S_2 , S_4 and S_5 are rules over genus-two curves. After the flop performed by blowing up two double points on $S_6 \cdot S_5$, the resulting geometry is

(5.61)

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A Crepantly resolvable geometries with smooth divisors and $b_3 = 0$

We already discussed one example of a completely smoothable model with $b_3 = 0$ in the main text, section 5.2. Here we will discuss the remaining theories in this class up until rank $r = 10$.

A.1 AD[D_{2n}, D_{2n}] theories

There is an interesting infinite series of IHS which are ‘fully smooth’ with $b_3 = 0$. In Type IIB, they give rise to the 4d generalised Argyres-Douglas $[D_{2n}, D_{2n}]$ ($n \geq 3$) theories. Note that for $n = 2$, the theory is the rank 1 E_6 SCFT. Consider the type IV($2, 2n-1, 2, 2n-1$) canonical singularity

$$x_1^{2n-1} + x_1 x_2^2 + x_3^{2n-1} + x_3 x_4^2 = 0. \quad (\text{A.1})$$

M-theory on this canonical singularity AD[D_{2n}, D_{2n}] results in a class of 5d SCFTs with the IR gauge theory description:

$$\text{SU}(n)_{\pm 3/2} + (2n+1)\mathbf{F}. \quad (\text{A.2})$$

The UV flavor symmetry is $G_F = \text{SO}(4n+2) \times \text{U}(1)$. Let us first summarize the properties of these singularities:

r	f	$G_F^{5\text{d}}$	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	\mathfrak{t}_2
$n-1$	$2n+2$	$\text{SO}(4n+2) \times \text{U}(1)$	$2n^2 + n + 1$	$2n^2 - n - 1$	$3n+1$	0	0	—

(A.3)

Now we check these properties from both deformation and resolution of \mathbf{X} .

A.1.1 Deformation and magnetic quiver

From the spectrum of the singularity, we can check that it is a Lagrangian theory with

$$f = 2n+2, \quad n_h = \frac{4}{3}n(1+2n^2), \quad n_v = \frac{1}{3}(n-1)(3+8n+8n^2), \quad (\text{A.4})$$

so that we find the virtual dimension of the 4d HB, $n_h - n_v = 3n+1$. This equals the actual dimension as computed from the geometry, $\hat{d}_H = r + f = 3n+1$. The 4d Lagrangian reads [6]:

$$\text{EQ}^{(4)}_{\text{AD}[D_{2n}, D_{2n}]} = \begin{array}{c} \square_1 \\ | \\ \bullet_n \\ | \\ \bullet_{2n-1} \end{array} \begin{array}{c} \bullet_1 \quad \bullet_2 \quad \bullet_3 \quad \cdots \quad \bullet_{2n-2} \quad \bullet_{2n-1} \quad \bullet_n \quad \bullet_1 \end{array} \quad (\text{A.5})$$

where the nodes are special unitary gauge groups. The 5d magnetic quiver is obtained by gauging all the special-unitary and flavor nodes in (A.5), and then modding out a common $\text{U}(1)$:

$$\text{MQ}^{(5)}_{\text{AD}[D_{2n}, D_{2n}]} = \begin{array}{c} \bullet_1 \\ | \\ \bullet_n \\ | \\ \bullet_{2n-1} \end{array} \begin{array}{c} \bullet_1 \quad \bullet_2 \quad \bullet_3 \quad \cdots \quad \bullet_{2n-2} \quad \bullet_{2n-1} \quad \bullet_n \quad \bullet_1 \end{array} \quad (\text{A.6})$$

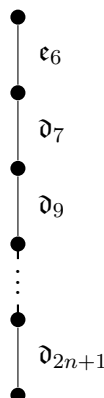


Figure 2. Hasse diagram for the $\text{AD}[D_{2n}, D_{2n}]$ singularity, which has the IR gauge theory description $\text{SU}(n)_{\pm 3/2} + (2n+1)\mathbf{F}$.

The Hasse diagram is shown in figure 2 (it was already computed in [11]). Note that the $\text{AD}[D_4, D_4]$ theory is also equivalent to the rank-1 Seiberg E_6 theory. The 3d mirror $\text{MQ}^{(4)}$ to (A.5) was recently discussed in [88], together with the magnetic quiver for any of the $\text{AD}[D_n, D_m]$ 4d SCFTs.

A.1.2 Resolution and 5d SCFT interpretation

The singularity (A.1) allows for a complete resolution, with the following sequence:

$$\begin{aligned} (x_1, x_2, x_3, x_4; \delta_1) \\ (x_2, x_4, \delta_i; \delta_{i+1}) \quad (i = 1, \dots, n-2). \end{aligned} \tag{A.7}$$

The resulting equation

$$x_1^{2n-1} \prod_{i=1}^{n-2} \delta_i^{2n-2i-2} + x_1 x_2^2 + x_3^{2n-1} \prod_{i=1}^{n-2} \delta_i^{2n-2i-2} + x_3 x_4^2 = 0 \tag{A.8}$$

is completely smooth, and all the exceptional divisors $S_i : \delta_i = 0$ are smooth as well. We can compute the non-vanishing triple intersection numbers among the different divisors, which are represented as

$$\tag{A.9}$$

The intersection numbers are consistent with the 5d IR gauge theory description (A.2) as well.

A.2 Rank 1

There are only three rank $r = 1$ 5d SCFTs in this category: the Seiberg rank 1 E_6 , E_7 and E_8 SCFTs, and we discussed them in [1]. They are described by the following type $\{1, 1\}$

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_1^2 + x_2^3 + x_3^6 + x_4^6$	1	6	E_6	11	5	$N + 6$	0	0	—
$x_1^2 + x_2^4 + x_3^4 + x_4^4$	1	7	E_7	17	10	$N + 7$	0	0	—
$x_1^2 + x_2^3 + x_3^6 + x_4^6$	1	8	E_8	29	21	$N + 8$	0	0	—

Table 7. All smooth rank 1 models with smooth divisors and $b_3 = 0$.

isolated hypersurface singularities:

$$\begin{aligned}
 \mathbf{X}_{E_6} &: x_1^3 + x_2^3 + x_3^3 + x_4^3 = 0 \\
 \mathbf{X}_{E_7} &: x_1^2 + x_2^4 + x_3^4 + x_4^4 = 0 \\
 \mathbf{X}_{E_8} &: x_1^2 + x_2^3 + x_3^6 + x_4^6 = 0.
 \end{aligned} \tag{A.10}$$

The resolution of these singularities consists of a single weighted blow-up in the ambient space:

$$\begin{aligned}
 \mathbf{X}_{E_6} &: (x_1, x_2, x_3, x_4; \delta_1) \\
 \mathbf{X}_{E_7} &: (x_1^{(2)}, x_2^{(1)}, x_3^{(1)}, x_4^{(1)}; \delta_1) \\
 \mathbf{X}_{E_8} &: (x_1^{(3)}, x_2^{(2)}, x_3^{(1)}, x_4^{(1)}; \delta_1).
 \end{aligned} \tag{A.11}$$

The exceptional divisors $S_1 : \delta_1 = 0$ are del Pezzo surfaces dP_6 , dP_7 and dP_8 respectively. We also list the full data of these theories in table 7. The $\text{EQ}^{(4)}$ for the rank 1 E_k theory are $\text{SU}(n)$ quivers in the shape of affine E_k Dynkin diagram, while the $\text{MQ}^{(5)}$ are $\text{U}(n)$ quivers in the same shape. The Hasse diagram is simply a single \mathfrak{e}_k , which corresponds to the minimal nilpotent orbit of E_k .

A.3 Rank 3: $\{4, 1\}\{3, 3, 9, 2\}$

The model is characterized by the following data:

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_1^3 + x_1x_2^3 + x_3^9 + x_3x_4^2$	3	8	$E_7 \times \text{U}(1)$	39	31	11	0	0	—

(A.12)

A.3.1 Deformations and magnetic quiver

From the CB spectrum we can infer the $\text{EQ}^{(4)}$ and $\text{MQ}^{(5)}$. The magnetic quiver for the 5d theory is

$$\text{MQ}^{(5)} = \begin{array}{c} \bullet \quad 1 \\ | \\ \bullet \quad 5 \\ | \\ \bullet \quad 9 \\ | \\ \bullet \quad 3 \quad \bullet \quad 6 \quad \bullet \quad 9 \quad \bullet \quad 7 \quad \bullet \quad 5 \quad \bullet \quad 3 \quad \bullet \quad 1 \end{array} \tag{A.13}$$

and the Hasse diagram is



(A.14)

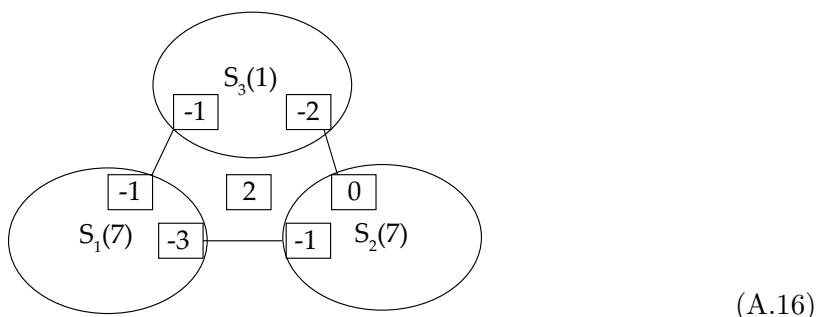
From the Hasse diagram, we can read off the flavor symmetry $G_F = E_7 \times U(1)$.

A.3.2 Resolution and 5d SCFT

The resolution sequence that completely resolves this model is

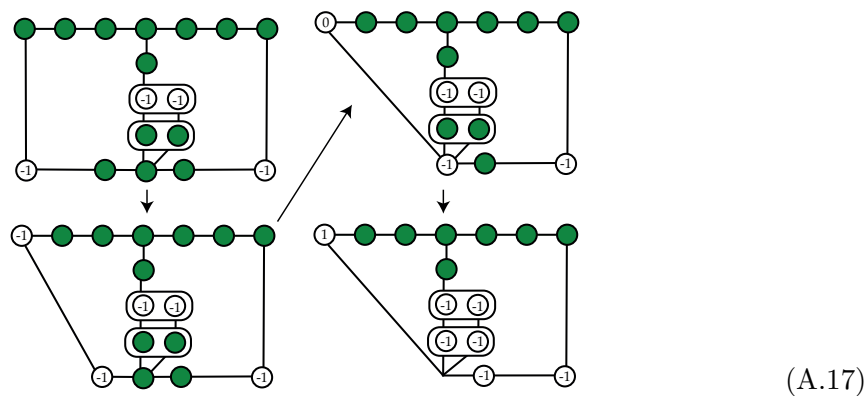
$$(x_1, x_2, x_3, x_4; \delta_1), (x_1^{(2)}, x_2^{(1)}, x_4^{(3)}, \delta_1^{(1)}; \delta_3), (x_1, x_4, \delta_1; \delta_2) \quad (\text{A.15})$$

with the intersection numbers



(A.16)

This theory does not have a 5d IR gauge theory description. Nonetheless, the theory is a descendant of 6d $(E_7, \text{SO}(7))$ conformal matter [79] with the following CFD transitions



(A.17)

The CFD of this 5d SCFT with $G_F = E_7 \times U(1)$ is at the bottom-right corner.

A.4 Rank 3: $\{12, 1\}\{7, 2, 3, 2\}$

The data of this equation is

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_1^7 + x_1x_2^2 + x_1x_3^3 + x_2x_4^2 + x_2x_3^2$	3	8	$E_6 \times \text{SU}(2) \times U(1)$	29	21	11	0	0	—

(A.18)

A.4.1 Deformations and magnetic quiver

The magnetic quiver again follows from the CB spectrum of $\mathcal{T}_{\mathbf{X}}^{4d}$:

$$\text{MQ}^{(5)} = \begin{array}{c} \bullet \quad 1 \\ | \\ \bullet \quad 4 \\ | \\ \bullet \quad 7 \\ | \\ \bullet \quad 1 \quad \bullet \quad 3 \quad \bullet \quad 5 \quad \bullet \quad 7 \quad \bullet \quad 5 \quad \bullet \quad 3 \quad \bullet \quad 1 \end{array} \quad (\text{A.19})$$

and the Hasse diagram is

$$\begin{array}{c} \bullet \\ \swarrow \quad \searrow \\ \bullet \quad \bullet \\ \swarrow \quad \searrow \quad \swarrow \quad \searrow \\ \bullet \quad \bullet \quad \bullet \quad \bullet \\ \swarrow \quad \searrow \\ \bullet \end{array} \quad \begin{array}{l} \mathfrak{e}_7 \quad \mathfrak{e}_6 \\ \mathfrak{e}_6 \quad \mathfrak{a}_1 \quad \mathfrak{e}_7 \quad \mathfrak{a}_7 \\ \mathfrak{a}_1 \quad \mathfrak{e}_6 \end{array} \quad (\text{A.20})$$

From the Hasse diagram, we can read off the flavor symmetry is $G_F = E_6 \times \text{SU}(2) \times \text{U}(1)$, which is different from the subgraph of the balanced nodes in the magnetic quiver.

A.4.2 Resolution and 5d SCFT

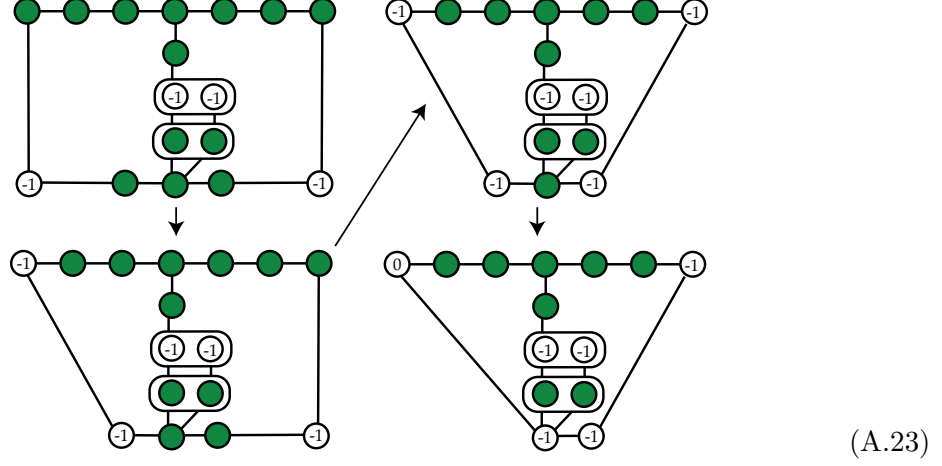
The resolution sequence is

$$(x_1, x_2, x_3, x_4; \delta_1), \quad (x_2^{(2)}, x_3^{(1)}, x_4^{(1)}, \delta_1^{(1)}; \delta_3), \quad (x_2, \delta_1; \delta_2). \quad (\text{A.21})$$

The non-vanishing triple intersection numbers are

$$\begin{array}{c} \text{---} \bigcirc \text{---} \\ \text{---} \bigcirc \text{---} \bigcirc \text{---} \\ \text{---} \bigcirc \text{---} \end{array} \quad \begin{array}{l} S_3(2) \\ S_1(7) \quad S_2(6) \end{array} \quad (\text{A.22})$$

This theory does not have a 5d IR gauge theory description. Nonetheless, the theory is a descendant of 6d ($E_7, SO(7)$) conformal matter as well, with the following CFD transitions



The theory on the bottom-right has the correct r_F and G_F . There is no valid BG-CFD embedding.

A.5 Rank 6: $\{18, 1\}\{2, 3, 2, 5\}$

The model with integral spectrum at this rank is:

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_2x_4^5 + x_3x_4^4 + x_1x_2^3 + x_2x_3^2 + x_1^2x_3$	6	6	$SU(5) \times U(1)^2$	36	30	12	0	0	—

(A.24)

A.5.1 Deformations and 4d SCFT

In this case the CB spectrum of \mathcal{T}_X^{4d} is

$$\Delta = \{2^4, 3^5, 4^4, 5^4, 6^3, 7^3, 8^2, 9^2, 10, 11, 13\}, \quad (\text{A.25})$$

which contains operators with $\Delta = 11$ and $\Delta = 13$, but there is no operator with $\Delta = 12$. The CB spectrum cannot match any set of Casimir operators of Lie groups, and there is no Lagrangian description of \mathcal{T}_X^{4d} . Again there is a conjectural generalised quiver that matches all the data of this model (flavor rank and CB spectrum)

$$[1] - SU(5) - SU(9) - D_2(SO(28)) - SO(12) \times Spin(6) - D_2(SO(16)) - [4] \quad (\text{A.26})$$

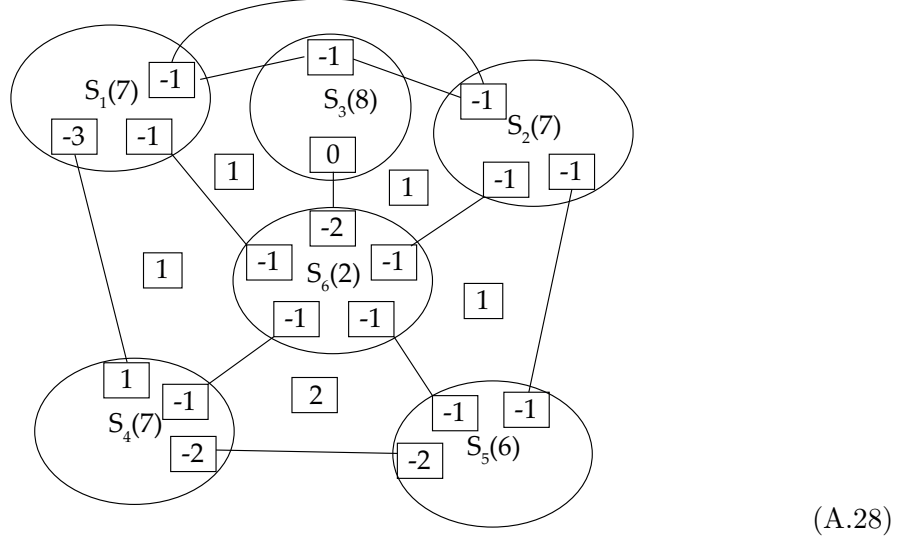
Note that $Spin(6)$ is gauged as a subgroup of the $SO(18) \times U(1)$ flavor of the $D_2(SO(16))$. Similarly, we gauge an $SO(12) \times SU(9)$ subgroup of the $SO(30) \times U(1)$ flavor symmetry of $D_2(SO(28))$.

A.5.2 Resolution and 5d SCFT

Resolution sequence that fully resolves this model is

$$\begin{aligned} (x_1, x_2, x_3, x_4; \delta_1), & \quad (x_1, x_2, x_3, \delta_1; \delta_4), & \quad (x_1^{(1)}, x_3^{(2)}, \delta_1^{(1)}, \delta_4^{(1)}; \delta_6) \\ (x_3, \delta_1; \delta_2), & \quad (x_1, \delta_2; \delta_3), & \quad (x_3, \delta_4; \delta_5). \end{aligned} \quad (\text{A.27})$$

The intersection relations among the compact surfaces $S_i : \delta_i = 0$ are



Note that $S_1 \cdot S_2 \cdot S_6 = S_1 \cdot S_3 \cdot S_6 = S_2 \cdot S_3 \cdot S_6 = 1$, the curve $S_1 \cdot S_2 = S_2 \cdot S_3 = S_1 \cdot S_3$ coincides, $S_1 \cdot S_2 \cdot S_3 = -1$. Since the curve $S_3 \cdot S_1$ is a section curve on S_3 and a ruling curve on S_1 , the geometry does not have a consistent ruling structure.

For the non-abelian part of flavor symmetry, it comes from the extra (-2) -curves on S_6 , which has the topology of gdP_7 . The curves on S_6 shown in (A.28) are written in the Picard group generators h, e_1, \dots, e_7 :

$$\begin{aligned}
 S_6 \cdot S_3 &= e_6 - e_7 \\
 S_6 \cdot S_2 &= e_7 \\
 S_6 \cdot S_5 &= h - e_6 - e_7 \\
 S_6 \cdot S_4 &= 2h - e_1 - e_2 - e_3 - e_4 - e_5 \\
 S_6 \cdot S_1 &= 3h - e_1 - e_2 - e_3 - e_4 - e_5 - 2e_6 - e_7.
 \end{aligned}
 \tag{A.29}$$

The maximal set of extra (-2) -curves with non-negative intersection numbers with the above curves are

$$\begin{aligned}
 C_1 &= e_1 - e_2 \\
 C_2 &= e_2 - e_3 \\
 C_3 &= e_3 - e_4 \\
 C_4 &= e_4 - e_5,
 \end{aligned}
 \tag{A.30}$$

which generate an $\text{SU}(5)$ flavor symmetry factor. Hence we have at least

$$G_F = \text{SU}(5) \times \text{U}(1)^2. \tag{A.31}$$

A.6 Rank 6: $\{12, 2\}\{13, 3, 2, 3\}$

The model has the following basic characteristics

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	\mathfrak{f}
$x_1^{13} + x_1x_2^3 + x_1x_3^2 + x_2x_3x_4 + x_2x_4^3$	6	9	$E_7 \times \text{U}(1)^2$	57	48	15	0	0	—

(A.32)

A.6.1 Deformations and magnetic quiver

From the CB spectrum we can identify the magnetic quiver as

$$\text{MQ}^{(5)} = \begin{array}{c} \bullet \text{ 1} \\ | \\ \bullet \text{ 7} \\ | \\ \bullet \text{ 13} \\ | \\ \bullet \text{ 1} \text{---} \bullet \text{ 5} \text{---} \bullet \text{ 9} \text{---} \bullet \text{ 10} \text{---} \bullet \text{ 7} \text{---} \bullet \text{ 4} \text{---} \bullet \text{ 1} \end{array} \quad (\text{A.33})$$

From the balanced nodes and the Hasse diagram, we can read off the non-abelian part of flavor symmetry $G_{F,nA} = E_7$. Hence the full UV flavor symmetry is $G_F = E_7 \times \text{U}(1)^2$.

A.6.2 Resolution and 5d SCFT

The following sequence of resolutions smoothes this model

$$\begin{aligned} (x_1, x_2, x_3, x_4; \delta_1), \quad & (x_2^{(1)}, x_3^{(2)}, x_4^{(1)}, \delta_1^{(1)}; \delta_4), \quad (x_2^{(2)}, x_3^{(3)}, x_4^{(1)}, \delta_4^{(1)}; \delta_6) \\ (x_3, \delta_1; \delta_2), \quad & (x_2, \delta_2; \delta_3), \quad (x_2, x_3, \delta_4; \delta_5). \end{aligned} \quad (\text{A.34})$$

The triple intersection numbers of $S_i : \delta_i = 0$ are given in the following diagram:

$$(\text{A.35})$$

For the non-Abelian part of flavor symmetry G_F , it comes from the extra (-2) -curves on S_6 . The topology of S_6 can be chosen as gdP_8 of type $\mathbf{E}_7 + \mathbf{A}_1$ [89], and the (-2) -curves arranged in an E_7 Dynkin diagram exactly gives rise to $G_{F,nA} = E_7$.

A.7 Rank 9: $\{4, 1\}\{4, 2, 4, 4\}$

We consider the type $\{4, 1\}\{4, 2, 4, 4\}$ singularity, with the data

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_1^4 + x_2^2 x_1 + x_3^4 + x_3 x_4^4$	9	5	$\text{SU}(2)^4 \times \text{U}(1)$	35	30	14	0	0	—

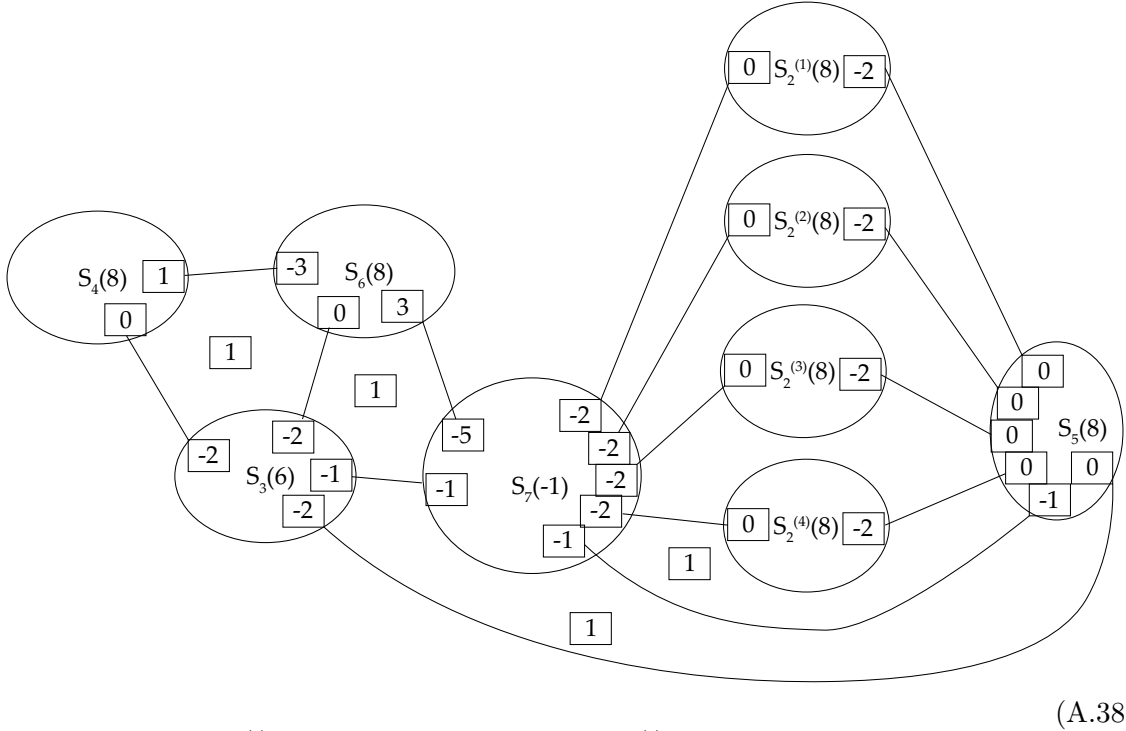
$$(\text{A.36})$$

In this case the CB spectrum of $\mathcal{T}_{\mathbf{X}}^{4d}$ contains operators with $\Delta = 13$ and $\Delta = 16$, but there is no operator with $\Delta = 14$ or 15. The CB spectrum cannot match any set of Casimir operators of Lie groups, and there is no Lagrangian description of $\mathcal{T}_{\mathbf{X}}^{4d}$. It would be interesting to identify similarly to the other models a generalised quiver.

The resolution sequence for this model is

$$\begin{aligned}
 & (x_1, x_2, x_3, x_4; \delta_1), \quad (x_1, x_2, x_3, \delta_1; \delta_4), \quad (x_2, \delta_1; \delta_2), \quad (x_1, \delta_1; \delta_3) \\
 & (\delta_1, \delta_2; \delta_5), \quad (\delta_2, \delta_4; \delta_6), \quad (\delta_2, \delta_6; \delta_7)
 \end{aligned} \tag{A.37}$$

The equation $\delta_1 = 0$ is empty, and the equation $\delta_2 = 0$ has four irreducible components, which correspond to compact divisors $S_2^{(1)}, \dots, S_2^{(4)}$. The triple intersection numbers of $S_i : \delta_i = 0$ are given in the following diagram:



Note that $S_5 \cdot S_7 \cdot S_2^{(i)} = 1$ for $i = 1, 2, 3, 4$. Since $S_2^{(i)}$ ($i = 1, 2, 3, 4$) are Hirzebruch surface \mathbb{F}_2 , and S_5 is \mathbb{F}_1 , $S_2^{(i)} \cdot S_5$ is a section curve on $S_2^{(i)} \cdot S_5$ but a ruling curve on S_5 . Hence the geometry does not have a ruling structure and IR gauge theory description.

For the non-abelian part of flavor symmetry $G_{F,nA}$, it comes from additional (-2) -curves on S_7 . Note that the curves on S_7 appeared in (A.38) can be written in terms of

Picard group generators h, e_1, \dots, e_{10} as

$$\begin{aligned}
 S_7 \cdot S_6 &= 2h - e_1 - e_2 - e_3 - e_4 - e_5 - e_6 - e_7 - e_8 - e_9 \\
 S_7 \cdot S_3 &= h - e_9 - e_{10} \\
 S_7 \cdot S_5 &= e_{10} \\
 S_7 \cdot S_2^{(1)} &= h - e_1 - e_2 - e_{10} \\
 S_7 \cdot S_2^{(2)} &= h - e_3 - e_4 - e_{10} \\
 S_7 \cdot S_2^{(3)} &= h - e_5 - e_6 - e_{10} \\
 S_7 \cdot S_2^{(4)} &= h - e_7 - e_8 - e_{10}
 \end{aligned} \tag{A.39}$$

The maximal set of additional (-2) -curves on S_7 with non-negative intersection numbers with the above curves are

$$\begin{aligned}
 C_1 &= e_1 - e_2 \\
 C_2 &= e_3 - e_4 \\
 C_3 &= e_5 - e_6 \\
 C_4 &= e_7 - e_8.
 \end{aligned} \tag{A.40}$$

These curves generate non-abelian flavor symmetry $G_{F,nA} = \text{SU}(2)^4$, hence we have (at least)

$$G_F = \text{SU}(2)^4 \times \text{U}(1). \tag{A.41}$$

A.8 Rank 10: $\{10, 1\}\{5, 2, 2, 6\}$

The model has the following data

F	r	f	G_F^{5d}	d_H	\hat{r}	\hat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2
$x_2x_4^6 + x_3^2x_4 + x_1^5 + x_2^2x_3$	10	4	SU(5)	50	46	14	0	0	—

(A.42)

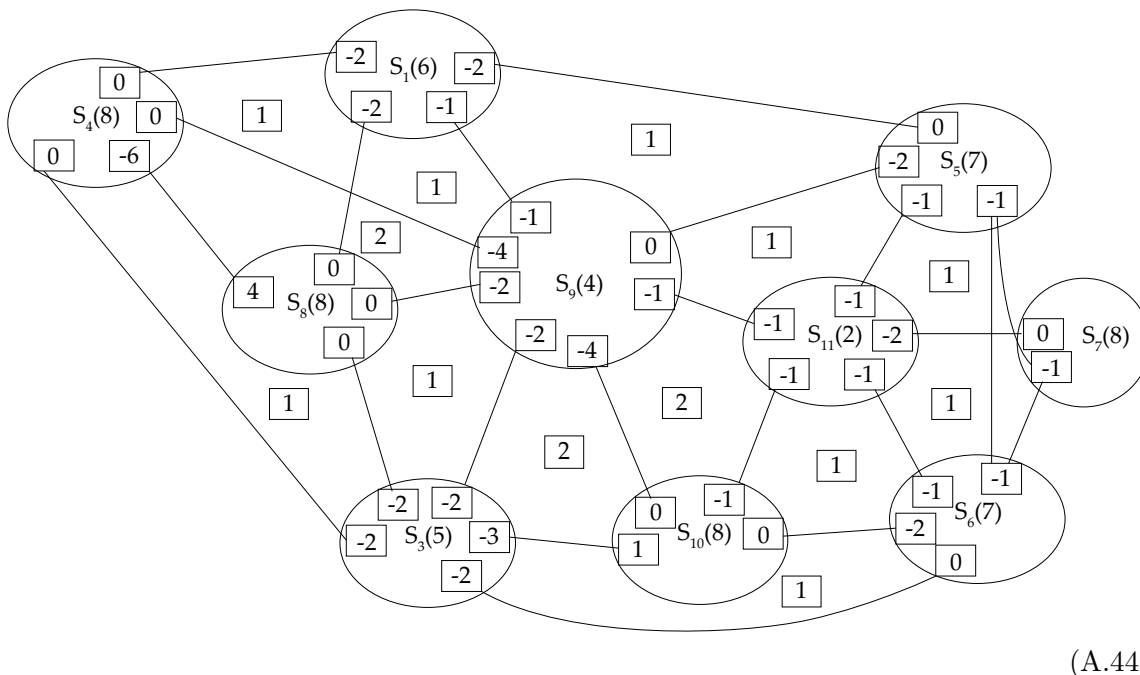
In this case the CB spectrum of \mathcal{T}_X^{4d} contains operators with $\Delta = 22$ and $\Delta = 25$, but there is no operator with $\Delta = 23$ or 24. The CB spectrum cannot match any set of Casimir operators of Lie groups, and there is no Lagrangian description of \mathcal{T}_X^{4d} .

Resolution sequence is

$$\begin{aligned}
 &(x_1, x_2, x_3, x_4; \delta_1), \quad (x_1^{(1)}, x_2^{(1)}, x_3^{(2)}, \delta_1^{(1)}; \delta_5), \quad (x_2^{(1)}, x_3^{(2)}, \delta_1^{(1)}, \delta_5^{(1)}; \delta_9) \\
 &(x_2^{(1)}, x_3^{(2)}, \delta_5^{(1)}, \delta_9^{(1)}; \delta_{11}), \quad (x_3, \delta_1; \delta_2), \quad (x_3, \delta_2; \delta_3), \quad (x_2, \delta_2; \delta_4), \quad (x_3, \delta_5; \delta_6) \\
 &(x_2, \delta_6; \delta_7), \quad (\delta_2, \delta_4; \delta_8).
 \end{aligned} \tag{A.43}$$

The equation $\delta_9 = 0$ is empty, and the triple intersection numbers among the remaining

10 compact divisors $S_i : \delta_i = 0$ are given in the following diagram:



Note that $S_1 \cdot S_4 \cdot S_8 = S_1 \cdot S_8 \cdot S_9 = 1$, $S_4 \cdot S_8 \cdot S_9 = 2$, $S_{11} \cdot S_5 \cdot S_6 = S_{11} \cdot S_5 \cdot S_7 = S_{11} \cdot S_6 \cdot S_7 = 1$, $S_5 \cdot S_6 \cdot S_7 = -1$. The curve $S_5 \cdot S_6 = S_5 \cdot S_7 = S_6 \cdot S_7$ coincides. Since $S_6 \cdot S_7$ is a section curve on S_7 but a ruling curve on S_5 and S_6 , the whole geometry does not have a ruling structure.

For the flavor symmetry G_F , note that the topology of S_{11} is exactly the same as the topology of S_6 in the rank-6 model (A.24). Following the same argument based on extra (-2) -curves on S_{11} , we get

$$G_F = \text{SU}(5). \quad (\text{A.45})$$

B Tables: smoothable models

Type	(a, b, c, d)	F	r	f	d_H	\hat{r}	$\widehat{d_H}$	$\Delta\mathcal{A}_r$	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾).
$\{4, 1\}$	$\{5, 2, 5, 2\}$	$x_1^5 + x_2^2 x_1 + x_3^5 + x_3 x_4^2$	2	8	22	14	10	0	0	0	
$\{4, 1\}$	$\{3, 3, 9, 2\}$	$x_3^9 + x_4^2 x_3 + x_1^3 + x_1 x_2^3$	3	8	39	31	11	0	0	0	
$\{4, 1\}$	$\{7, 2, 7, 2\}$	$x_1^7 + x_2^2 x_1 + x_3^7 + x_3 x_4^2$	3	10	37	27	13	0	0	0	
$\{12, 1\}$	$\{7, 2, 3, 2\}$	$x_1^7 + x_3^3 x_1 + x_2^2 x_1 + x_2 x_3^2 + x_2 x_4^2$	3	8	29	21	11	0	0	0	
$\{4, 1\}$	$\{9, 2, 9, 2\}$	$x_1^9 + x_2^2 x_1 + x_3^9 + x_3 x_4^2$	4	12	56	44	16	0	0	0	
$\{4, 1\}$	$\{3, 3, 3, 3\}$	$x_1^3 + x_2^3 x_1 + x_3^3 + x_3 x_4^3$	5	7	28	21	12	0	0	0	$2^4 3^5 4^4 5^3 6^2 7^2, 9$
$\{4, 1\}$	$\{11, 2, 11, 2\}$	$x_1^{11} + x_2^2 x_1 + x_3^{11} + x_3 x_4^2$	5	14	79	65	19	0	0	0	
$\{4, 1\}$	$\{13, 2, 13, 2\}$	$x_1^{13} + x_2^2 x_1 + x_3^{13} + x_3 x_4^2$	6	16	106	90	22	0	0	0	

Table 8. Rank $r > 1$ Model with smooth divisors and $b_3 = 0$. The red colored models have a generalised $D_p G$ quiver representation.

Type	(a, b, c, d)	F	r	f	d_H	\widehat{d}_H	ΔA_r	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{12, 2\}$	$\{13, 3, 2, 3\}$	$x_1^{13} + x_2^3 x_1 + x_3^2 x_1 + x_2 x_3^3 + x_2 x_3 x_4$	6	9	57	48	0	0	0	
$\{18, 1\}$	$\{2, 3, 2, 5\}$	$x_2 x_4^5 + x_3 x_4^4 + x_1 x_2^3 + x_2 x_3^2 + x_1^2 x_3$	6	6	36	30	0	0	0	$2^4, 3^5, 4^4, 5^4, 6^3, 7^3, 8^2, 9^2, 10, 11, 13$
$\{4, 1\}$	$\{15, 2, 15, 2\}$	$x_1^{15} + x_2^2 x_1 + x_3^{15} + x_3 x_4^2$	7	18	137	119	0	0	0	
$\{4, 1\}$	$\{17, 2, 17, 2\}$	$x_1^{17} + x_2^2 x_1 + x_3^{17} + x_3 x_4^2$	8	20	172	152	0	0	0	
$\{4, 1\}$	$\{4, 2, 4, 4\}$	$x_1^4 + x_2^2 x_1 + x_3^4 + x_3 x_4^4$	9	5	35	30	0	0	0	$2^4, 3^2, 4^5, 5^3, 6^3, 7^2, 8^3, 9^2, 10^2, 12^2, 13, 16$
$\{4, 1\}$	$\{19, 2, 19, 2\}$	$x_1^{19} + x_2^2 x_1 + x_3^{19} + x_3 x_4^2$	9	22	211	189	0	0	0	
$\{4, 1\}$	$\{21, 2, 21, 2\}$	$x_1^{21} + x_2^2 x_1 + x_3^{21} + x_3 x_4^2$	10	24	254	230	0	0	0	
$\{10, 1\}$	$\{5, 2, 2, 6\}$	$x_2 x_4^6 + x_3^2 x_4 + x_1^5 + x_2^2 x_3$	10	4	50	46	0	0	0	$2^3, 3^4, 4^4, 5^3, 6^3, 7^3, 8^3, 9^3, 10^3, 11^2, 12^2, 13^2, 14^2, 15^2, 16, 17, 18, 19, 20, 22, 25$

Table 9. Rank $r > 1$ Model with smooth divisors and $b_3 = 0$ (cont.); orange shows models which do not have a $D_p^b(G)$ trinion representation in 4d.

Type	(a, b, c, d)	F	r	f	d_H	\widehat{r}	$\widehat{d_H}$	$\Delta\mathcal{A}_r$	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{1, 1\}$	$\{2, 3, 6, 12\}$	$x_4^{12} + x_3^6 + x_2^3 + x_1^2$	2	8	59	51	10	-1	2	\mathbb{Z}_2^2	
$\{1, 1\}$	$\{2, 4, 4, 8\}$	$x_4^8 + x_2^4 + x_3^4 + x_1^2$	2	7	35	28	9	-1	2	\mathbb{Z}_2^2	
$\{1, 1\}$	$\{2, 5, 5, 5\}$	$x_2^5 + x_3^5 + x_4^5 + x_1^2$	2	0	32	32	2	-6	12	\mathbb{Z}_2^{12}	
$\{1, 1\}$	$\{3, 3, 3, 6\}$	$x_4^6 + x_1^3 + x_2^3 + x_3^3$	2	6	23	17	8	-1	2	\mathbb{Z}_2^2	
$\{1, 1\}$	$\{2, 3, 6, 18\}$	$x_4^{18} + x_3^6 + x_2^3 + x_1^2$	3	8	89	81	11	-2	4	\mathbb{Z}_3^2	
$\{1, 1\}$	$\{2, 3, 9, 9\}$	$x_3^9 + x_4^9 + x_2^3 + x_1^2$	3	0	64	64	3	-7	14	\mathbb{Z}_2^{14}	
$\{1, 1\}$	$\{2, 4, 4, 12\}$	$x_4^{12} + x_2^4 + x_3^4 + x_1^2$	3	7	53	46	10	-2	4	\mathbb{Z}_3^2	
$\{1, 1\}$	$\{3, 3, 3, 9\}$	$x_4^9 + x_1^3 + x_2^3 + x_3^3$	3	6	35	29	9	-2	4	\mathbb{Z}_3^2	
$\{1, 1\}$	$\{3, 4, 4, 4\}$	$x_2^4 + x_3^4 + x_4^4 + x_1^2$	3	0	27	27	3	-6	12	\mathbb{Z}_3^6	$2^6, 3^7, 5^3, 6^6, 8, 9^3, 12$
$\{2, 1\}$	$\{2, 7, 7, 3\}$	$x_2^7 + x_3^7 + x_3x_4^3 + x_1^2$	3	0	45	45	3	-6	12	\mathbb{Z}_2^{12}	
$\{2, 1\}$	$\{4, 4, 4, 2\}$	$x_1^4 + x_2^4 + x_3^4 + x_3x_4^2$	3	3	24	21	6	-3	6	\mathbb{Z}_2^6	

Table 10. Rank $r = 2, 3$ Models with smooth divisors and $b_3 > 0$. Blue indicates models with trinion representations in 4d.

Type	(a, b, c, d)	F	r	f	d_H	\widehat{r}	\widehat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{1, 1\}$	$\{2, 3, 6, 24\}$	$x_4^{24} + x_3^6 + x_2^3 + x_1^2$	4	8	119	111	12	-3	6	\mathbb{Z}_4^2	
$\{1, 1\}$	$\{2, 4, 4, 16\}$	$x_4^{16} + x_2^4 + x_3^4 + x_1^2$	4	7	71	64	11	-3	6	\mathbb{Z}_4^2	
$\{1, 1\}$	$\{2, 4, 6, 6\}$	$x_3^6 + x_4^6 + x_2^4 + x_1^2$	4	5	40	35	9	-2	4	\mathbb{Z}_2^4	
$\{1, 1\}$	$\{3, 3, 3, 12\}$	$x_4^{12} + x_1^3 + x_2^3 + x_3^3$	4	6	47	41	10	-3	6	\mathbb{Z}_4^2	
$\{2, 1\}$	$\{3, 6, 6, 2\}$	$x_2^6 + x_3^6 + x_1^3 + x_3 x_4^2$	4	2	36	34	6	-4	8	\mathbb{Z}_2^8	
$\{7, 1\}$	$\{5, 5, 2, 2\}$	$x_1^5 + x_2^5 + x_2 x_3^2 + x_3 x_4^2$	4	4	30	26	8	-2	4	\mathbb{Z}_2^4	
$\{1, 1\}$	$\{2, 3, 6, 30\}$	$x_4^{30} + x_3^6 + x_2^3 + x_1^2$	5	8	149	141	13	-4	8	\mathbb{Z}_5^2	
$\{1, 1\}$	$\{2, 3, 10, 10\}$	$x_3^{10} + x_4^{10} + x_2^3 + x_1^2$	5	0	81	81	5	-8	16	\mathbb{Z}_3^8	$2^7, 3^8, 5^6, 6^9, 8^5, 9^8, 11^4, 12^7, 14^3, 15^6, 17^2, 18^5, 20, 21^4, 24^3, 27^2, 30$
$\{1, 1\}$	$\{2, 4, 4, 20\}$	$x_4^{20} + x_2^4 + x_3^4 + x_1^2$	5	7	89	82	12	-4	8	\mathbb{Z}_5^2	
$\{1, 1\}$	$\{3, 3, 3, 15\}$	$x_4^{15} + x_1^3 + x_2^3 + x_3^3$	5	6	59	53	11	-4	8	\mathbb{Z}_5^2	

Table 11. Rank $r = 4, 5$ Models with smooth divisors and $b_3 > 0$. Blue indicates models that have a $D_p^b(G)$ trinion representation.

Type	(a, b, c, d)	F	r	f	d_H	\widehat{d}_H	ΔA_r	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{1, 1\}$	$\{2, 3, 6, 36\}$	$x_4^{36} + x_3^6 + x_2^3 + x_1^2$	6	8	179	171	14	-5	10 \mathbb{Z}_6^2	
$\{1, 1\}$	$\{2, 4, 4, 24\}$	$x_4^{24} + x_2^4 + x_3^4 + x_1^2$	6	7	107	100	13	-5	10 \mathbb{Z}_6^2	
$\{1, 1\}$	$\{2, 5, 6, 6\}$	$x_3^6 + x_4^6 + x_2^5 + x_1^2$	6	0	50	50	6	-8	16 \mathbb{Z}_5^4	$2^3, 3^4, 4^5, 5^4, 7^2, 8^3, 9^4, 10^5, 12, 13^2, 14^3, 15^6, 18, 19^2, 20^3, 24, 25^2, 30$
$\{1, 1\}$	$\{3, 3, 3, 18\}$	$x_4^{18} + x_1^3 + x_2^3 + x_3^3$	6	6	71	65	12	-5	10 \mathbb{Z}_6^2	
$\{2, 1\}$	$\{2, 6, 6, 4\}$	$x_2^6 + x_3^6 + x_3 x_4^4 + x_1^2$	6	1	48	47	7	-6	12 \mathbb{Z}_4	$2^4, 3^5, 4^5, 5, 6^3, 7^4, 8^5, 9, 10^2, 11^3, 12^4, 14, 15^2, 16^3, 19, 20^2, 24$
$\{2, 1\}$	$\{2, 8, 8, 3\}$	$x_2^8 + x_3^8 + x_3 x_4^3 + x_1^2$	6	1	60	59	7	-6	12 \mathbb{Z}_3^6	$2^6, 3^7, 4, 5^5, 6^7, 7, 8^4, 9^6, 10, 11^3, 12^5, 14^2, 15^4, 17, 18^3, 21^2, 24$
$\{2, 1\}$	$\{4, 5, 5, 2\}$	$x_2^5 + x_3^5 + x_1^4 + x_3 x_4^2$	6	0	36	36	6	-6	12 \mathbb{Z}_4^4	$2^4, 3^5, 4^5, 6^2, 7^4, 8^5, 10, 11^2, 12^4, 15, 16^2, 20$
$\{4, 1\}$	$\{5, 2, 5, 3\}$	$x_1^5 + x_2^2 x_1 + x_3^5 + x_3 x_4^3$	6	2	34	32	8	-4	8 \mathbb{Z}_3^4	$2^5, 3^6, 4, 5^4, 6^5, 7, 8^2, 9^4, 11, 12^2, 15$
$\{7, 1\}$	$\{8, 4, 2, 2\}$	$x_1^8 + x_2^4 + x_2 x_3^2 + x_3 x_4^2$	6	3	40	37	9	-2	4 \mathbb{Z}_2^4	$\text{Sp}(2) \quad \text{Sp}(8)$
$\{1, 1\}$	$\{2, 3, 6, 42\}$	$x_4^{42} + x_3^6 + x_2^3 + x_1^2$	7	8	209	201	15	-6	12 \mathbb{Z}_7^2	
$\{1, 1\}$	$\{2, 4, 4, 28\}$	$x_4^{28} + x_2^4 + x_3^4 + x_1^2$	7	7	125	118	14	-6	12 \mathbb{Z}_7^2	
$\{1, 1\}$	$\{3, 3, 3, 21\}$	$x_4^{21} + x_1^3 + x_2^3 + x_3^3$	7	6	83	77	13	-6	12 \mathbb{Z}_7^2	

Table 12. Rank $r = 6, 7$ Models with smooth divisors and $b_3 > 0$. Orange labels models with integral spectrum, but no known trinion representation in 4d.

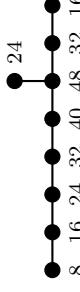
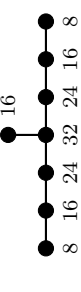
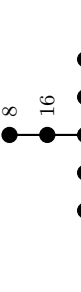
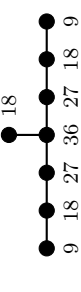
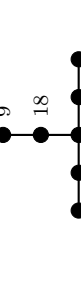
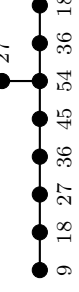
Type	(a, b, c, d)	F	r	f	d_H	\hat{r}	$\widehat{d_H}$	ΔA_r	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{1, 1\}$	$\{2, 3, 6, 48\}$	$x_4^{48} + x_3^6 + x_2^3 + x_1^2$	8	8	239	231	16	-7	14	\mathbb{Z}_8^2	
$\{1, 1\}$	$\{2, 3, 9, 12\}$	$x_4^{12} + x_3^9 + x_2^3 + x_1^2$	8	4	90	86	12	-2	4	\mathbb{Z}_3^2	$2^5, 3^4, 4^4, 5^5, 6^5, 7^3, 8^5, 9^5, 10^3, 11^3, 12^5, 13^3, 14^3, 15^3, 16^3, 17^3, 18^3, 19, 20^3, 21^3, 22, 23, 24^3, 25, 26, 27, 28, 30, 32, 33, 36$
$\{1, 1\}$	$\{2, 4, 4, 32\}$	$x_4^{32} + x_3^4 + x_2^4 + x_1^2$	8	7	143	136	15	-7	14	\mathbb{Z}_8^2	
$\{1, 1\}$	$\{2, 4, 6, 8\}$	$x_4^8 + x_3^6 + x_2^4 + x_1^2$	8	3	54	51	11	-2	4	\mathbb{Z}_3^2	$2^5, 3^4, 4^3, 5^4, 6^5, 7^2, 8^4, 9^4, 10^2, 11^2, 12^4, 13, 14^2, 15^2, 16, 17, 18^2, 20, 21, 24$
$\{1, 1\}$	$\{3, 3, 3, 24\}$	$x_4^{24} + x_3^3 + x_2^3 + x_3^3$	8	6	95	89	14	-7	14	\mathbb{Z}_8^2	
$\{2, 1\}$	$\{3, 6, 3, 3\}$	$x_2^6 + x_1^3 + x_3^3 + x_3 x_4^3$	8	2	36	34	10	-2	4	\mathbb{Z}_3^2	$2^5, 3^4, 4^2, 5^3, 6^5, 7, 8^3, 9^3, 10, 11, 12^3, 14, 15, 18$
$\{2, 1\}$	$\{3, 6, 9, 2\}$	$x_3^9 + x_4^2 x_3 + x_2^6 + x_1^3$	8	6	53	47	14	-1	2	\mathbb{Z}_2^2	$2^5, 3^4, 4^7, 5^3, 6^5, 7^4, 8^3, 9^3, 10^4, 11, 12^3, 13, 14, 15, 16, 18$
$\{1, 1\}$	$\{2, 4, 4, 36\}$	$x_4^{36} + x_2^4 + x_3^4 + x_1^2$	9	7	161	154	16	-8	16	\mathbb{Z}_9^2	
$\{1, 1\}$	$\{3, 3, 3, 27\}$	$x_4^{27} + x_1^3 + x_2^3 + x_3^3$	9	6	107	101	15	-8	16	\mathbb{Z}_9^2	
$\{2, 1\}$	$\{2, 3, 6, 45\}$	$x_3 x_4^{45} + x_3^6 + x_2^3 + x_1^2$	9	8	269	261	17	-8	16	\mathbb{Z}_9^2	
$\{2, 1\}$	$\{4, 4, 8, 2\}$	$x_3^8 + x_4^2 x_3 + x_1^4 + x_2^4$	9	3	42	39	12	-3	6	\mathbb{Z}_6^2	$2^8, 4^8, 5^3, 6^6, 8^6, 9, 10^3, 12^3, 14, 16$
$\{12, 1\}$	$\{7, 2, 3, 3\}$	$x_1^7 + x_3^3 x_1 + x_2^2 x_1 + x_2 x_4^3 + x_2 x_3^2$	9	3	44	41	12	-2	Unknown	\mathbb{Z}_3^2	$2^4, 3^5, 4^2, 5^4, 6^4, 7^2, 8^3, 9^4, 10, 11^2, 12^3, 13, 14, 15^2, 17, 18, 21$

Table 13. Rank $r = 8, 9$ Models with smooth divisors and $b_3 > 0$.

Type	(a, b, c, d)	F	r	f	d_H	\hat{r}	$\widehat{d_H}$	ΔA_r	b_3	t_2	Magnetic Quiver (MQ ⁽⁵⁾)
$\{1, 1\}$	$\{2, 4, 4, 40\}$	$x_4^{40} + x_2^4 + x_3^4 + x_1^2$	10	7	179	172	17	-9	18	\mathbb{Z}_{10}^2	
$\{1, 1\}$	$\{3, 3, 3, 30\}$	$x_4^{30} + x_1^3 + x_2^3 + x_3^3$	10	6	119	113	16	-9	18	\mathbb{Z}_{10}^2	
$\{2, 1\}$	$\{2, 3, 6, 50\}$	$x_3 x_4^{50} + x_3^6 + x_2^3 + x_1^2$	10	8	299	291	18	-9	18	\mathbb{Z}_{10}^2	
$\{2, 1\}$	$\{2, 5, 5, 6\}$	$x_3 x_4^6 + x_2^5 + x_3^5 + x_1^2$	10	0	52	52	10	-2	4	\mathbb{Z}_2^4	$2^7, 4^5, 6^7, 8^6, 10^4, 12^6, 14^4, 16^2, 18^4, 20^2, 22, 24^2, 26, 30$
$\{2, 1\}$	$\{3, 6, 12, 2\}$	$x_3^{12} + x_4^2 x_3 + x_2^6 + x_1^3$	10	2	66	64	12	-4	8	\mathbb{Z}_2^8	$2^9, 4^9, 5^2, 6^8, 8^8, 9, 10^6, 12^6, 13, 14^4, 16^4, 18^2, 20^2, 22, 24$

Table 14. Rank $r = 10$ Models with smooth divisors and $b_3 > 0$.

Type	(a, b, c, d)	F	Divs	r	f	d_H	\hat{r}	\widehat{d}_H	ΔA_r	b_3	t_2	CB Spectrum
$\{1, 1\}$	$\{2, 3, 6, 6\}$	$x_3^6 + x_4^6 + x_3^2 + x_1^2$	Smt	1	8	29	21	9	0	0	0	$\{2^8, 3^6, 4^4, 5^2, 6\}$
$\{1, 1\}$	$\{2, 3, 6, 7\}$	$x_4^7 + x_3^6 + x_3^2 + x_1^2$	Sgl	1	0	30	30	1	-1	2	\mathbb{Z}_7^2	$\{\frac{7}{6}, \frac{4}{3}, \frac{3}{2}, \frac{2}{2}, \frac{3}{3}, \frac{5}{6}, \frac{11}{2}, \frac{7}{3}, \frac{7}{3}, \frac{5}{3}, \frac{5}{3}, \frac{8}{3}, \frac{17}{6}, 3, \frac{7}{2}, \frac{11}{3}, \frac{23}{4}, \frac{14}{3}, \frac{14}{3}, \frac{29}{5}, \frac{35}{6}, 6, 7\}$
$\{1, 1\}$	$\{2, 4, 4, 4\}$	$x_2^4 + x_3^4 + x_4^4 + x_1^2$	Smt	1	7	17	10	8	0	0	0	$\{2^6, 3^3, 4\}$
$\{1, 1\}$	$\{2, 4, 4, 5\}$	$x_4^5 + x_2^4 + x_3^4 + x_1^2$	Sgl	1	0	18	18	1	-1	2	\mathbb{Z}_5^2	$\{\frac{5}{4}, \frac{5}{4}, \frac{3}{2}, \frac{3}{2}, \frac{3}{4}, \frac{7}{2}, \frac{5}{2}, \frac{5}{2}, \frac{5}{2}, \frac{11}{4}, \frac{11}{4}, 3, \frac{15}{4}, \frac{15}{4}, 5\}$
$\{1, 1\}$	$\{3, 3, 3, 3\}$	$x_1^3 + x_2^3 + x_3^3 + x_4^3$	Smt	1	6	11	5	7	0	0	0	$\{2, 2, 2, 2, 3\}$
$\{1, 1\}$	$\{3, 3, 3, 4\}$	$x_4^4 + x_1^3 + x_2^3 + x_3^3$	Sgl	1	0	12	12	1	-1	2	\mathbb{Z}_4^2	$\{\frac{4}{3}, \frac{4}{3}, \frac{5}{3}, \frac{5}{3}, \frac{5}{3}, \frac{5}{3}, 2, \frac{8}{3}, \frac{8}{3}, 3, 4\}$
$\{2, 1\}$	$\{2, 3, 7, 5\}$	$x_3^7 + x_4^5 x_3 + x_2^3 + x_1^2$	Sgl	1	0	29	29	1	0	0	0	$\{\frac{32}{31}, \frac{38}{31}, \frac{44}{31}, \frac{48}{31}, \frac{50}{31}, \frac{54}{31}, \frac{58}{31}, \frac{60}{31}, \frac{66}{31}, \frac{68}{31}, \frac{72}{31}, \frac{74}{31}, \frac{78}{31}, \frac{80}{31}, \frac{84}{31}, \frac{102}{31}, \frac{104}{31}, \frac{108}{31}, \frac{110}{31}, \frac{114}{31}, \frac{120}{31}, \frac{138}{31}, \frac{140}{31}, \frac{144}{31}, \frac{150}{31}, \frac{174}{31}, \frac{180}{31}, \frac{210}{31}\}$
$\{2, 1\}$	$\{2, 4, 5, 3\}$	$x_3^5 + x_4^3 x_3 + x_2^4 + x_1^2$	Sgl	1	1	17	16	2	0	0	0	$\{\frac{14}{13}, \frac{17}{13}, \frac{18}{13}, \frac{20}{13}, \frac{21}{13}, \frac{24}{13}, \frac{28}{13}, \frac{29}{13}, \frac{30}{13}, \frac{32}{13}, \frac{33}{13}, \frac{36}{13}, \frac{44}{13}, \frac{45}{13}, \frac{48}{13}, \frac{60}{13}\}$
$\{2, 1\}$	$\{3, 3, 4, 2\}$	$x_3^4 + x_4^2 x_3 + x_1^3 + x_2^3$	Sgl	1	2	11	9	3	0	0	0	$\{\frac{8}{7}, \frac{10}{7}, \frac{12}{7}, \frac{15}{7}, \frac{16}{7}, \frac{18}{7}, \frac{24}{7}\}$
$\{7, 1\}$	$\{2, 5, 3, 3\}$	$x_2^5 + x_3^3 x_2 + x_3 x_4^3 + x_1^2$	Sgl	1	0	17	17	1	0	0	0	$\{\frac{20}{19}, \frac{24}{19}, \frac{26}{19}, \frac{28}{19}, \frac{30}{19}, \frac{32}{19}, \frac{36}{19}, \frac{42}{19}, \frac{44}{19}, \frac{46}{19}, \frac{48}{19}, \frac{50}{19}, \frac{54}{19}, \frac{56}{19}, \frac{66}{19}, \frac{68}{19}, \frac{72}{19}, \frac{90}{19}\}$
$\{7, 1\}$	$\{3, 4, 2, 2\}$	$x_2^4 + x_3^2 x_2 + x_1^3 + x_3 x_4^2$	Sgl	1	0	11	11	1	0	0	0	$\{\frac{14}{13}, \frac{17}{13}, \frac{18}{13}, \frac{20}{13}, \frac{21}{13}, \frac{24}{13}, \frac{30}{13}, \frac{32}{13}, \frac{33}{13}, \frac{36}{13}, \frac{48}{13}\}$
$\{11, 1\}$	$\{4, 2, 2, 2\}$	$x_1^4 + x_2^2 x_1 + x_2 x_3^2 + x_3 x_4^2$	Sgl	1	1	11	10	2	0	0	0	$\{\frac{10}{9}, \frac{4}{3}, \frac{13}{9}, \frac{14}{9}, \frac{16}{9}, \frac{20}{9}, \frac{7}{3}, \frac{22}{9}, \frac{8}{3}, \frac{32}{9}\}$

Table 15. Rank $r = 1$: smoothable models. ‘Divs’ indicates whether all exceptional divisors of the resolution are smooth (Smt) or some are singular (Sgl). We use black color for the ones with a known Lagrangian \mathcal{J}_X^{4d} , blue color for trinomial models, red color for generalised $D_p^b(G)$ models and orange color for the unknown ones.

Type	(a, b, c, d)	F	Divs	r	f	d_H	\widehat{d}_H	Δ_r	b_3	t_2	CB Spectrum	
$\{1, 1\}$	$\{2, 3, 6, 18\}$	$x_4^{18} + x_3^6 + x_2^3 + x_1^2$	Smt	3	8	89	81	11	-2	4	\mathbb{Z}_2^3	$\{2^9, 3^9, 4^8, 5^8, 6^8, 7^6, 8^6, 10^4, 11^4, 12^4, 13^2, 14^2, 15^2, 16, 17, 18\}$
$\{1, 1\}$	$\{2, 3, 6, 19\}$	$x_4^{19} + x_3^6 + x_2^3 + x_1^2$	Sgl	3	0	90	90	3	-3	6	\mathbb{Z}_{19}^2	$\left\{ \frac{7}{6}, \frac{4^2}{3}, \frac{3^2}{2}, \frac{5^2}{3}, \frac{11}{2}, \frac{13}{3}, \frac{7^2}{6}, \frac{5^2}{3}, \frac{8^2}{2}, \frac{17}{3}, \frac{19}{6}, \frac{10^2}{3}, \frac{7^2}{2}, \frac{11^2}{3}, \frac{23}{6}, \frac{4}{3}, \frac{13^2}{2}, \frac{9^2}{3}, \frac{14^2}{6}, \frac{29}{3}, \frac{5}{6}, \frac{16^2}{3}, \frac{11^2}{2}, \frac{17^2}{3}, \frac{35}{6}, \frac{6}{3}, \frac{19^2}{3}, \frac{13^2}{2}, \frac{20^2}{3}, \frac{41}{6}, \frac{7}{2}, \frac{15^2}{3}, \frac{23^2}{6}, \frac{47}{3}, \frac{17^2}{6}, \frac{26^2}{3}, \frac{53}{9}, \frac{19^2}{2}, \frac{29^2}{6}, \frac{59}{10}, \frac{32^2}{6}, \frac{65}{6}, \frac{11}{3}, \frac{35^2}{2}, \frac{71}{6}, \frac{12}{3}, \frac{38^2}{6}, \frac{77}{3}, \frac{13}{6}, \frac{83}{3}, \frac{14}{6}, \frac{89}{6}, \frac{15}{6}, \frac{95}{6}, \frac{16}{6}, \frac{17}{6}, \frac{18}{6}, \frac{19}{6} \right\}$
$\{1, 1\}$	$\{2, 3, 7, 14\}$	$x_4^{14} + x_3^7 + x_2^3 + x_1^2$	Sgl	3	0	78	78	3	-3	6	\mathbb{Z}_3^3	$\left\{ \frac{3^5}{2}, \frac{2^5}{3}, \frac{3^6}{2}, \frac{7^4}{2}, \frac{9^6}{2}, \frac{5^4}{2}, \frac{6^6}{2}, \frac{13^3}{2}, \frac{15^5}{2}, \frac{8^3}{2}, \frac{9^5}{2}, \frac{19^2}{2}, \frac{21^4}{2}, \frac{11^2}{2}, \frac{12^4}{2}, \frac{25}{2}, \frac{27^3}{2}, \frac{14}{2}, \frac{15^3}{2}, \frac{33^2}{2}, \frac{18^2}{2}, \frac{39}{2}, \frac{21}{2} \right\}$
$\{1, 1\}$	$\{2, 3, 9, 9\}$	$x_3^9 + x_4^9 + x_2^3 + x_1^2$	Smt	3	0	64	64	3	-7	14	\mathbb{Z}_2^{14}	$\{2^{13}, 4^{13}, 6^{11}, 8^9, 10^7, 12^5, 14^3, 16^2, 18\}$
$\{1, 1\}$	$\{2, 4, 4, 12\}$	$x_4^{12} + x_2^4 + x_3^4 + x_1^2$	Smt	3	7	53	46	10	-2	4	\mathbb{Z}_3^3	$\{2^8, 3^8, 4^6, 5^6, 6^6, 7^3, 8^3, 9^3, 10, 11, 12\}$
$\{1, 1\}$	$\{2, 4, 4, 13\}$	$x_4^{13} + x_2^4 + x_3^4 + x_1^2$	Sgl	3	0	54	54	3	-3	6	\mathbb{Z}_{13}^2	$\left\{ \frac{5^2}{4}, \frac{3^3}{2}, \frac{7^2}{4}, \frac{9^2}{4}, \frac{5^3}{4}, \frac{11^2}{4}, \frac{3}{4}, \frac{13^2}{4}, \frac{7^3}{4}, \frac{15^2}{4}, \frac{4}{4}, \frac{9^3}{4}, \frac{19^2}{4}, \frac{5}{4}, \frac{11^3}{4}, \frac{23^2}{4}, \frac{6}{4}, \frac{13^3}{4}, \frac{27^2}{4}, \frac{7}{4}, \frac{31^2}{4}, \frac{8}{4}, \frac{35^2}{4}, \frac{9}{4}, \frac{39^2}{4}, \frac{10}{4}, \frac{11}{4}, \frac{12}{4}, \frac{13}{4} \right\}$
$\{1, 1\}$	$\{3, 3, 3, 9\}$	$x_3^9 + x_1^3 + x_2^3 + x_3^3$	Smt	3	6	35	29	9	-2	4	\mathbb{Z}_3^2	$\{2^7, 3^7, 4^4, 5^4, 6^4, 7, 8, 9\}$
$\{1, 1\}$	$\{3, 3, 3, 10\}$	$x_4^{10} + x_1^3 + x_2^3 + x_3^3$	Sgl	3	0	36	36	3	-3	6	\mathbb{Z}_{10}^2	$\left\{ \frac{4^3}{3}, \frac{5^3}{3}, \frac{2}{3}, \frac{7^3}{3}, \frac{8^3}{3}, \frac{10^3}{3}, \frac{11^3}{3}, \frac{4}{3}, \frac{14^3}{3}, \frac{5}{3}, \frac{17^3}{3}, \frac{6}{3}, \frac{20^3}{3}, \frac{7}{3}, \frac{8}{3}, \frac{9}{3}, \frac{10}{3} \right\}$
$\{1, 1\}$	$\{3, 4, 4, 4\}$	$x_2^4 + x_3^4 + x_4^4 + x_1^3$	Smt	3	0	27	27	3	-6	12	\mathbb{Z}_3^3	$\{2^6, 3^7, 5^3, 6^6, 8, 9, 12\}$
$\{2, 1\}$	$\{2, 5, 7, 4\}$	$x_3^7 + x_4^4 x_3 + x_2^5 + x_1^2$	Sgl	3	0	44	44	3	-1	2	\mathbb{Z}_5^2	$\left\{ \frac{5}{4}, \frac{3^2}{2}, \frac{7}{4}, \frac{2}{4}, \frac{5^2}{4}, \frac{11^2}{4}, \frac{3^2}{4}, \frac{13}{4}, \frac{15}{4}, \frac{4^2}{4}, \frac{17}{4}, \frac{9}{4}, \frac{5^2}{4}, \frac{21}{4}, \frac{25^2}{4}, \frac{13^2}{2}, \frac{27}{4}, \frac{7}{2}, \frac{15^2}{4}, \frac{31}{4}, \frac{8}{4}, \frac{35}{4}, \frac{9}{4}, \frac{10^2}{4}, \frac{41}{4}, \frac{21}{4}, \frac{45}{4}, \frac{23}{4}, \frac{25}{4}, \frac{55}{4}, \frac{14}{4}, \frac{15}{4}, \frac{35}{2} \right\}$
$\{2, 1\}$	$\{2, 7, 3, 9\}$	$x_3 x_4^9 + x_2^7 + x_3^3 + x_1^2$	Sgl	3	0	75	75	3	0	0		$\left\{ \frac{20}{19}, \frac{22}{19}, \frac{24}{19}, \frac{30}{19}, \frac{32}{19}, \frac{34}{19}, \frac{36}{19}, \frac{42}{19}, \frac{44}{19}, \frac{46}{19}, \frac{48}{19}, \frac{50}{19}, \frac{52}{19}, \frac{56}{19}, \frac{58}{19}, \frac{60}{19}, \frac{62}{19}, \frac{70}{19}, \frac{72}{19}, \frac{74}{19}, \frac{78}{19}, \frac{80}{19}, \frac{84}{19}, \frac{86}{19}, \frac{88}{19}, \frac{90}{19}, \frac{98}{19}, \frac{100}{19}, \frac{102}{19}, \frac{104}{19}, \frac{106}{19}, \frac{108}{19}, \frac{112}{19}, \frac{116}{19}, \frac{126}{19}, \frac{128}{19}, \frac{130}{19}, \frac{132}{19}, \frac{134}{19}, \frac{140}{19}, \frac{142}{19}, \frac{144}{19}, \frac{154}{19}, \frac{156}{19}, \frac{158}{19}, \frac{160}{19}, \frac{162}{19}, \frac{168}{19}, \frac{170}{19}, \frac{182}{19}, \frac{184}{19}, \frac{186}{19}, \frac{188}{19}, \frac{196}{19}, \frac{198}{19}, \frac{210}{19}, \frac{212}{19}, \frac{214}{19}, \frac{216}{19}, \frac{218}{19}, \frac{220}{19}, \frac{222}{19}, \frac{224}{19}, \frac{226}{19}, \frac{228}{19}, \frac{230}{19}, \frac{232}{19}, \frac{234}{19}, \frac{236}{19}, \frac{238}{19}, \frac{240}{19}, \frac{242}{19}, \frac{252}{19}, \frac{14}{19}, \frac{268}{19}, \frac{270}{19}, \frac{294}{19}, \frac{296}{19}, \frac{322}{19}, \frac{324}{19}, \frac{350}{19}, \frac{378}{19} \right\}$
$\{2, 1\}$	$\{2, 7, 7, 3\}$	$x_2^7 + x_3^7 + x_3 x_4^3 + x_1^2$	Smt	3	0	45	45	3	-6	12	\mathbb{Z}_2^{12}	$\{2^{12}, 4^{11}, 6^9, 8^6, 10^4, 12^2, 14\}$

Table 18. Rank $r = 3$: smoothable models. ‘Divs’ indicates whether all exceptional divisors of the resolution are smooth (Smt) or some are singular (Sgl).

Type	(a, b, c, d)	F	Divs	r	f	d_H	\widehat{d}_H	$\Delta\mathcal{A}_r$	b_3	t_2	CB Spectrum
$\{2, 1\}$	$\{3, 5, 9, 2\}$	$x_3^9 + x_4^2 x_3 + x_2^5 + x_1^3$	Sgl	3	0	40	40	-1	2	\mathbb{Z}_5^2	$\left\{ \frac{5^2}{4}, \frac{3^2}{2}, \frac{7^3}{4}, 2, \frac{5^3}{2}, \frac{11^2}{2}, 3^2, \frac{13}{4}, \frac{15^2}{4}, 4^3, \frac{17}{4}, \frac{9}{2}, 5^2, \frac{21^2}{4}, \frac{25^3}{2}, \frac{13}{2}, \frac{27}{2}, \frac{15^2}{2}, \frac{31}{2}, \frac{35}{2}, 9, 10, \frac{45}{2} \right\}$
$\{2, 1\}$	$\{3, 8, 4, 2\}$	$x_2^8 + x_3^4 + x_1^3 + x_3 x_4^2$	Sgl	3	0	35	35	-2	4	\mathbb{Z}_3^4	$\left\{ \frac{3^4}{2}, 2^4, \frac{3^5}{2}, \frac{7^3}{2}, \frac{9^4}{2}, 5^2, 6^4, \frac{13}{2}, \frac{15^3}{2}, 8, 9, 9, \frac{21}{2}, 12 \right\}$
$\{2, 1\}$	$\{4, 4, 4, 2\}$	$x_1^4 + x_2^4 + x_3^4 + x_3 x_4^2$	Smt	3	3	24	21	-3	6	\mathbb{Z}_2^6	$\{2^8, 3^2, 4^6, 5, 6^3, 8\}$
$\{4, 1\}$	$\{3, 3, 9, 2\}$	$x_3^9 + x_4^2 x_3 + x_1^3 + x_1 x_2^3$	Smt	3	8	39	31	0	0	0	$\{2^7, 3^7, 4^5, 5^5, 6^3, 7^2, 8, 9\}$
$\{4, 1\}$	$\{3, 3, 10, 2\}$	$x_3^{10} + x_4^2 x_3 + x_1^3 + x_1 x_2^3$	Sgl	3	1	39	38	4	0	0	$\left\{ \frac{20}{19}, \frac{24}{19}, \frac{26}{19}, \frac{30}{19}, \frac{32}{19}, \frac{36}{19}, \frac{39}{19}, \frac{42}{19}, \frac{44}{19}, \frac{46}{19}, \frac{48}{19}, \frac{50}{19}, \frac{54}{19}, \frac{59}{19}, \frac{62}{19}, \frac{64}{19}, \frac{66}{19}, \frac{68}{19}, \frac{72}{19}, \frac{80}{19}, \frac{82}{19}, \frac{84}{19}, \frac{86}{19}, \frac{90}{19}, \frac{100}{19}, \frac{102}{19}, \frac{104}{19}, \frac{108}{19}, \frac{120}{19}, \frac{122}{19}, \frac{126}{19}, \frac{140}{19}, \frac{144}{19}, \frac{162}{19}, \frac{180}{19} \right\}$
$\{4, 1\}$	$\{3, 5, 4, 2\}$	$x_1 x_2^5 + x_3^4 + x_1^3 + x_3 x_4^2$	Sgl	3	1	33	32	4	0	0	$\left\{ \frac{12}{11}, \frac{14}{11}, \frac{18}{11}, \frac{19}{11}, \frac{20}{11}, \frac{24}{11}, \frac{26}{11}, \frac{27}{11}, \frac{28}{11}, \frac{30}{11}, \frac{32}{11}, \frac{34}{11}, \frac{35}{11}, \frac{40}{11}, \frac{42}{11}, \frac{43}{11}, \frac{48}{11}, \frac{50}{11}, \frac{56}{11}, \frac{58}{11}, \frac{59}{11}, \frac{60}{11}, \frac{64}{11}, \frac{72}{11}, \frac{74}{11}, \frac{75}{11}, \frac{80}{11}, \frac{90}{11}, \frac{104}{11}, \frac{120}{11}, \frac{144}{11}, \frac{162}{11}, \frac{180}{11} \right\}$
$\{4, 1\}$	$\{7, 2, 7, 2\}$	$x_1^7 + x_2^2 x_1 + x_3^7 + x_3 x_4^2$	Smt	3	10	37	27	13	0	0	$\{2^8, 3^7, 4^5, 5^3, 6^2, 7\}$
$\{4, 1\}$	$\{7, 2, 8, 2\}$	$x_3^8 + x_4^2 x_3 + x_1^7 + x_1 x_2^2$	Sgl	3	2	37	35	5	0	0	$\left\{ \frac{16}{15}, \frac{6}{5}, \frac{4}{3}, \frac{22^2}{15}, \frac{8}{5}, \frac{26}{15}, \frac{28}{15}, \frac{31}{15}, \frac{32}{15}, \frac{34}{15}, \frac{38}{15}, \frac{8}{3}, \frac{14}{5}, \frac{47}{15}, \frac{16}{3}, \frac{10^2}{3}, \frac{52}{15}, \frac{18}{5}, \frac{21}{5}, \frac{64^2}{15}, \frac{22}{15}, \frac{68}{15}, \frac{14}{3}, \frac{16}{5}, \frac{82}{15}, \frac{28}{3}, \frac{32}{5}, \frac{98}{15}, \frac{112}{15} \right\}$
$\{7, 1\}$	$\{2, 5, 3, 7\}$	$x_3 x_4^7 + x_2^5 + x_2 x_3^3 + x_1^2$	Sgl	3	0	47	47	3	0	0	$\left\{ \frac{16}{15}, \frac{6}{5}, \frac{4}{3}, \frac{28}{15}, \frac{8}{5}, \frac{26}{15}, \frac{32}{15}, \frac{34}{15}, \frac{38}{15}, \frac{44}{15}, \frac{46}{15}, \frac{56}{15}, \frac{58}{15}, \frac{59}{15}, \frac{62}{15}, \frac{64}{15}, \frac{66}{15}, \frac{68}{15}, \frac{72}{15}, \frac{80}{15}, \frac{82}{15}, \frac{84}{15}, \frac{86}{15}, \frac{90}{15}, \frac{104}{15}, \frac{112}{15}, \frac{120}{15}, \frac{122}{15}, \frac{124}{15}, \frac{126}{15}, \frac{140}{15}, \frac{144}{15}, \frac{166}{15}, \frac{188}{15}, \frac{192}{15} \right\}$
$\{7, 1\}$	$\{3, 5, 7, 2\}$	$x_2 x_3^7 + x_4^2 x_3 + x_2^5 + x_1^3$	Sgl	3	0	39	39	3	0	0	$\left\{ \frac{20}{19}, \frac{24}{19}, \frac{26}{19}, \frac{30}{19}, \frac{32}{19}, \frac{36}{19}, \frac{42}{19}, \frac{44}{19}, \frac{47}{19}, \frac{48}{19}, \frac{50}{19}, \frac{54}{19}, \frac{56}{19}, \frac{59}{19}, \frac{62}{19}, \frac{64}{19}, \frac{66}{19}, \frac{68}{19}, \frac{72}{19}, \frac{74}{19}, \frac{75}{19}, \frac{78}{19}, \frac{84}{19}, \frac{90}{19}, \frac{92}{19}, \frac{96}{19}, \frac{102}{19}, \frac{104}{19}, \frac{108}{19}, \frac{112}{19}, \frac{120}{19}, \frac{122}{19}, \frac{126}{19}, \frac{138}{19}, \frac{140}{19}, \frac{144}{19}, \frac{162}{19}, \frac{168}{19}, \frac{186}{19}, \frac{210}{19} \right\}$
$\{7, 1\}$	$\{3, 6, 2, 3\}$	$x_2^6 + x_3^2 x_2 + x_1^3 + x_3 x_4^3$	Sgl	3	2	30	28	5	0	0	$\left\{ \frac{5^2}{4}, \frac{3^2}{2}, 2, \frac{9^2}{2}, \frac{5^2}{2}, \frac{11^2}{2}, 3^2, \frac{7}{2}, \frac{15}{4}, \frac{17^2}{4}, \frac{9^2}{2}, \frac{21}{2}, \frac{23}{2}, \frac{29}{2}, \frac{31}{2}, 6^2, \frac{29}{2}, \frac{15}{2}, 9 \right\}$
$\{7, 1\}$	$\{4, 7, 2, 2\}$	$x_2^7 + x_3^2 x_2 + x_1^4 + x_3 x_4^2$	Sgl	3	0	30	30	-1	2	\mathbb{Z}_4^2	$\left\{ \frac{4^3}{3}, \frac{5^3}{3}, 2^2, \frac{8^3}{3}, 3^3, \frac{10}{3}, 4^3, \frac{13^2}{3}, \frac{14}{3}, \frac{16^3}{3}, \frac{17}{3}, \frac{20^2}{3}, 7, 8, \frac{28}{3} \right\}$

Table 19. Rank $r = 3$: smoothable models. ‘Divs’ indicates whether all exceptional divisors of the resolution are smooth (Smt) or some are singular (Sgl).

Type	(a, b, c, d)	F	Divs	r	f	d_H	\hat{r}	$\widehat{d_H}$	$\Delta\mathcal{A}_r$	b_3	t_2	CB Spectrum
$\{10, 1\}$	$\{3, 2, 5, 3\}$	$x_4x_3^5 + x_2^2x_3 + x_1^3 + x_2x_4^3$	Sgl	3	0	30	30	3	0	0	0	$\{\frac{11}{10}, \frac{6}{5}, \frac{7}{5}, \frac{3}{2}, \frac{17}{10}, \frac{9}{5}, \frac{21}{10}, \frac{13}{5}, \frac{27}{5}, \frac{29}{10}, \frac{3}{5}, \frac{16}{10}, \frac{33}{10}, \frac{18}{5}, \frac{39}{10}, \frac{21}{5}, \frac{22}{5}, \frac{9}{2}, \frac{47}{10}, \frac{24}{5}, \frac{27}{5}, \frac{57}{10}, \frac{31}{5}, \frac{63}{10}, \frac{15}{5}, \frac{39}{10}, \frac{93}{10}\}$
$\{10, 1\}$	$\{5, 2, 2, 3\}$	$x_1^5 + x_2x_4^3 + x_2^2x_3 + x_3^2x_4$	Sgl	3	0	24	24	3	0	0	0	$\{\frac{9}{8}, \frac{5}{4}, \frac{11}{8}, \frac{3}{4}, \frac{7}{8}, \frac{15}{8}, \frac{17}{8}, \frac{19}{8}, \frac{5}{4}, \frac{11}{3}, \frac{25}{4}, \frac{13}{4}, \frac{27}{8}, \frac{15}{4}, \frac{35}{8}, \frac{37}{8}, \frac{39}{8}, \frac{5}{2}, \frac{45}{8}, \frac{25}{4}, \frac{13}{2}, \frac{65}{8}\}$
$\{11, 1\}$	$\{4, 5, 2, 2\}$	$x_1x_2^5 + x_3^2x_2 + x_1^4 + x_3x_4^2$	Sgl	3	1	29	28	4	0	0	0	$\{\frac{10}{9}, \frac{4}{3}, \frac{13}{9}, \frac{14}{9}, \frac{16}{9}, \frac{17}{9}, \frac{20}{9}, \frac{7}{3}, \frac{22}{9}, \frac{8}{3}, \frac{25}{9}, \frac{26}{9}, \frac{28}{9}, \frac{32}{9}, \frac{11}{3}, \frac{34}{9}, \frac{4}{3}, \frac{37}{9}, \frac{40}{9}, \frac{44}{9}, \frac{5}{3}, \frac{46}{9}, \frac{16}{3}, \frac{56}{9}, \frac{19}{3}, \frac{20}{3}, \frac{68}{9}, \frac{80}{9}\}$
$\{11, 1\}$	$\{5, 2, 2, 3\}$	$x_1^5 + x_2^2x_1 + x_3x_4^3 + x_2x_3^2$	Sgl	3	2	24	22	5	0	0	0	$\{\frac{5}{4}, \frac{3}{2}, \frac{2}{2}, \frac{9}{2}, \frac{5}{2}, \frac{11}{2}, \frac{3}{2}, \frac{7}{2}, \frac{15}{2}, \frac{4}{4}, \frac{17}{4}, \frac{9}{2}, \frac{21}{4}, \frac{23}{4}, \frac{6}{2}, \frac{15}{2}\}$
$\{11, 1\}$	$\{6, 3, 2, 2\}$	$x_1^6 + x_2^3x_1 + x_2x_3^2 + x_3x_4^2$	Sgl	3	1	25	24	4	0	0	0	$\{\frac{10}{9}, \frac{4}{3}, \frac{13}{9}, \frac{14}{9}, \frac{16}{9}, \frac{17}{9}, \frac{20}{9}, \frac{22}{9}, \frac{8}{3}, \frac{25}{9}, \frac{26}{9}, \frac{28}{9}, \frac{32}{9}, \frac{34}{9}, \frac{4}{3}, \frac{37}{9}, \frac{40}{9}, \frac{46}{9}, \frac{16}{3}, \frac{49}{9}, \frac{52}{9}, \frac{20}{3}, \frac{8}{9}, \frac{9}{9}, \frac{4}{9}, \frac{3}{9}, \frac{9}{9}, \frac{3}{9}, \frac{8}{9}, \frac{8}{9}\}$
$\{11, 1\}$	$\{8, 2, 4, 2\}$	$x_1^8 + x_2^2x_1 + x_2x_3^4 + x_3x_4^2$	Sgl	3	1	37	36	4	0	0	0	$\{\frac{18}{17}, \frac{20}{17}, \frac{22}{17}, \frac{24}{17}, \frac{25}{17}, \frac{26}{17}, \frac{28}{17}, \frac{30}{17}, \frac{32}{17}, \frac{36}{17}, \frac{40}{17}, \frac{41}{17}, \frac{42}{17}, \frac{44}{17}, \frac{46}{17}, \frac{48}{17}, \frac{54}{17}, \frac{56}{17}, \frac{57}{17}, \frac{58}{17}, \frac{60}{17}, \frac{62}{17}, \frac{64}{17}, \frac{72}{17}, \frac{73}{17}, \frac{74}{17}, \frac{76}{17}, \frac{78}{17}, \frac{80}{17}, \frac{92}{17}, \frac{94}{17}, \frac{110}{17}, \frac{112}{17}, \frac{128}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}, \frac{17}{17}\}$
$\{12, 1\}$	$\{7, 2, 3, 2\}$	$x_1^7 + x_3^3x_1 + x_2^2x_1 + x_2x_3^2 + x_2x_4^2$	Smt	3	8	29	21	11	0	0	0	$\{2^6, 3^6, 4^4, 5^3, 6, 7\}$
$\{19, 1\}$	$\{2, 2, 3, 4\}$	$x_2x_4^4 + x_3^3x_4 + x_1x_2^2 + x_1^2x_3$	Sgl	3	2	25	23	5	0	0	0	$\{\frac{7}{6}, \frac{4}{3}, \frac{5}{2}, \frac{11}{6}, \frac{13}{6}, \frac{7}{3}, \frac{5}{3}, \frac{8}{3}, \frac{17}{6}, \frac{3}{6}, \frac{19}{6}, \frac{7}{2}, \frac{11}{2}, \frac{23}{2}, \frac{4}{3}, \frac{13}{3}, \frac{5}{6}, \frac{31}{6}, \frac{16}{3}, \frac{17}{3}, \frac{47}{6}\}$

Table 20. Rank $r = 3$: smoothable models. ‘Divs’ indicates whether all exceptional divisors of the resolution are smooth (Smt) or some are singular (Sgl). (Continued)

Type	(a, b, c, d)	F	Divs	r	f	d_H	\hat{r}	$\widehat{d_H}$	$\Delta\mathcal{A}_r$	b_3	t_2
$\{1, 1\}$	$\{2, 3, 6, 24\}$	$x_4^{24} + x_3^6 + x_2^3 + x_1^2$	Smt	4	8	119	111	12	-3	6	\mathbb{Z}_4^2
$\{1, 1\}$	$\{2, 3, 6, 25\}$	$x_4^{25} + x_3^6 + x_2^3 + x_1^2$	Sgl	4	0	120	120	4	-4	8	\mathbb{Z}_{25}^2
$\{1, 1\}$	$\{2, 4, 4, 16\}$	$x_4^{16} + x_2^4 + x_3^4 + x_1^2$	Smt	4	7	71	64	11	-3	6	\mathbb{Z}_4^2
$\{1, 1\}$	$\{2, 4, 4, 17\}$	$x_4^{17} + x_2^4 + x_3^4 + x_1^2$	Sgl	4	0	72	72	4	-4	8	\mathbb{Z}_{17}^2
$\{1, 1\}$	$\{2, 4, 6, 6\}$	$x_3^6 + x_4^6 + x_2^4 + x_1^2$	Smt	4	5	40	35	9	-2	4	\mathbb{Z}_2^4
$\{1, 1\}$	$\{3, 3, 3, 12\}$	$x_4^{12} + x_1^3 + x_2^3 + x_3^3$	Smt	4	6	47	41	10	-3	6	\mathbb{Z}_4^2
$\{1, 1\}$	$\{3, 3, 3, 13\}$	$x_4^{13} + x_1^3 + x_2^3 + x_3^3$	Sgl	4	0	48	48	4	-4	8	\mathbb{Z}_{13}^2
$\{2, 1\}$	$\{3, 4, 5, 3\}$	$x_3^5 + x_4^3 x_3 + x_2^4 + x_1^3$	Sgl	4	0	33	33	4	-1	2	\mathbb{Z}_4^2
$\{2, 1\}$	$\{3, 5, 15, 2\}$	$x_3^{15} + x_4^2 x_3 + x_2^5 + x_1^3$	Sgl	4	8	68	60	12	0	0	0
$\{2, 1\}$	$\{3, 5, 16, 2\}$	$x_3^{16} + x_4^2 x_3 + x_2^5 + x_1^3$	Sgl	4	0	68	68	4	0	0	0
$\{2, 1\}$	$\{3, 6, 6, 2\}$	$x_2^6 + x_3^6 + x_1^3 + x_3 x_4^2$	Smt	4	2	36	34	6	-4	8	\mathbb{Z}_2^8
$\{4, 1\}$	$\{9, 2, 9, 2\}$	$x_1^9 + x_2^2 x_1 + x_3^9 + x_3 x_4^2$	Smt	4	12	56	44	16	0	0	0
$\{4, 1\}$	$\{9, 2, 10, 2\}$	$x_3^{10} + x_4^2 x_3 + x_1^9 + x_1 x_2^2$	Sgl	4	2	56	54	6	0	0	0
$\{7, 1\}$	$\{2, 7, 4, 4\}$	$x_2^7 + x_3^4 x_2 + x_3 x_4^4 + x_1^2$	Sgl	4	0	45	45	4	-1	2	\mathbb{Z}_4^2
$\{7, 1\}$	$\{2, 12, 3, 4\}$	$x_2^{12} + x_3^3 x_2 + x_3 x_4^4 + x_1^2$	Sgl	4	1	60	59	5	0	0	0
$\{7, 1\}$	$\{4, 8, 2, 2\}$	$x_2^8 + x_3^2 x_2 + x_1^4 + x_3 x_4^2$	Sgl	4	3	36	33	7	0	0	0
$\{7, 1\}$	$\{5, 5, 2, 2\}$	$x_1^5 + x_2^5 + x_2 x_3^2 + x_3 x_4^2$	Smt	4	4	30	26	8	-2	4	\mathbb{Z}_2^4
$\{8, 1\}$	$\{2, 16, 3, 5\}$	$x_2^{16} + x_4^5 x_2 + x_3^3 x_2 + x_1^2 + x_3^2 x_4^2$	Sgl	4	9	76	67	13	0	0	0
$\{10, 1\}$	$\{2, 3, 10, 4\}$	$x_4 x_3^{10} + x_2^3 x_3 + x_2 x_4^4 + x_1^2$	Sgl	4	0	60	60	4	0	0	0
$\{10, 1\}$	$\{4, 2, 2, 6\}$	$x_2 x_4^6 + x_3^2 x_4 + x_1^4 + x_2^2 x_3$	Sgl	4	0	36	36	4	0	0	0
$\{11, 1\}$	$\{7, 3, 2, 2\}$	$x_1^7 + x_2^3 x_1 + x_2 x_3^2 + x_3 x_4^2$	Sgl	4	3	30	27	7	0	0	0
$\{11, 1\}$	$\{10, 2, 5, 2\}$	$x_1^{10} + x_2^2 x_1 + x_2 x_3^5 + x_3 x_4^2$	Sgl	4	1	56	55	5	0	0	0
$\{19, 1\}$	$\{2, 2, 5, 3\}$	$x_4 x_3^5 + x_1^2 x_3 + x_2 x_4^3 + x_1 x_2^2$	Sgl	4	2	31	29	6	0	0	0

Table 21. Rank $r = 4$: smoothable models. ‘Divs’ indicates whether all exceptional divisors of the resolution are smooth (Smt) or some are singular (Sgl). The CB spectra are rather expansive and we include them instead in the ancillary mathematica file.

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References

- [1] C. Closset, S. Schäfer-Nameki and Y.-N. Wang, *Coulomb and Higgs branches from canonical singularities. Part 0*, *JHEP* **02** (2021) 003 [[arXiv:2007.15600](#)] [[INSPIRE](#)].
- [2] A.D. Shapere and C. Vafa, *BPS structure of Argyres-Douglas superconformal theories*, [hep-th/9910182](#) [[INSPIRE](#)].
- [3] D. Xie and S.-T. Yau, *4d $N = 2$ SCFT and singularity theory. Part I. Classification*, [arXiv:1510.01324](#) [[INSPIRE](#)].
- [4] N. Seiberg, *Five-dimensional SUSY field theories, nontrivial fixed points and string dynamics*, *Phys. Lett. B* **388** (1996) 753 [[hep-th/9608111](#)] [[INSPIRE](#)].
- [5] D.R. Morrison and N. Seiberg, *Extremal transitions and five-dimensional supersymmetric field theories*, *Nucl. Phys. B* **483** (1997) 229 [[hep-th/9609070](#)] [[INSPIRE](#)].
- [6] C. Closset, S. Giacomelli, S. Schäfer-Nameki and Y.-N. Wang, *5d and 4d SCFTs: canonical singularities, trinions and S-dualities*, *JHEP* **05** (2021) 274 [[arXiv:2012.12827](#)] [[INSPIRE](#)].
- [7] C. Closset, S. Schäfer-Nameki and Y.-N. Wang, *Coulomb and Higgs branches from canonical singularities. Part II*, to appear.
- [8] S.S.-T. Yau and Y. Yu, *Classification of 3-dimensional isolated rational hypersurface singularities with c^* -action*, *Rocky Mountain J. Math.* **35** (2005) 1795.
- [9] I.C. Davenport and I.V. Melnikov, *Landau-Ginzburg skeletons*, *JHEP* **05** (2017) 050 [[arXiv:1608.04259](#)] [[INSPIRE](#)].
- [10] M. Buican and H. Jiang, *1-form symmetry, isolated $N = 2$ SCFTs, and Calabi-Yau threefolds*, *JHEP* **12** (2021) 024 [[arXiv:2106.09807](#)] [[INSPIRE](#)].
- [11] A. Bourget et al., *The Higgs mechanism — Hasse diagrams for symplectic singularities*, *JHEP* **01** (2020) 157 [[arXiv:1908.04245](#)] [[INSPIRE](#)].
- [12] A. Bourget, J.F. Grimminger, A. Hanany, M. Sperling and Z. Zhong, *Magnetic quivers from brane webs with $O5$ planes*, *JHEP* **07** (2020) 204 [[arXiv:2004.04082](#)] [[INSPIRE](#)].
- [13] A. Bourget, J.F. Grimminger, A. Hanany, M. Sperling, G. Zafrir and Z. Zhong, *Magnetic quivers for rank 1 theories*, *JHEP* **09** (2020) 189 [[arXiv:2006.16994](#)] [[INSPIRE](#)].
- [14] A. Bourget, J.F. Grimminger, A. Hanany, R. Kalveks, M. Sperling and Z. Zhong, *Magnetic lattices for orthosymplectic quivers*, *JHEP* **12** (2020) 092 [[arXiv:2007.04667](#)] [[INSPIRE](#)].
- [15] M. van Beest, A. Bourget, J. Eckhard and S. Schäfer-Nameki, *(Symplectic) leaves and (5d Higgs) branches in the Poly(go)nesian tropical rain forest*, *JHEP* **11** (2020) 124 [[arXiv:2008.05577](#)] [[INSPIRE](#)].
- [16] M. Van Beest, A. Bourget, J. Eckhard and S. Schäfer-Nameki, *(5d RG-flow) trees in the tropical rain forest*, *JHEP* **03** (2021) 241 [[arXiv:2011.07033](#)] [[INSPIRE](#)].
- [17] M. Akhond, F. Carta, S. Dwivedi, H. Hayashi, S.-S. Kim and F. Yagi, *Five-brane webs, Higgs branches and unitary/orthosymplectic magnetic quivers*, *JHEP* **12** (2020) 164 [[arXiv:2008.01027](#)] [[INSPIRE](#)].
- [18] A. Bourget, J.F. Grimminger, A. Hanany, M. Sperling and Z. Zhong, *Branes, quivers, and the affine Grassmannian*, [arXiv:2102.06190](#) [[INSPIRE](#)].
- [19] A. Bourget, J.F. Grimminger, M. Martone and G. Zafrir, *Magnetic quivers for rank 2 theories*, *JHEP* **03** (2022) 208 [[arXiv:2110.11365](#)] [[INSPIRE](#)].

- [20] A. Bourget, J.F. Grimminger, A. Hanany, R. Kalveks and Z. Zhong, *Higgs branches of U/SU quivers via brane locking*, [arXiv:2111.04745](#) [[INSPIRE](#)].
- [21] M. Akhond, F. Carta, S. Dwivedi, H. Hayashi, S.-S. Kim and F. Yagi, *Factorised 3d $N = 4$ orthosymplectic quivers*, *JHEP* **05** (2021) 269 [[arXiv:2101.12235](#)] [[INSPIRE](#)].
- [22] M. van Beest and S. Giacomelli, *Connecting 5d Higgs branches via Fayet-Iliopoulos deformations*, *JHEP* **12** (2021) 202 [[arXiv:2110.02872](#)] [[INSPIRE](#)].
- [23] S. Nawata, M. Sperling, H.E. Wang and Z. Zhong, *Magnetic quivers and line defects — on a duality between 3d $N = 4$ unitary and orthosymplectic quivers*, *JHEP* **02** (2022) 174 [[arXiv:2111.02831](#)] [[INSPIRE](#)].
- [24] M. Sperling and Z. Zhong, *Balanced B and D-type orthosymplectic quivers — magnetic quivers for product theories*, [arXiv:2111.00026](#) [[INSPIRE](#)].
- [25] A. Collinucci and R. Valandro, *A string theory realization of special unitary quivers in 3 dimensions*, *JHEP* **11** (2020) 157 [[arXiv:2008.10689](#)] [[INSPIRE](#)].
- [26] A. Collinucci, M. De Marco, A. Sangiovanni and R. Valandro, *Higgs branches of 5d rank-zero theories from geometry*, *JHEP* **10** (2021) 018 [[arXiv:2105.12177](#)] [[INSPIRE](#)].
- [27] M. De Marco and A. Sangiovanni, *Higgs branches of rank-0 5d theories from M-theory on (A_j, A_l) and (A_k, D_n) singularities*, *JHEP* **03** (2022) 099 [[arXiv:2111.05875](#)] [[INSPIRE](#)].
- [28] M. Reid et al., *Canonical 3-folds*, in *Journées de géométrie algébrique d'Angers*, Siythoff and Noordhoff, Alphen aan den Rijn, The Netherlands (1980), p. 273.
- [29] V.I. Arnol'd, *Critical points of smooth functions and their normal forms*, *Russian Math. Surv.* **30** (1975) 1.
- [30] S. Katz, P. Mayr and C. Vafa, *Mirror symmetry and exact solution of 4D $N = 2$ gauge theories: 1*, *Adv. Theor. Math. Phys.* **1** (1998) 53 [[hep-th/9706110](#)] [[INSPIRE](#)].
- [31] J.N. Mather and S.S.-T. Yau, *Classification of isolated hypersurface singularities by their moduli algebras*, *Invent. Math.* **69** (1982) 243.
- [32] A.D. Shapere and Y. Tachikawa, *Central charges of $N = 2$ superconformal field theories in four dimensions*, *JHEP* **09** (2008) 109 [[arXiv:0804.1957](#)] [[INSPIRE](#)].
- [33] Y. Wang, D. Xie, S.S.T. Yau and S.-T. Yau, *4d $N = 2$ SCFT from complete intersection singularity*, *Adv. Theor. Math. Phys.* **21** (2017) 801 [[arXiv:1606.06306](#)] [[INSPIRE](#)].
- [34] M. Caorsi and S. Cecotti, *Geometric classification of 4d $N = 2$ SCFTs*, *JHEP* **07** (2018) 138 [[arXiv:1801.04542](#)] [[INSPIRE](#)].
- [35] P.C. Argyres and M. Martone, *Towards a classification of rank r $N = 2$ SCFTs. Part II. Special Kähler stratification of the Coulomb branch*, *JHEP* **12** (2020) 022 [[arXiv:2007.00012](#)] [[INSPIRE](#)].
- [36] S. Cecotti, M. Del Zotto, M. Martone and R. Moscrop, *The characteristic dimension of four-dimensional $N = 2$ SCFTs*, [arXiv:2108.10884](#) [[INSPIRE](#)].
- [37] M. Reid et al., *Minimal models of canonical 3-folds*, in *Algebraic varieties and analytic varieties*, Mathematical Society of Japan, Japan (1983), p. 131.
- [38] M. Reid et al., *Young person's guide to canonical singularities*, *Alg. Geom. Bowdoin* **46** (1985) 345.

- [39] K.A. Intriligator, D.R. Morrison and N. Seiberg, *Five-dimensional supersymmetric gauge theories and degenerations of Calabi-Yau spaces*, *Nucl. Phys. B* **497** (1997) 56 [[hep-th/9702198](#)] [[INSPIRE](#)].
- [40] P. Jefferson, H.-C. Kim, C. Vafa and G. Zafrir, *Towards classification of 5d SCFTs: single gauge node*, [arXiv:1705.05836](#) [[INSPIRE](#)].
- [41] C. Closset, M. Del Zotto and V. Saxena, *Five-dimensional SCFTs and gauge theory phases: an M-theory/type IIA perspective*, *SciPost Phys.* **6** (2019) 052 [[arXiv:1812.10451](#)] [[INSPIRE](#)].
- [42] C. Closset and M. Del Zotto, *On 5d SCFTs and their BPS quivers. Part I. B-branes and brane tilings*, [arXiv:1912.13502](#) [[INSPIRE](#)].
- [43] V. Saxena, *Rank-two 5d SCFTs from M-theory at isolated toric singularities: a systematic study*, *JHEP* **04** (2020) 198 [[arXiv:1911.09574](#)] [[INSPIRE](#)].
- [44] C. Closset and H. Magureanu, *The U-plane of rank-one 4d $N = 2$ KK theories*, *SciPost Phys.* **12** (2022) 065 [[arXiv:2107.03509](#)] [[INSPIRE](#)].
- [45] M. Caibar, *Minimal models of canonical singularities and their cohomology*, Ph.D. thesis, University of Warwick, U.K. (1999).
- [46] M. Caibar, *Minimal models of canonical 3-fold singularities and their Betti numbers*, *Int. Math. Res. Not.* **2005** (2005) 1563.
- [47] V.V. Batyrev, *Dual polyhedra and mirror symmetry for Calabi-Yau hypersurfaces in toric varieties*, *J. Alg. Geom.* **3** (1994) 493 [[alg-geom/9310003](#)] [[INSPIRE](#)].
- [48] C. Lawrie and S. Schäfer-Nameki, *The Tate form on steroids: resolution and higher codimension fibers*, *JHEP* **04** (2013) 061 [[arXiv:1212.2949](#)] [[INSPIRE](#)].
- [49] F. Apruzzi, C. Lawrie, L. Lin, S. Schäfer-Nameki and Y.-N. Wang, *Fibers add flavor. Part I. Classification of 5d SCFTs, flavor symmetries and BPS states*, *JHEP* **11** (2019) 068 [[arXiv:1907.05404](#)] [[INSPIRE](#)].
- [50] D. Gaiotto, A. Kapustin, N. Seiberg and B. Willett, *Generalized global symmetries*, *JHEP* **02** (2015) 172 [[arXiv:1412.5148](#)] [[INSPIRE](#)].
- [51] D.R. Morrison, S. Schäfer-Nameki and B. Willett, *Higher-form symmetries in 5d*, *JHEP* **09** (2020) 024 [[arXiv:2005.12296](#)] [[INSPIRE](#)].
- [52] F. Albertini, M. Del Zotto, I. García Etxebarria and S.S. Hosseini, *Higher form symmetries and M-theory*, *JHEP* **12** (2020) 203 [[arXiv:2005.12831](#)] [[INSPIRE](#)].
- [53] M. Del Zotto, I. García Etxebarria and S.S. Hosseini, *Higher form symmetries of Argyres-Douglas theories*, *JHEP* **10** (2020) 056 [[arXiv:2007.15603](#)] [[INSPIRE](#)].
- [54] L. Bhardwaj and S. Schäfer-Nameki, *Higher-form symmetries of 6d and 5d theories*, *JHEP* **02** (2021) 159 [[arXiv:2008.09600](#)] [[INSPIRE](#)].
- [55] R.C. Randell, *The homology of generalized Brieskorn manifolds*, *Topology* **14** (1975) 347.
- [56] P. Orlik, *On the homology of weighted homogeneous manifolds*, in *Proceedings of the second conference on compact transformation groups*, H.T. Ku, L.N. Mann, J.L. Sicks and J.C. Su eds., Springer, Berlin, Heidelberg, Germany (1972), p. 260.
- [57] P. Candelas and X.C. de la Ossa, *Comments on conifolds*, *Nucl. Phys. B* **342** (1990) 246 [[INSPIRE](#)].

- [58] B.G. Cooper, *On the monodromy at isolated singularities of weighted homogeneous polynomials*, *Trans. Amer. Math. Soc.* **269** (1982) 149.
- [59] O. Aharony, N. Seiberg and Y. Tachikawa, *Reading between the lines of four-dimensional gauge theories*, *JHEP* **08** (2013) 115 [[arXiv:1305.0318](#)] [[INSPIRE](#)].
- [60] M. Del Zotto, J.J. Heckman, D.S. Park and T. Rudelius, *On the defect group of a 6D SCFT*, *Lett. Math. Phys.* **106** (2016) 765 [[arXiv:1503.04806](#)] [[INSPIRE](#)].
- [61] I. García Etxebarria, B. Heidenreich and D. Regalado, *IIB flux non-commutativity and the global structure of field theories*, *JHEP* **10** (2019) 169 [[arXiv:1908.08027](#)] [[INSPIRE](#)].
- [62] L. Bhardwaj, M. Hubner and S. Schäfer-Nameki, *1-form symmetries of 4d $N = 2$ class S theories*, *SciPost Phys.* **11** (2021) 096 [[arXiv:2102.01693](#)] [[INSPIRE](#)].
- [63] L. Bhardwaj, M. Hubner and S. Schäfer-Nameki, *Liberating confinement from Lagrangians: 1-form symmetries and lines in 4d $N = 1$ from 6d $N = (2, 0)$* , *SciPost Phys.* **12** (2022) 040 [[arXiv:2106.10265](#)] [[INSPIRE](#)].
- [64] S.S. Hosseini and R. Moscrop, *Maruyoshi-Song flows and defect groups of $D_p^b(G)$ theories*, *JHEP* **10** (2021) 119 [[arXiv:2106.03878](#)] [[INSPIRE](#)].
- [65] M. Kreuzer and H. Skarke, *On the classification of quasihomogeneous functions*, *Commun. Math. Phys.* **150** (1992) 137 [[hep-th/9202039](#)] [[INSPIRE](#)].
- [66] Wolfram Research Inc., *Mathematica*, version 12.3.1, <https://www.wolfram.com/mathematica/>, U.S.A. (2021).
- [67] S. Cecotti, A. Neitzke and C. Vafa, *R-twisting and 4d/2d correspondences*, [arXiv:1006.3435](#) [[INSPIRE](#)].
- [68] W. Lerche, C. Vafa and N.P. Warner, *Chiral rings in $N = 2$ superconformal theories*, *Nucl. Phys. B* **324** (1989) 427 [[INSPIRE](#)].
- [69] S. Gukov, C. Vafa and E. Witten, *CFT's from Calabi-Yau four folds*, *Nucl. Phys. B* **584** (2000) 69 [*Erratum ibid.* **608** (2001) 477] [[hep-th/9906070](#)] [[INSPIRE](#)].
- [70] D.G. Markushevich, *Canonical singularities of three-dimensional hypersurfaces*, *Math. U.S.S.R.-Izvestiya* **26** (1986) 315.
- [71] P. Jefferson, S. Katz, H.-C. Kim and C. Vafa, *On geometric classification of 5d SCFTs*, *JHEP* **04** (2018) 103 [[arXiv:1801.04036](#)] [[INSPIRE](#)].
- [72] P. Argyres, M. Lotito, Y. Lü and M. Martone, *Geometric constraints on the space of $N = 2$ SCFTs. Part III. Enhanced Coulomb branches and central charges*, *JHEP* **02** (2018) 003 [[arXiv:1609.04404](#)] [[INSPIRE](#)].
- [73] J.A. Minahan and D. Nemeschansky, *An $N = 2$ superconformal fixed point with E_6 global symmetry*, *Nucl. Phys. B* **482** (1996) 142 [[hep-th/9608047](#)] [[INSPIRE](#)].
- [74] S. Cecotti, M. Del Zotto and S. Giacomelli, *More on the $N = 2$ superconformal systems of type $D_p(G)$* , *JHEP* **04** (2013) 153 [[arXiv:1303.3149](#)] [[INSPIRE](#)].
- [75] L. Bhardwaj and P. Jefferson, *Classifying 5d SCFTs via 6d SCFTs: rank one*, *JHEP* **07** (2019) 178 [*Addendum ibid.* **01** (2020) 153] [[arXiv:1809.01650](#)] [[INSPIRE](#)].
- [76] J. Tian and Y.-N. Wang, *5D and 6D SCFTs from C^3 orbifolds*, [arXiv:2110.15129](#) [[INSPIRE](#)].

- [77] D.R. Heath-Brown, *The density of rational points on Cayley’s cubic surface*, [math.NT/0210333](#).
- [78] F. Apruzzi, C. Lawrie, L. Lin, S. Schäfer-Nameki and Y.-N. Wang, *5d superconformal field theories and graphs*, *Phys. Lett. B* **800** (2020) 135077 [[arXiv:1906.11820](#)].
- [79] F. Apruzzi, C. Lawrie, L. Lin, S. Schäfer-Nameki and Y.-N. Wang, *Fibers add flavor. Part II. 5d SCFTs, gauge theories, and dualities*, *JHEP* **03** (2020) 052 [[arXiv:1909.09128](#)] [[INSPIRE](#)].
- [80] F. Apruzzi, S. Schäfer-Nameki and Y.-N. Wang, *5d SCFTs from decoupling and gluing*, *JHEP* **08** (2020) 153 [[arXiv:1912.04264](#)] [[INSPIRE](#)].
- [81] L. Di Pietro and Z. Komargodski, *Cardy formulae for SUSY theories in $d = 4$ and $d = 6$* , *JHEP* **12** (2014) 031 [[arXiv:1407.6061](#)] [[INSPIRE](#)].
- [82] M. Buican and T. Nishinaka, *On the superconformal index of Argyres-Douglas theories*, *J. Phys. A* **49** (2016) 015401 [[arXiv:1505.05884](#)] [[INSPIRE](#)].
- [83] A. Arabi Ardehali, *High-temperature asymptotics of supersymmetric partition functions*, *JHEP* **07** (2016) 025 [[arXiv:1512.03376](#)] [[INSPIRE](#)].
- [84] C. Beem and L. Rastelli, *Vertex operator algebras, Higgs branches, and modular differential equations*, *JHEP* **08** (2018) 114 [[arXiv:1707.07679](#)] [[INSPIRE](#)].
- [85] S. Giacomelli, *RG flows with supersymmetry enhancement and geometric engineering*, *JHEP* **06** (2018) 156 [[arXiv:1710.06469](#)] [[INSPIRE](#)].
- [86] S. Giacomelli, N. Mekareeya and M. Sacchi, *New aspects of Argyres-Douglas theories and their dimensional reduction*, *JHEP* **03** (2021) 242 [[arXiv:2012.12852](#)] [[INSPIRE](#)].
- [87] M.J. Kang, C. Lawrie and J. Song, *Infinitely many 4D $N = 2$ SCFTs with $a = c$ and beyond*, *Phys. Rev. D* **104** (2021) 105005 [[arXiv:2106.12579](#)] [[INSPIRE](#)].
- [88] F. Carta, S. Giacomelli, N. Mekareeya and A. Mininno, *Conformal manifolds and 3d mirrors of (D_n, D_m) theories*, *JHEP* **02** (2022) 014 [[arXiv:2110.06940](#)] [[INSPIRE](#)].
- [89] U. Derenthal, *Singular del Pezzo surfaces whose universal torsors are hypersurfaces*, *Proc. Lond. Math. Soc.* **108** (2013) 638.